

# Density Functional Calculations of the Mechanical, Electronic and Dynamical Properties of Antiperovskite $\text{Ca}_3\text{BO}$ (B = Pb, Ge, Sn)

Phillip Wilfsen Otieno Nyawere

Department of Physical and Biological Sciences, Kabarak University, Nakuru, Kenya

Email: pnyawere@gmail.com

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## Abstract

An analysis of mechanical, electronic and dynamical properties of antiperovskite  $\text{Ca}_3\text{BO}$  (B = Pb, Ge, Sn) in cubic phase space group Pm-3m (221) has been studied using first principle density functional theory (DFT). Ground state energy computation was done using the Projector Augmented Wave (PAW) Pseudo Potentials and the Plane Wave (PW) basis set. The Generalized Gradient Approximation (GGA) was used for the exchange correlation. The open source code QUANTUM ESPRESSO (QE) was used in this study in which plane wave basis sets are applied for the expansion of the electronic structure wave function. Thermo\_pw as a post-processing code was used for the computation of mechanical properties including bulk modulus and elastic constants with their derivatives. The lattice parameters are here calculated to be 4.87 Å, 4.86 Å and 4.84 Å for  $\text{Ca}_3\text{BO}$  (B = Pb, Ge, Sn) respectively which compares well with other works. This also shows that the three crystals are similar in size and in most of their properties. In addition to this, projected density of states and band structure are also computed both showing that these materials are of semi-metallic nature and are stable in cubic phase. Phonon modes at gamma are also reported.

## Keywords

Antiperovskite, Elastic Constants, Generalized Gradient Approximation

## 1. Introduction

Inverse Perovskites or antiperovskite compounds have two types of anion and cation positions reversed. Their general formula is given as  $\text{A}_3\text{BO}$  where A sites

are occupied by an electropositive ion while B sites are occupied by different anions whose size is different from A. Antiperovskite has displayed interesting properties which include magnetism, ionic conductivity, superconductivity and thermoelectricity, [1] among others. The discovery of superconducting properties of cubic antiperovskite  $\text{MgCN}_{13}$  compound [2] has motivated research on more of these antiperovskites. These findings have prompted a new round of development of functional antiperovskites. A new class of antiperovskite compounds  $\text{A}_3\text{BO}$  where  $\text{A} = \text{Ca}, \text{Sr}, \text{Ba}$  and  $\text{B} = \text{Pb}$  have been reported to crystallize into the cubic inverse perovskite structure. Because of the varied properties of Antiperovskite materials, interests have shifted to their unusual high charge carrier mobility which leads to many technological applications [3]. There has been in the recent past growing research interest in band topology, leading to a broad classification of non-trivial topological metals classified as Weyl Semimetal, Dirac Semimetal and Nodal line Semimetal (NLS)<sup>[4]</sup>. These classified metals have their conduction band and valence band crossing due to band inversion.

Three Dimensional (3D) Dirac fermion systems is a class of topological materials in which bulk conduction and valence bands with linear energy momentum dispersion relations meet at finite points or along curves in K-space. In order to verify the presence of 3D Dirac fermions theoretically, it is necessary to study electronic structure and in particular band structure to identify a Dirac point in 3D momentum space. The band structure of the inverse perovskite  $\text{Ca}_3\text{PbO}$ , through soft x-ray angle-resolved photoemission spectroscopy and cone like band dispersions are observed which is in close agreement with the predictions of electronic structure calculations [5].

The band structure calculation of  $\text{Ca}_3\text{BO}$  and its family using Wein2k have found that three-dimensional massive Dirac electrons exist at their Fermi energy. The present work is based on first principle calculations and demonstrates inverse perovskite phase of  $\text{Ca}_3\text{BO}$  ( $\text{B} = \text{Pb}, \text{Sn}, \text{Ge}$ ) as the new candidates for realizing Dirac fermionic system. An investigation of structural, electronic and dynamical properties of  $\text{Ca}_3\text{BO}$  is reported using DFT-based simulation code Quantum Espresso (QE). The electronic band structure and the projected density of states (PDOS) have also been studied. The present computed results are found to be in good agreement with the already available experimental and theoretical results.

The rest of this paper is arranged such that Section II is Computational methodology, Section III is discussing results and Section IV are Conclusions.

