

A Full Predictor-Corrector Finite Element Method for the One-Dimensional Heat Equation with Time-Dependent Singularities

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Abstract

The energy norm convergence rate of the finite element solution of the heat equation is reduced by the time-regularity of the exact solution. This paper presents an adaptive finite element treatment of time-dependent singularities on the one-dimensional heat equation. The method is based on a Fourier decomposition of the solution and an extraction formula of the coefficients of the singularities coupled with a predictor-corrector algorithm. The method recovers the optimal convergence rate of the finite element method on a quasi-uniform mesh refinement. Numerical results are carried out to show the efficiency of the method.

Keywords

Singularities, Finite Element Methods, Heat Equation, Predictor-Corrector Algorithm

1. Introduction

The heat equation is a partial differential equation known to model problems in thermodynamics, Brownian motions and particle diffusion. In the last seven decades, the increasing availability of powerful computers has been a major propulsion to the implementation of numerical methods for partial differential equations. Concerning the numerical solution of the heat equation, the usual methods are the finite difference methods (see [1] [2] [3]), the methods of lines (see [4] [5] [6]) and more recently the space-time finite element methods (FEMs) (see [7]-[12]).

The singular behavior of solutions of a partial differential equation is known to affect the convergence rate of the numerical approximation. The singularity can be geometric (due to the existence of reentrant corners or edges, the change of boundary conditions), or can be directly related to a singularity of a component of the equation (sign-changing coefficient, component singular at a point...). The goal of this paper is to use classical FEM to solve the one-dimensional heat problem: Find u such that:

$$\frac{\partial u}{\partial t} - k_w^2 \frac{\partial^2 u}{\partial x^2} = f \qquad \text{in} (0, L) \times (0, T]$$

$$u(0, t) = u(L, t) = 0 \qquad t \in (0, T) \qquad (1)$$

$$u(x, 0) = 0 \qquad x \in (0, L)$$

where L > 0, T > 0, $k_w > 0$ and the source term f(x,t) is some function singular at a point $t_0 \in [0,T]$.

The space-time FEM is known to give an approximation of the solution of Problem (1) with optimal convergence rates, but the method either uses discontinuous Galerkin methods (see [12]) that increase the number of degrees of freedom, or a particular interpolation error estimate in order to guarantee the optimal convergence (see [7]). Those techniques are efficient but involve particular adaptations' techniques.

Singular complement methods and predictor-corrector FEM (PCFEM) have been introduced to solve the problem of geometric singularities of elliptic boundary value problems (see [13]-[23]) and more recently on geometric singularities of hyperbolic problems (see [24]). The treatment of time-dependent singularities by means of PCFEM is also introduced in [25] using the method of lines. The main advantage of PCFEM is its application on any standard type of mesh (structured or unstructured) and on usual interpolation operators. This paper presents a PCFEM for the one-dimensional heat Problem (1). The method relies on a Fourier decomposition of the solution where computable formulas of the singular functions are derived and the solution is recomputed with a corrector algorithm in order to cancel the effect of the singularity caused by the source term f.

This paper is organized as follows: Section 2 presents the initial boundary value problem with the function spaces, the existence and uniqueness of the solution, a stability result and the Fourier decomposition of the solution with an extraction formula of the coefficients of the singularity. Section 3 presents the FEM, the predictor-corrector algorithm and some error estimates. Section 4 gives some numerical results that show the efficiency of the method and Section 5 is the conclusion of the paper.

2. Initial Boundary Value Problem

2.1. Function Spaces and Stability Results

In order to study Problem (1), it is crucial to introduce the following spaces:

1) Given a domain $\Omega \subset \mathbb{R}^n$, n = 1, 2, the space $L^2(\Omega)$ of classes of Lebesgue-measurable and square-integrable functions over Ω with the norm:

$$\left\|u\right\|_{L^{2}(\Omega)} = \left(\int_{\Omega} \left|u\left(\boldsymbol{x}\right)\right|^{2} \mathrm{d}\boldsymbol{x}\right)^{1/2}$$

where x can be a vector or a scalar depending on n.

