

Antiproton Production with a Fixed Target and Search for Superheavy Particles at the LHC

Alexey B. Kurepin^{1*}, Nikolay A. Kurepin², Konstantin A. Skazytkin²

¹Institute for Nuclear Research, Moscow, Russia

²Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia

Email: *kurepin@inr.ru

How to cite this paper: Kurepin, A.B., Kurepin, N.A. and Skazytkin, K.A. (2022) Antiproton Production with a Fixed Target and Search for Superheavy Particles at the LHC. *Journal of Modern Physics*, 13, 1093-1098.

<https://doi.org/10.4236/jmp.2022.137062>

Received: May 9, 2022

Accepted: July 15, 2022

Published: July 18, 2022

Copyright © 2022 by author(s) and

Scientific Research Publishing Inc.

This work is licensed under the Creative Commons Attribution International

License (CC BY 4.0).

<http://creativecommons.org/licenses/by/4.0/>



Open Access

Abstract

A proposal for an experiment to measure the cross section of antiproton production in a proton-nuclear collision in a kinematically forbidden region for nucleon-nucleon interaction on a fixed LHC target is considered. It is shown that this process can be separated from the kinematically allowed production process using the existing detectors of the ALICE facility at a proton energy of 7 TeV with a fixed nuclear target. Assuming the scale dependence of the cross section, the data obtained can be used to estimate the subthreshold cross section for the production of superheavy particles with a mass of several tens of TeV in the LHC lead nucleus beam.

Keywords

Antiproton, Fixed Target, Collider, Heavy Ion Collision, Superheavy Particles

1. Introduction

Among the processes of particle production in nuclear collisions at high energy beams, one can highlight those where production occurs under kinematic conditions that are forbidden in the nucleon-nucleon interaction. At present, this problem is of great importance for studying the possibility of particle production with beams of lead nuclei for a mass exceeding 14 TeV, which is attainable in a nucleon-nucleon collision at the Large Hadron Collider (LHC). Actually, the energy at the center of mass in the collision of lead nuclei at the LHC at a beam energy of 2.76 TeV per nucleon is 1150 TeV. The existence of such superheavy particles with a mass much less than the Planck mass is forbidden in the Grand Unification Theory. However, when trying to solve the problem of the hierarchy of interactions in models of superstrings, super symmetry and introducing additional dimensions, the creation of such particles is allowed. For example, in the

model of extra dimensions with a compactification radius of one Fermi, particles with a mass of several tens of TeV can be produced [1] [2].

Obviously, for the production of superheavy particles, the cross section of the “subthreshold” process will be very small. Theoretical estimates of this cross section are apparently impossible at present. However, an analogy can be drawn with the subthreshold production of antiprotons at intermediate energies, since the subthreshold process is obviously associated with the correlation of nucleons or quarks in colliding nuclei and could be weakly dependent on energy. Nevertheless, such a study must also be carried out at higher energies.

This investigation can be carried out in the study of the production process under kinematic conditions that are forbidden in the nucleon-nucleon interaction. However, obtaining the necessary data in the collider operation mode is impossible. In this case, the produced antiprotons have an energy of several hundred GeV, and the spectrometry of antiprotons at such energies with existing detectors does not provide the required accuracy. As first shown in this article, it is possible to measure the production of antiprotons in a proton beam with an energy of 7 TeV at the LHC on a fixed target of heavy nuclei ($\sqrt{s} = 114.7$ GeV). Such a study was started at the U-70 accelerator in Protvino [3] at an energy of 19 GeV ($\sqrt{s} = 6.05$ GeV) per nucleon on carbon nuclei. However, the measurement of antiproton production was carried out at small forward angles, which leads to large antiproton momenta and less accuracy. The measurement of the transverse momentum of antiprotons proposed in this article makes it possible to obtain better accuracy at a much higher energy.