## 2. Computational Methodology

Calculations in this work were done in the framework of density functional theory using for the exchange-correlation functional, the generalized gradient approximation method of Perdew-Burke Ernzerhof (PBE) [6]. The pseudopotentials were extracted from Quantum espresso database [7]. The total energy convergence in the iterative solution of the Kohn-Sham equations [8] was set at

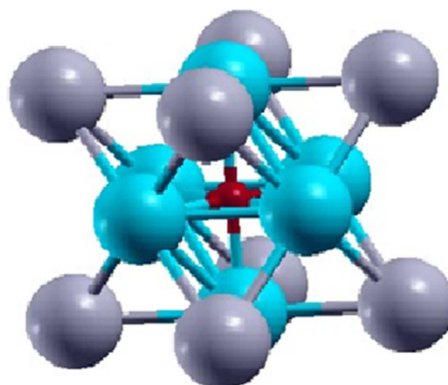
$2 \times 10^{-8}$  Ry. Optimized cell dimensions were fitted in Murnaghan fitting methodology of second order. The k-points and the kinetic energy cut-off values were properly optimized to convergence at the ground state energy. The valence configuration used for  $\text{Ca}_3\text{PbO}$  was 4s and 3p for Ca,  $6s^2$  for Pb and  $2s^2$  for O. Similar criteria is used for the other two compounds where Pb is replaced with Ge and Sn. The Brillouin sampling was based on the Monkhost scheme [9] [10] [11]. The K-point mesh in the irreducible high symmetry points in the Brillouin zone used was  $13 \times 13 \times 13$ . For the elastic constant calculation, the “quasi-static” approximation was used where the elastic constants were computed at zero absolute temperature and saved as elastic constants.

### 3. Results

#### 3.1. Structure Properties

**Figure 1** shows the optimized unit cell of  $\text{Ca}_3\text{PbO}$  with Wyckoff positions given as Ca (0, 0.5, 0.5), Ca (0.5, 0, 0.5), Ca (0.5, 0.5, 0), O (0.5, 0.5, 0.5) and Pb (0, 0, 0). The structures of  $\text{Ca}_3\text{SnO}$  and  $\text{Ca}_3\text{GeO}$  are all similar to that of  $\text{Ca}_3\text{PbO}$  with the position of Pb, Sn and Ge being the same. The convergence was achieved at a pressure of zero GPa. The equilibrium lattice constants, bulk modulus and derived bulk modulus have been obtained by minimizing the total crystal energy. This minimum energy was calculated for different values of lattice constants using second order Birch-Murnaghan equation of states. The calculated values are represented in **Table 1** for cubic phase of space group  $\text{Pm}\bar{3}\text{m}$  (#221) for  $\text{Ca}_3\text{BO}$  (B = Pb, Ge, Sn). The lattice constants are  $4.87\text{\AA}$ ,  $4.86\text{\AA}$ , and  $4.84\text{\AA}$  respectively for  $\text{Ca}_3\text{BO}$  (B = Pb, Ge, Sn) compounds which are found to compare well with other studies.

The three crystals show very close similarities in their sizes because of the similarity in the atoms in each crystal. The value of lattice parameters for all the three crystals are in agreement with other theoretical and experimental studies available. The values of bulk modulus and volume of single crystal also agrees



**Figure 1.** Optimized crystal structure of  $\text{Ca}_3\text{PbO}$  as visualized in xcrysden (Colour online) shows grey for O, light blue for Ca and red for Pb respectively. The central red at (0, 0, 0) is the position of Pb. It is this position that varies for  $\text{Ca}_3\text{SnO}$  and  $\text{Ca}_3\text{GeO}$  to be occupied by Sn and Ge respectively.

with the data of other studies.