2) Given an interval I = (a,b) and a Hilbert space V, $C^k(a,b;V)$ will denote the space of bounded *k*-continuously differentiable functions of the form $t \mapsto u(\cdot,t) \in V$ with the norm:

$$\left\|u\right\|_{\mathcal{C}^{k}}=\sup_{t\in I, |\beta|\leq k}\left\|D^{\beta}u\left(\cdot,t\right)\right\|_{V}$$

where $\beta = (\beta_1, \dots, \beta_n) \in \mathbb{N}^n$ is a multi-index, $|\beta| = \beta_1 + \dots + \beta_n$ and $D^{\beta} \varphi = \frac{\partial^{|\beta|} \varphi}{\partial x_1^{\beta_1} \cdots \partial x_n^{\beta_n}}.$

3) Given an interval I = (a,b), a Hilbert space V and $0 < \alpha < 1$, define the space $C^{k,\alpha}(a,b;V)$ be the space of k-continuously differentiable functions of the form $t \mapsto u(\cdot,t) \in V$ such that $D^{\beta}u(\cdot,t)$, $|\beta| = k$, is α -Hölder. The space $C^{k,\alpha}(a,b;V)$ is equipped with the norm:

$$\|u\|_{\mathcal{C}^{k,\alpha}} = \|u\|_{\mathcal{C}^{k}} + \max_{|\beta|=k} \sup_{a \le t_{1} < t_{2} \le b} \frac{\|D^{\beta}u(\cdot,t_{1}) - D^{\beta}u(\cdot,t_{2})\|_{V}}{|t_{1} - t_{2}|^{\alpha}}.$$

4) For a given domain Ω , the space $H^m(\Omega)$, $m \in \mathbb{N}$, $m \ge 1$ denotes the space of Lebesgue-measurable classes of functions u such that $u \in L^2(\Omega)$ and $D^{\beta}u \in L^2(0,L)$ for $0 < |\beta| \le m$, equipped with the norm:

$$\left\|u\right\|_{H^{m}(\Omega)} = \left(\int_{\Omega|\beta| \leq m} \left|D^{\beta}u\right|^{2} \mathrm{d}x\right)^{\frac{1}{2}}.$$

The space $H_0^1(\Omega) = \{ u \in H^1(\Omega) : u(\mathbf{x}) = 0 \text{ on } \partial\Omega \}$ where $\partial\Omega$ represents the boundary of Ω . The space $H_0^1(\Omega)$ is equipped with the norm of $H^1(\Omega)$.

The space $H^{-1}(\Omega)$ is the dual space of $H^1_0(\Omega)$ in the sense of distributions.

5) For an interval I = (a,b) and a domain Ω , the space $L^2(a,b;H_0^1(\Omega))$ is the completion of $\mathcal{C}(a,b;H_0^1(\Omega))$ with respect to the norm:

$$\left\|u\right\|_{L^{2}\left(a,b;H_{0}^{1}\left(\Omega\right)\right)}=\left(\int_{a}^{b}\left(\int_{\Omega\left|\beta\right|\leq1}\left|D^{\beta}u\left(\mathbf{x},t\right)\right|^{2}\mathrm{d}\mathbf{x}\right)\mathrm{d}t\right)^{\frac{1}{2}}.$$

The space $L^2(a,b;H^{-1}(\Omega))$ with the norm:

$$\|u\|_{L^{2}(a,b;H^{-1}(\Omega))} = \left(\int_{a}^{b} \|u\|_{H^{-1}(\Omega)}^{2} \mathrm{d}t\right)^{\frac{1}{2}},$$

where

$$\|u\|_{H^{-1}(\Omega)} = \sup_{v \in H^{1}_{0}(\Omega)} \left(\frac{|u(v)|}{\|v\|_{H^{1}_{0}(\Omega)}} \right).$$

The space $H^1(a,b;H^{-1}(\Omega))$ with the norm:

$$\left\|u\right\|_{H^{1}\left(a,b;H^{-1}\left(\Omega\right)\right)} = \left(\int_{a}^{b} \left(\left\|u\right\|_{H^{-1}\left(\Omega\right)}^{2} + \left\|\frac{\partial u}{\partial t}\right\|_{H^{-1}\left(\Omega\right)}^{2}\right) \mathrm{d}t\right)^{\frac{1}{2}}.$$