2. Subthreshold Production of Antiprotons in Nuclear Collisions

The study of the production of antiprotons in proton-nuclear and nucleus-nucleus collisions at energies below the production threshold in a nucleon-nucleon collision has been the subject of a significant number of works [4] [5] [6] [7] [8]. The measurements were carried out with proton and nuclear beams at JINR, BNL, GSI and KEK. In all experiments, the values of the production cross sections were obtained, which significantly exceeded the estimates obtained when taking into account the lowering of the production threshold due to the Fermi motion of nucleons in the nucleus [9]. A unified phenomenological description of all experimental data was obtained in the generalized parton model [10] [11]. The model is based on the introduction of the parton distribution parameter not only in the target nucleus (x), but also in the incident nucleus (z). Due to the conservation of the 4-momentum, we obtain the following results:

$$(zP_1 + xP_2 - P_d)^2 = (zP'_1 + xP'_2 + P_i)^2 \quad (1)$$

where P_1 , P_2 and P_d —are the 4-momenta of the incident and fixed nucleons in the nuclei and of the produced antiproton respectively, and P_i is the 4-momentum of an additional proton for conservation of the baryon number. For the maximum values of the momentum of the produced antiproton:

$$\vec{p}'_1 = \vec{p}'_2 = \vec{p}'_i = 0 \quad (2)$$

We have: $(P'_1)^2 = m^2$ and $(P'_2)^2 = m^2$, where m —is the mass of a nucleon, \vec{p}'_1 , \vec{p}'_2 , \vec{p}'_i momenta of incident and fixed nucleons in nuclei and an additional proton.

From Equation (1), the mutual dependence of the parameters x and z can be obtained [10]. The production of antiprotons in nucleus-nucleus collisions is now possible not only at small values of the parton parameters, but also in the kinematically forbidden region for nucleon-nucleon interactions at $x > 1$ and $z > 1$.

In reference [11], a universal dependence of the scaling type of all subthreshold data on the production of antiprotons was obtained at x in the range of 1 - 4 and at z values equal to 1 for incident protons, 1.3 for deuterons, 2 for carbon nuclei, and 3 for more heavy nuclei at energies from 2 to 6 GeV per nucleon:

$$(A_1 A_2)^{-0.43} \cdot E_1 \frac{d^3 \sigma}{dp^3} [\text{mb} \cdot \text{GeV}^{-2} \cdot \text{c}^3 \cdot \text{sr}^{-1}] = 0.57 \exp(-x/0.158) \quad (3)$$

where E_1 is the total energy of antiproton and A_1 and A_2 are the mass numbers of colliding nuclei.

This dependence has an exponential form and describes the data for proton-nuclear collisions well. For nucleus-nuclear collisions, the deviations of individual data from the curve are larger, but the range of cross sections reaches four orders of magnitude.

Assuming a weak dependence of the obtained dependence of the subthreshold production cross section on energy, the yield of superheavy particles with a mass of 16 TeV on the LHC proton beam in the collision of lead nuclei was estimated [12]. The obtained value of approximately 70 particles per year allows planning the corresponding experiment. However, for more reasonable estimates, it is necessary to determine the dependence of the production cross section on the scaling parameters in the subthreshold process, at values $x > 1$, at high energies closer to the LHC energies.

In the next section, we analyse the possibility of measuring the cross section for antiproton production in the kinematically forbidden region on a fixed target of the LHC collider.

3. Production of Antiprotons in a Kinematically Forbidden Region at a Fixed Target of the LHC Collider

To determine the possibility of measuring the yield of antiprotons in the kinematically forbidden region on a fixed target of the LHC collider [13] the kinematics of antiproton production at an energy of 7 TeV on bismuth nuclei were calculated. In **Figure 1**, the magnitudes of the maximum transverse momentum values of antiprotons in the center-of-mass system are given for the parameters $x = 1$ and $x = 2$, depending on the pseudorapidity. These magnitudes differ significantly. Thus, it is possible to separate the kinematically allowed process from the forbidden process in the nucleon-nucleon collision.