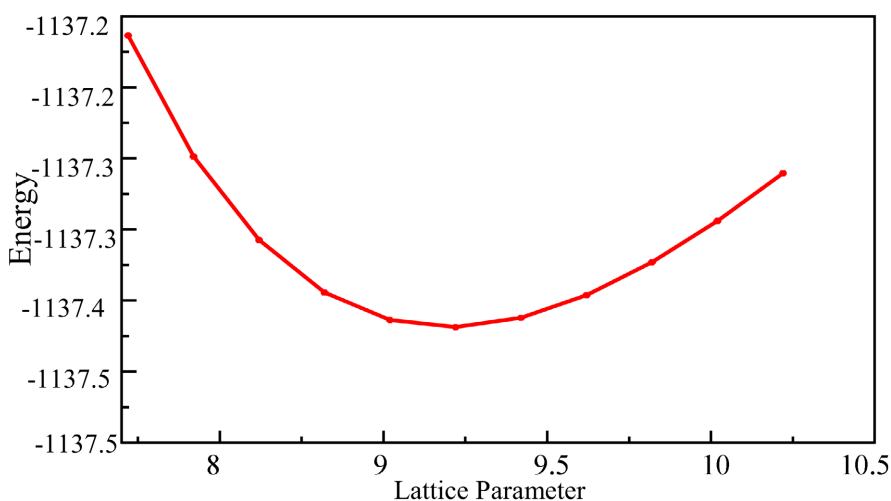
**Figure 2** shows the optimization of energy against lattice parameter. The lattice parameter was found to be 9.199 Bohr (4.87 Å) with volume of 115.35Å<sup>3</sup> for Ca<sub>3</sub>PbO. The optimization follows for all the three crystals.

### 3.2. Electronic Structure Properties

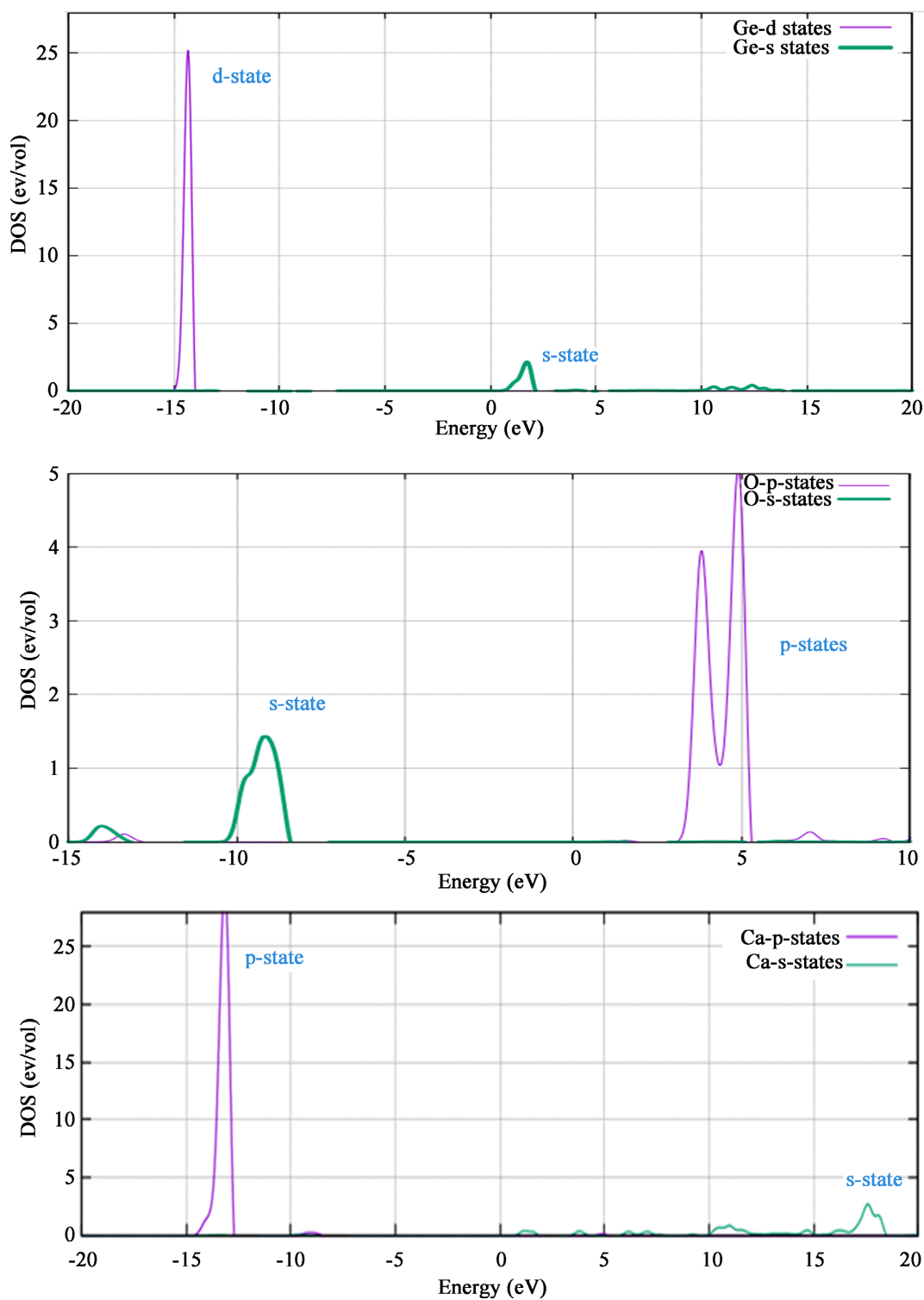
In this section, the partial density of states and band structures of Ca<sub>3</sub>BO (B = Pb, Ge, Sn) are discussed. The Fermi energy for all these crystals was set at zero and they all show semi-metallic behaviour. In **Figure 3**, Ca<sub>3</sub>PbO shows that above the Fermi level O-s states and Pb-s states are dominant even though other states also appear. Below the Fermi level Ca-s-p states are more dominant. Ca<sub>3</sub>GeO has Ge-s and Ca-s states in the conduction band and for Ca<sub>3</sub>SnO there are Sn-s and Ca-s states. It is clearly seen that the electrons of metallic atoms are majorly responsible for the conduction of energy in the compounds. Ca<sub>3</sub>PbO has