6) The space

$$X = \left\{ v \in L^{2}\left(0, T; H_{0}^{1}(0, L)\right) \cap H^{1}\left(0, T; H^{-1}(0, L)\right) : v(x, 0) = 0 \ \forall x \in (0, L) \right\}.$$

7) The space $Y = \left\{ v \in L^2(0,T; H^1_0(0,L)) : v(x,0) = 0 \ \forall x \in (0,L) \right\}.$

The variational formulation of Problem (1) is: Find $u \in X$ such that:

$$a(u,v) = \int_0^T \int_0^L f v dx \quad \forall v \in H_0^1(0,L)$$
(2)

where $a(u,v) = \int_0^T \int_0^L \frac{\partial u}{\partial t} v + k_w^2 \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} dx dt$.

Theorem 1. ([26] p. 382 (Theorem 11.3))

Assume that $f \in L^2(0,T;H^{-1}(0,L))$. Then, Problem (1) has a unique solution $u \in L^2(0,T;H_0^1(0,L)) \cap H^1(0,T;H^{-1}(0,L))$.

Remark. From the Sobolev embedding theorem (see [27], p. 98]), the space $H^1(0,T;H_0^1(0,L)) \subset C^{0,1/2}(0,T;H^{-1}(0,L))$, so the solution *u* of Problem (1) is at least $\frac{1}{2}$ -Hölder continuous in time.

Define the linear continuous map named as the Newton potential (see [7], p. 5]) $\mathcal{N}: L^2(0,T; H^{-1}(0,L)) \to L^2(0,T; H^1_0(0,L))$ by $\mathcal{N}\varphi$ such that $\|\varphi\|_{L^2(0,T; H^{-1}(0,L))} = k_w^2 \int_0^T \int_0^L \varphi \mathcal{N}\varphi \, dx \, dt$ and $\|\mathcal{N}\varphi\|_{L^2(0,T; H^1_0(0,L))} = \|\varphi\|_{L^2(0,T; H^{-1}(0,L))}$ for any $\varphi \in L^2(0,T; H^{-1}(0,L))$, if one uses the Cauchy-Schwartz inequality and the

Hölder inequality,

$$\begin{aligned} a(u,v) &= \int_{0}^{T} \int_{0}^{L} \frac{\partial u}{\partial t} v + k_{w}^{2} \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} dx dt \\ &\leq \left(\left\| \frac{\partial u}{\partial t} \right\|_{L^{2}(0,T;H^{-1}(0,L))} + k_{w}^{2} \left\| u \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))} \right) \right\| v \|_{L^{2}(0,T;H^{1}_{0}(0,L))} \\ &\leq \left(\sqrt{2} + k_{w}^{2} \right) \left(\left\| \frac{\partial u}{\partial t} \right\|_{L^{2}(0,T;H^{-1}(0,L))}^{2} + \left\| u \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))}^{2} \right)^{1/2} \left\| v \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))} \\ &\leq \left(\sqrt{2} + k_{w}^{2} \right) \left\| u \right\|_{L^{2}(0,T;H^{1}_{0}(0,L)) \cap H^{1}(0,T;H^{-1}(0,L))} \left\| v \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))}. \end{aligned}$$

Hence, $a(\cdot, \cdot)$ is bounded and

$$\|u\|_{L^{2}(0,T;H_{0}^{1}(0,L))\cap H^{1}(0,T;H^{-1}(0,L))}^{2} = \left\|\frac{\partial u}{\partial t}\right\|_{L^{2}(0,T;H^{-1}(0,L))}^{2} + \|u\|_{L^{2}(0,T;H_{0}^{1}(0,L))}^{2}$$

$$= \left\|\mathcal{N}\frac{\partial u}{\partial t}\right\|_{L^{2}(0,T;H_{0}^{1}(0,L))}^{2} + \|u\|_{L^{2}(0,T;H_{0}^{1}(0,L))}^{2}$$
(3)

so if $u \in X$, one has:

$$\begin{aligned} a(u,u) &= \int_0^T \int_0^L \frac{\partial u}{\partial t} u + k_w^2 \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} dx dt \\ &= \frac{1}{2} \int_0^T \frac{d}{dt} \|u\|_{L^2(0,L)}^2 dt + k_w^2 \|u\|_{L^2(0,T;H_0^1(0,L))}^2 \\ &= \frac{1}{2} \|u(\cdot,T)\|_{L^2(0,L)}^2 + k_w^2 \|u\|_{L^2(0,T;H_0^1(0,L))} \ge k_w^2 \|u\|_{L^2(0,T;H_0^1(0,L))} \end{aligned}$$

and using the definition of \mathcal{N} , the Cauchy-Schwartz inequality and (3),

$$\begin{aligned} a\left(u, \mathcal{N}\frac{\partial u}{\partial t}\right) &= \int_{0}^{T} \int_{0}^{L} \frac{\partial u}{\partial t} \mathcal{N}\frac{\partial u}{\partial t} + k_{w}^{2} \frac{\partial u}{\partial x} \frac{\partial}{\partial x} \left(\mathcal{N}\frac{\partial u}{\partial t}\right) dx dt \\ &\geq \frac{1}{k_{w}^{2}} \left\|\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T;H^{-1}(0,L)\right)}^{2} - k_{w}^{2} \left\|u\right\|_{L^{2}\left(0,T;H_{0}^{1}(0,L)\right)} \left\|\mathcal{N}\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T;H_{0}^{1}(0,L)\right)} \\ &\geq \frac{1}{k_{w}^{2}} \left\|\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T;H^{-1}(0,L)\right)}^{2} - \frac{k_{w}^{2}}{2} \left(\left\|u\right\|_{L^{2}\left(0,T;H_{0}^{1}(0,L)\right)}^{2} + \left\|\mathcal{N}\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T;H_{0}^{1}(0,L)\right)}^{2} \right) \end{aligned}$$
(4)
$$&= \frac{1}{k_{w}^{2}} \left\|\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T;H^{-1}(0,L)\right)}^{2} - \frac{k_{w}^{2}}{2} \left\|u\|_{L^{2}\left(0,T;H_{0}^{1}(0,L)\right)}^{2} \right. \end{aligned}$$

Hence, from (3) and (4),

$$\begin{aligned} a\left(u, u + \mathcal{N}\frac{\partial u}{\partial t}\right) &= a\left(u, u\right) + a\left(u, \mathcal{N}\frac{\partial u}{\partial t}\right) \\ &\geq k_w^2 \left\|u\right\|_{L^2\left(0, T; H_0^1(0, L)\right)}^2 + \frac{1}{k_w^2} \left\|\frac{\partial u}{\partial t}\right\|_{L^2\left(0, T; H^{-1}(0, L)\right)}^2 - \frac{k_w^2}{2} \left\|u\right\|_{L^2\left(0, T; H_0^1(0, L)\right)}^2 \\ &\geq \frac{k_w^2}{2} \left\|u\right\|_{L^2\left(0, T; H_0^1(0, L)\right)}^2 + \frac{1}{k_w^2} \left\|\frac{\partial u}{\partial t}\right\|_{L^2\left(0, T; H^{-1}(0, L)\right)}^2 \\ &\geq \max\left(\frac{k_w^2}{2}, \frac{1}{k_w^2}\right) \left\|u\right\|_{L^2\left(0, T; H_0^1(0, L)\right)}^2 \cap H^1(0, T; H^{-1}(0, L)) \end{aligned}$$

and from (3),

$$\begin{aligned} \left\| u + \mathcal{N} \frac{\partial u}{\partial t} \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))}^{2} &\leq 2 \left[\left\| u \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))}^{2} + \left\| \mathcal{N} \frac{\partial u}{\partial t} \right\|_{L^{2}(0,T;H^{1}_{0}(0,L))}^{2} \right] \\ &= 2 \left\| u \right\|_{L^{2}(0,T;H^{1}_{0}(0,L)) \cap H^{1}(0,T;H^{-1}(0,L))}^{2}. \end{aligned}$$