The identification and measurement of the transverse momentum in the ALICE installation is carried out by the TPC projection camera. The fixed target will be located at a distance of 480 cm from the IP. The geometry of the TPC and the beam tube limits the range of possible angles from 5 to 28 degrees. **Figure 2** and **Figure 3** show the corresponding intervals of pseudorapidity for $x = 1$ and $x = 2$. The required transverse momentum ranges of 3 - 20 GeV are available for TPC measurements [14].

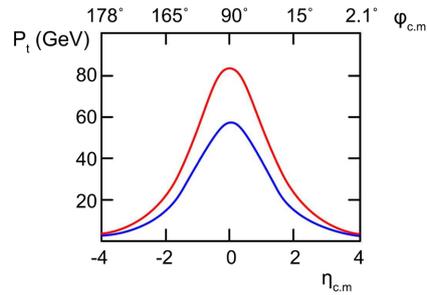


Figure 1. Dependence of the transverse momentum of antiprotons on the pseudorapidity in the center of mass. Blue line with $x = 1$, red line with $x = 2$.

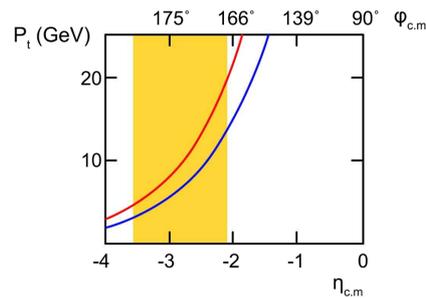


Figure 2. Dependence of the transverse momentum of antiprotons on the pseudorapidity in the center of mass. Blue line with $x = 1$, red line with $x = 2$. The area available with a fixed target and $x = 1$ is marked in yellow.

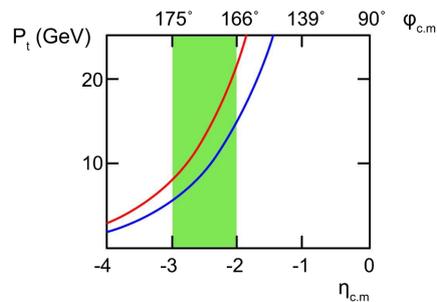


Figure 3. Dependence of the transverse momentum of antiprotons on the pseudorapidity in the center of mass. Blue line with $x = 1$, red line with $x = 2$. The area available with a fixed target and $x = 2$ is marked in green.

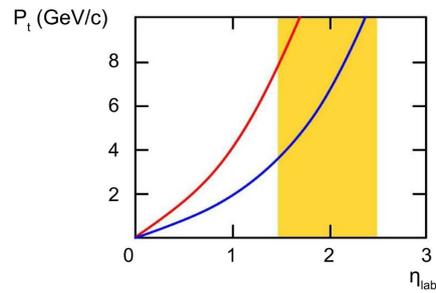


Figure 4. Dependence of the transverse momentum of antiprotons on the pseudorapidity in a laboratory system. Blue line with $x = 1$, red line with $x = 2$. The area available with a fixed target is marked in yellow.

Table 1. Parameter x , antiproton production cross sections and antiproton yield as a function of the antiproton transverse momentum.

P_t	4	6	8	GeV
E_d	8.5	12.8	17	GeV
x	1.1	1.6	2.18	-
σ_{inv}	8×10^{-3}	6×10^{-4}	8×10^{-6}	$\text{mb} \cdot \text{GeV}^{-2} \cdot \text{c}^3 \cdot \text{sr}^{-1}$
N_d	25×10^3	3×10^3	50	1/hour

Figure 4 shows the dependence of the transverse momentum on the pseudorapidity in the laboratory system.

When analysing the measured cross sections for the production of antiprotons, the parameters x and z must satisfy the following relation:

$$x = \frac{z \cdot E \cdot E_d \cdot (1 - \cos(\theta)) - M^2}{(z \cdot E - E_d)m} \tag{4}$$

where E is the beam energy, E_d is the antiproton energy, M is the antiproton mass, m is the nucleon mass and θ is the production angle.