**Table 1.** Calculated lattice parameter (*a*), volume (*V*), bulk modulus (*B*) and pressure derivative of bulk modulus (*dB/dP*).

Material	<i>a</i> (Å)	<i>V</i> (Å <sup>3</sup> )	<i>B</i> (GPa)	<i>dB/dP</i>
Ca <sub>3</sub> PbO	4.87	115.35	51.79	4.33 (present)
	4.84 [12]	113.37	--	4.37 [14]
	4.73 [13]	105.82	64.59	
Ca <sub>3</sub> GeO	4.86	114.79	57.71	4.31 (present)
	4.73 [15]	105.7 [16]	57.71	
Ca <sub>3</sub> SnO	4.84	113.38	55.08	3.67 (present)
	4.86 [17]	114.5 [17]		
	4.83 [18]	112.88		



**Figure 2.** Optimized energy versus lattice parameter for Ca<sub>3</sub>PbO. The minimum energy corresponding to the lattice parameter is shown. The behaviour of convergence is similar for all the other two crystals.



**Figure 3.** Projected density of states in electrons per electron-volt for the compound  $\text{Ca}_3\text{BO}$ . S-states, d-states and p-states are labelled.

more electrons on the conduction band making it better conductor compared to the rest. This could be attributed to the position of Pb in the periodic table that makes it contribute more valence electrons.

From **Figure 3**, it is seen that s-states of Ca and Ge are dominant in the conduction band but p-states are the ones dominant for Oxygen in the conduction band. This is expected to be the case for  $\text{Ca}_3\text{SnO}$  and  $\text{Ca}_3\text{PbO}$  compounds where oxygen is expected to behave in the similar way as in  $\text{Ca}_3\text{GeO}$ .

From **Figure 4**, the band structures are plotted and all the three compounds show semi-metallic nature with direct band gap for all the compounds. The high symmetry points are all similar with strikingly similar band structures. The only differences is the Fermi energy of each of the compounds for example for Fermi energy for  $\text{Ca}_3\text{SnO}$  is 8.64 eV and that of  $\text{Ca}_3\text{PbO}$  is 9.16 eV.

### 3.3. Mechanical Properties

The mechanical properties of  $\text{Ca}_3\text{PbO}$  were also simulated using thermo\_pw code which is a post-processing code in quantum espresso to find the stability nature of this material. The cubic crystal has only three independent parameters which are  $C_{11}$ ,  $C_{12}$  and  $C_{44}$ . For a solid under strain the elastic energy is given as [19],

$$\frac{\Delta E}{V} = \frac{1}{2} \sum_{i=1}^6 \sum_{j=1}^6 c_{ij} e_i e_j \quad (1)$$

where  $V$  is the undistorted lattice cell volume,  $\Delta E$  is the change in energy from the strain vectors  $e = (e_1, e_2, e_3, e_4, e_5, e_6)$  and  $c$  is the matrix of the elastic constants [20]. The Born's stability criteria [21] at zero pressure requires that the elastic constants are related such that:  $C_{11} > 0$ ,  $C_{11} - C_{12} > 0$ ,  $C_{44} > 0$ ,  $C_{11} + 2C_{12} > 0$  and  $C_{11} > B > C_{12}$ . These constants are calculated by using the obtained second-order elastic constants of Equation (2); the results are given in **Table 2**.