At a high energy of 7 TeV of a proton beam, the parameter x is practically independent of the parameter z and beam energy:

$$x \approx \frac{E_d \cdot (1 - \cos(\theta))}{m} \tag{5}$$

For the parameters of the planned experiment: $\theta = 28^\circ$, $\Delta p = 1 \text{ GeV}$, $\Delta\Omega = 0.1 \text{ sr}$, the values of parameter x , as a function of the antiproton transverse momentum, are given in **Table 1**. The same table shows the production cross section calculated by formula (3) and antiproton yield at a luminosity of $10^{30} \text{ cm}^{-2} \cdot \text{cek}^{-1}$.

4. Conclusion

Investigation of antiproton production in the kinematically forbidden region in the nucleon-nucleon interaction on a fixed target of the LHC collider is possible with the existing ALICE facility detectors. The data obtained on the dependence

of the subthreshold production cross section on the scaling parameter $x > 1$ can be used to estimate the yield of superheavy particle production with the LHC lead nucleus beam. The results of the proposed measurements will allow studying the dependence of the scaling effect on energy, as well as obtaining new data on collective effects in nuclei.

Acknowledgements

We thank the members of the fixed-target ALICE study group and the AFTER study group for helpful discussions and interest in the work.

Conflicts of Interest

The authors declare no conflicts of interest regarding the publication of this paper.

References

- [1] Arkani-Hamed, N., Dimopoulos, S. and Dvali, G. (1998) *Physics Letters B*, **429**, 263. [https://doi.org/10.1016/S0370-2693\(98\)00466-3](https://doi.org/10.1016/S0370-2693(98)00466-3)
- [2] Rubakov, V.A. (2003) *Physics-Uspekhi*, **46**, 211. <https://doi.org/10.1070/PU2003v046n02ABEH001355>
- [3] Afonin, A.G., *et al.* (2020) *Physics of Atomic Nuclei*, **83**, 140. <https://doi.org/10.31857/S0044002720020014>
- [4] Baldin, A.A., *et al.* (1988) *JETP Letters*, **48**, 137-140.
- [5] Baldin, A.A., *et al.* (1990) *Nuclear Physics A*, **519**, 407-411. [https://doi.org/10.1016/0375-9474\(90\)90644-2](https://doi.org/10.1016/0375-9474(90)90644-2)
- [6] Carroll, J.B., *et al.* (1989) *Physical Review Letters*, **62**, 1829-1832. <https://doi.org/10.1103/PhysRevLett.62.1829>
- [7] Chiba, J., *et al.* (1993) *Nuclear Physics A*, **553**, 771-774. [https://doi.org/10.1016/0375-9474\(93\)90696-U](https://doi.org/10.1016/0375-9474(93)90696-U)
- [8] Schroeter, A., *et al.* (1993) *Nuclear Physics A*, **553**, 775-778. [https://doi.org/10.1016/0375-9474\(93\)90697-V](https://doi.org/10.1016/0375-9474(93)90697-V)
- [9] Shor, A., Perez-Mendez, V. and Ganezer, K. (1990) *Nuclear Physics A*, **514**, 717-733. [https://doi.org/10.1016/0375-9474\(90\)90019-I](https://doi.org/10.1016/0375-9474(90)90019-I)
- [10] Stavinski, V.S. (1988) JINR Rapid Comm. № 18-86, p. 5.
- [11] Kurepin, A.B., Shileev, K.A. and Topilskaya, N.S. (1997) *Genshiryoku Kenkyu, Tokyo*, **41**, 177-182.
- [12] Kurepin, A. (2021) *Journal of Modern Physics*, **12**, 433-439. <https://doi.org/10.4236/jmp.2021.124030>
- [13] Hadjidakis, C., *et al.* (2021) *Physics Reports*, **911**, 1-83. <https://doi.org/10.1016/j.physrep.2021.01.002>
- [14] Acharya, S., *et al.*, ALICE Collaboration (2021) *Journal of High Energy Physics*, No. 5, Article 290. [https://doi.org/10.1007/JHEP05\(2021\)290](https://doi.org/10.1007/JHEP05(2021)290)