$$C_{ij} = \frac{1}{V_0} \frac{\partial^2 E_{tot}}{\partial \xi_i \partial \xi_j} \quad (2)$$

By using Born's stability rule, it can be seen that these compounds are stable. This is important in mechanical application where pressure is to be applied on the crystal like in thermoelectric energy generation [22].

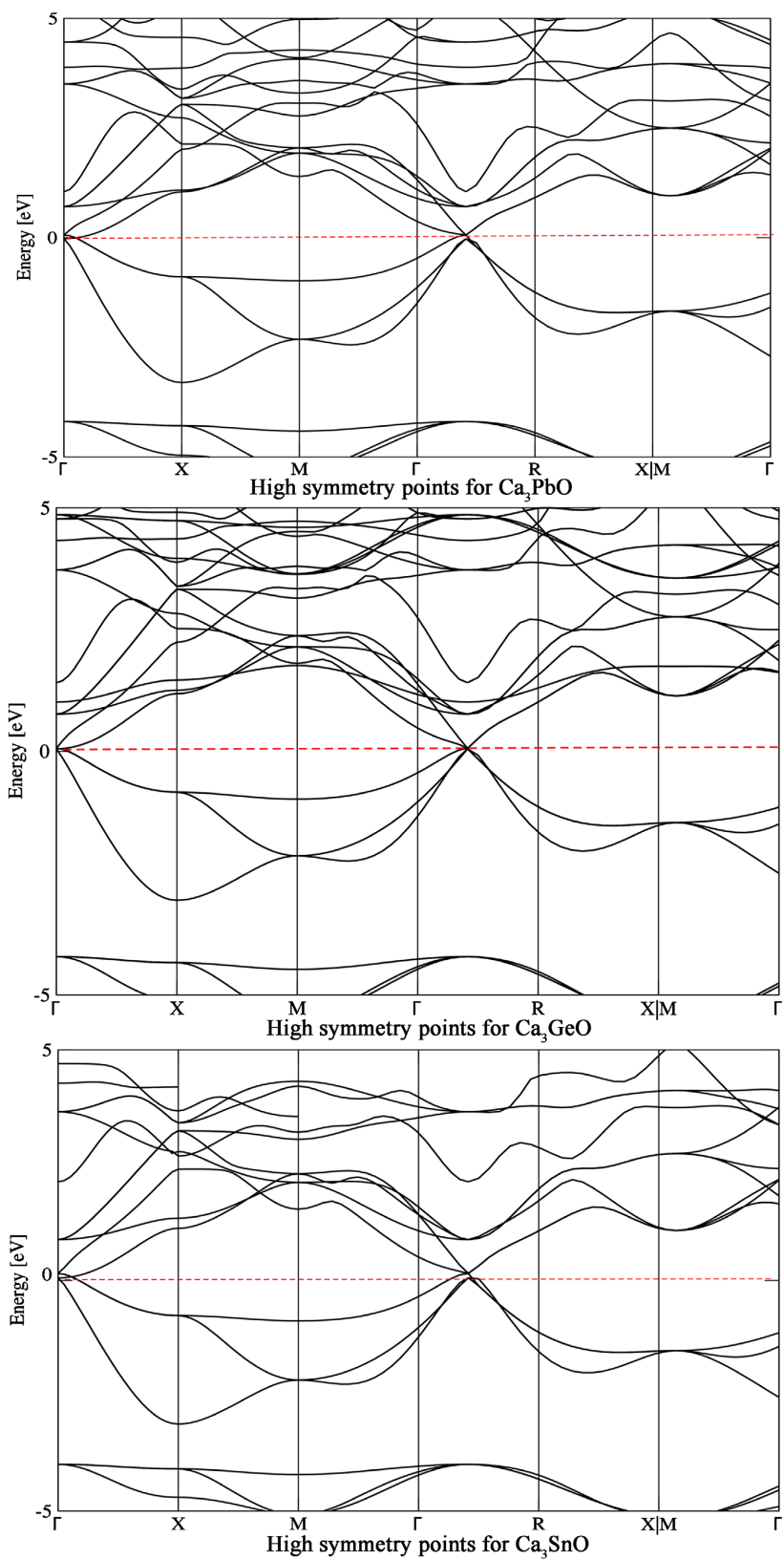
The Zener anisotropy factor ( $A$ ), Poisson ratio ( $\delta$ ), and Young's modulus ( $E$ ) that are important parameters in understanding elastic properties are also calculated using these formulas [12];

$$A = \frac{2C_{44}}{C_{11} - C_{12}} \quad (3)$$

$$\delta = \frac{1}{2} \left\{ \frac{B - \frac{2}{3}G}{B + \frac{1}{3}G} \right\} \quad (4)$$

$$E = \frac{9GB}{G + 3B} \quad (5)$$

where  $G$  is the isotropic shear modulus as a function of crystal orientation and is given as  $G = (G_V + G_R)/2$ , herein  $G_V$  is Voigt's shear modulus (which is related



**Figure 4.** Bands of antiperovskite  $\text{Ca}_3\text{PbO}$  against the high symmetry points. The Fermi energy has been set at the origin. The diagram shows semi-metallic band behaviour. This behaviour is also depicted for the  $\text{Ca}_3\text{SnO}$  and  $\text{Ca}_3\text{GeO}$  crystals.

**Table 2.** Elastic constants, bulk modulus and stability.

Material	$C_{11}$ (GPa)	$C_{12}$ (GPa)	$C_{44}$ (GPa)	$B$ (GPa)	Stability
Ca <sub>3</sub> PbO	97.68	28.71	49.09	51.7	Stable
Ca <sub>3</sub> GeO	118.1	27.51	49.53	57.7	Stable
Ca <sub>3</sub> SnO	103.32	30.96	50.49	55.08	Stable

**Table 3.** Voigt's shear modulus ( $G_V$ ), Reuss's Shear modulus ( $G_R$ ), Voigt's Young's modulus ( $E_V$ ), Reuss Young's modulus ( $E_R$ ), Poisson ratio ( $n$ ), Zener anisotropy factor ( $A$ ) and compressibility ( $\beta$ ).

	Ca <sub>3</sub> PbO	Ca <sub>3</sub> SnO	Ca <sub>3</sub> GeO
$G_V$ (GPa)	41.98	44.77	47.83
$G_R$ (GPa)	42.62	43.59	47.74
$G/B$	0.82	0.8	0.83
$E_V$ (GPa)	101.46	103.49	112.44
$E_R$ (GPa)	100.29	103.49	112.2
$n_{\max}$	0.181	0.186	0.175
$n_{\min}$	0.176	0.18	0.175
$A$	1.42	1.39	1.09
$\beta$ (GPa)	0.0064	0.0058	0.0061

to the upper bound of  $G$  values) and  $G_R$  is Reuss's shear modulus (which is related to the lower bound of  $G$  values) and can be written as  $G_V = (C_{11} - C_{12} + 3C_{44})/5$  and  $5/G_R = 4/(C_{11} - C_{12}) + 3/C_{44}$ , respectively.

**Table 3** shows calculated values of Voigt's isotropic shear modulus, Reuss's shear modulus, ratio of  $G/B$ , Young's modulus  $E$ , maximum and minimum Poisson's ration ( $n_{\max}$ ,  $n_{\min}$ ) and compressibility factor  $\beta$  for all the three crystals. It is seen that the isotropic shear modulus are such that the value of  $E_V$  and  $E_R$  increases from Ca<sub>3</sub>PbO, Ca<sub>3</sub>SnO and Ca<sub>3</sub>GeO respectively. The ratio  $G/B$  is called Pugh's ratio which shows whether a material is brittle or ductile. For  $G/B > 0.5$  shows brittle nature of the material but  $G/B < 0.5$  shows that the material is ductile. For the three crystals,  $G/B$  values calculated are 0.82, 0.80 and 0.83 for Ca<sub>3</sub>PbO, Ca<sub>3</sub>SnO and Ca<sub>3</sub>GeO respectively. This shows that these crystals are brittle. It can also be concluded that Ca<sub>3</sub>GeO has the largest Young's modulus which is an indication that it is stiffer compared to the other two compounds.

For Poisson's ratio less than 0.25, the material shows strong covalent bonds and value larger than 0.25 indicates ionic bond [23] [24] [25] [26] [27]. As can be seen from our computations all our Poisson's ratios are less than 0.25 which is an indication that all these atoms have strong covalent bonds. For bulk materials, elastic anisotropy is determined by Zener anisotropy factor  $A$  which if is 1, shows that the material is isotropic otherwise the material is anisotropic. All three materials show factor of  $A$  as 1.42, 1.39 and 1.09 for Ca<sub>3</sub>PbO, Ca<sub>3</sub>SnO and



$\text{Ca}_3\text{GeO}$  respectively. It can therefore be concluded that these compounds are anisotropic. This therefore implies that direction of pressure application is critical for maximum performance of these materials in thermoelectric design.

The compressibility is a measure of elasticity defined from the relation;

$$\beta = \frac{C_{11} - C_{12}}{\Omega} \quad (6)$$

$$\Omega = (C_{11} + C_{12})C_{11} - 2C_{12}^2 \quad (7)$$

The compressibility values are found to be  $0.0064 \text{ GPa}^{-1}$ ,  $0.0058 \text{ GPa}^{-1}$  and  $0.0061 \text{ GPa}^{-1}$  for  $\text{Ca}_3\text{PbO}$ ,  $\text{Ca}_3\text{SnO}$  and  $\text{Ca}_3\text{GeO}$  respectively. Compressibility is the inverse of bulk modulus and therefore high value of  $\beta$  means low value of bulk modulus.

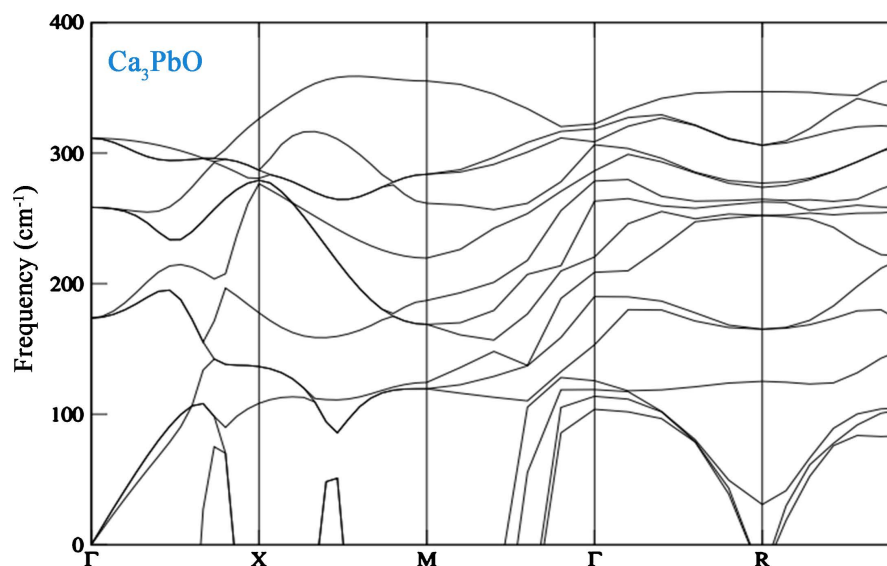
### 3.4. Dynamical Properties

Phonon dispersion relations are depicted in **Figures 5-7** below. The vibration frequency is directly proportional to the energy.

The primitive cell of the atoms has five atoms which gives 15 phonon branches with 3 acoustic and 12 optical modes for all the crystals. In all the crystals, there is no gap between the acoustic and optical branches of these modes. The energy of dispersion is for  $\text{Ca}_3\text{SnO}$  is the least followed by  $\text{Ca}_3\text{GeO}$  and then  $\text{Ca}_3\text{PbO}$ . This is not proportional to atomic mass unit of the metal in the crystal. **Figure 7** can be related to **Figure 4** where there is lack of overlap between the conduction and valence band. This shows that tin will require more energy for the conduction electrons to fully participate in conduction compared to the other two crystals.

## 4. Conclusions

Study of antiperovskite  $\text{Ca}_3\text{BO}$  (B = Pb, Ge, Sn) using density functional theory



**Figure 5.** Dispersion relation of  $\text{Ca}_3\text{PbO}$  with the highest frequency at about  $350 \text{ cm}^{-1}$ .

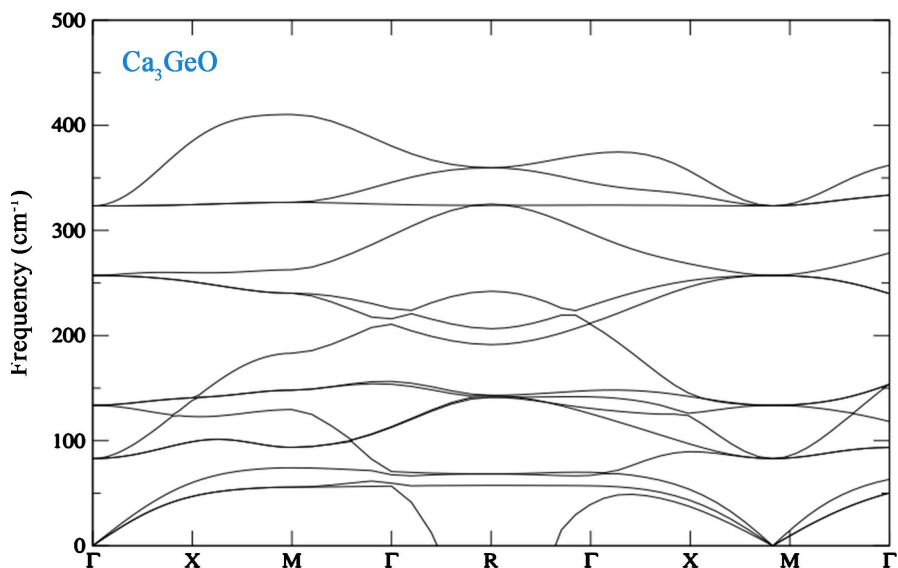


Figure 6. Dispersion relation of  $\text{Ca}_3\text{GeO}$  with highest frequency at about  $400\text{ cm}^{-1}$ .

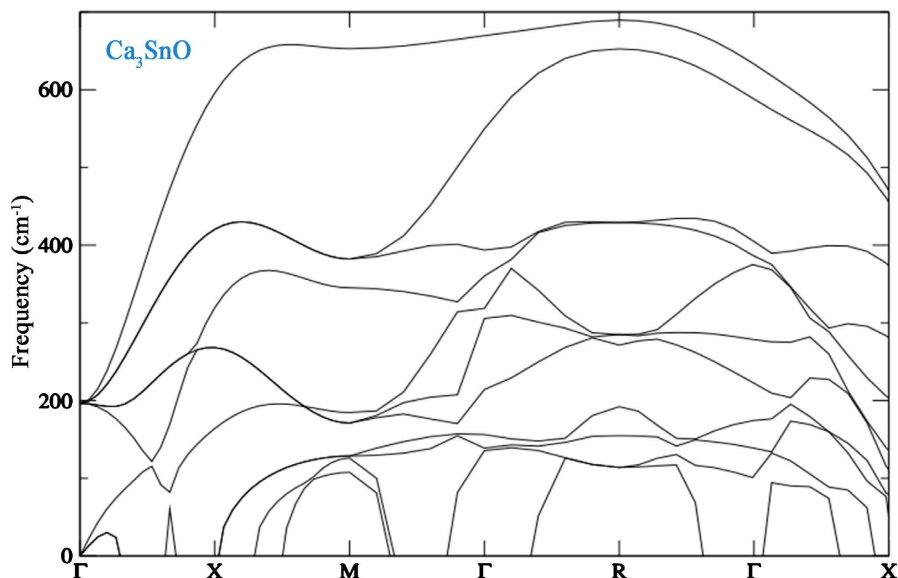


Figure 7. Dispersion relation of  $\text{Ca}_3\text{SnO}$  with the highest frequency at  $620\text{ cm}^{-1}$ .

is here reported. The electronic structure studies show that these materials are conductors with similarities in properties as depicted by their band structures and projected density of states. There are also indicating stability from elastic constants study and derived properties. Phonon dispersion also points at the metallic nature of compounds  $\text{CaBO}$ . These crystals are therefore stable in their cubic states and are suitable for pressure-induced applications including thermoelectric applications

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## Conflicts of Interest

The author declares no conflicts of interest regarding the publication of this paper.

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