Then,

$$a\left(u, u + \mathcal{N}\frac{\partial u}{\partial t}\right) \\ \geq \frac{1}{\sqrt{2}} \max\left(\frac{k_{w}^{2}}{2}, \frac{1}{k_{w}^{2}}\right) \|u\|_{L^{2}\left(0,T; H^{1}_{0}(0,L)\right) \cap H^{1}\left(0,T; H^{-1}(0,L)\right)} \left\|u + \mathcal{N}\frac{\partial u}{\partial t}\right\|_{L^{2}\left(0,T; H^{1}_{0}(0,L)\right)}$$
(5)

and one has the stability condition:

$$\|u\|_{L^{2}(0,T;H^{1}_{0}(0,L))\cap H^{1}(0,T;H^{-1}(0,L))} \leq \sup_{v\in L^{2}(0,T;H^{1}_{0}(0,L)), v\neq 0} \frac{a(u,v)}{\|v\|_{L^{2}(0,T;H^{1}_{0}(0,L))}}.$$
(6)

2.2. Fourier Decomposition of the Solution

Assuming that the solution u of (1) is an element of $C^{0,\alpha}(0,T;H_0^1(0,L))$, $1/2 \le \alpha < 1$, and has a line singularity at a point $t_0 \in [0,T]$, based on the boundary conditions one can set

 $u(x,t) = \sum_{k=1}^{\infty} v_k(t^{\alpha}) \sin(k\pi x/L) := \sum_{k=1}^{\infty} \gamma_k(t) \sin(k\pi x/L)$ and obtain use separation of variables to obtain the following result.

Theorem 2. The solution *u* of Problem (1) can be written as

 $u(x,t) = \sum_{k=1}^{\infty} \gamma_k(t) s_k(x) \text{ where } s_k(x) = \sin(k\pi x/L) \text{, and setting } \beta_k = \frac{k\pi}{L} k_w,$ 1) If $t_0 = 0$,

$$\gamma_{k}\left(t\right) = \left[\int_{0}^{t^{\alpha}} \frac{1}{\alpha} f_{k}\left(\tau^{1/\alpha}\right) \tau^{\frac{1-\alpha}{\alpha}} \mathrm{e}^{\beta_{k}^{2}\tau^{1/\alpha}} \mathrm{d}\tau\right] \mathrm{e}^{-\beta_{k}^{2}t}$$
(7)

and $\frac{\mathrm{d}\gamma_{k}}{\mathrm{d}t}(t) = f_{k}(t) - \beta_{k}^{2}\gamma_{k}(t).$ 2) If $t_{0} > 0$,

$$\gamma_{k}\left(t\right) = \left[\int_{\left(t_{0}-t\right)^{\alpha}}^{t^{\alpha}} \frac{1}{\alpha} f_{k}\left(t_{0}-\tau^{1/\alpha}\right) \tau^{\frac{1-\alpha}{\alpha}} \mathrm{e}^{-\beta_{k}^{2} \tau^{1/\alpha}} \mathrm{d}\tau\right] \mathrm{e}^{\beta_{k}^{2}\left(t_{0}-t\right)}$$
(8)

and $\frac{d\gamma_k}{dt}(t) = f_k(t) - \beta_k^2 \gamma_k(t)$.

3. Finite Element Discretization

3.1. Finite ELement Method

Let \mathcal{T}_h be a triangulation of the domain $(0, L) \times (0, T)$ with quasi-uniform meshes of diameter *h* and define:

$$X_{h} = \left\{ v \in X : v_{|K} \in \mathcal{P}^{1}\left(\left(0, L \right) \times \left(0, T \right) \right) \forall K \in \mathcal{T}_{h} \right\}$$

with the norm of $L^{2}(0,T;H_{0}^{1}(0,L)) \cap H^{1}(0,T;H^{-1}(0,L))$.

The Galerkin approximation consists to find $u_h \in X_h$ such that:

$$a(u_h, v_h) = \int_0^T \int_0^L f v_h \, \mathrm{d}x \, \mathrm{d}t \quad \forall v \in X_h.$$
(9)

The variational Problem (9) has a unique solution according to Theorem 1 and since $X_h \subset X \subset L^2(0,T; H_0^1(0,L)) \cap H^1(0,T; H^{-1}(0,L))$, for any $u_h \in X_h$, from the inequality (5),

$$\begin{aligned} u_{h} \|_{X_{h}} &= \left\| u_{h} \right\|_{L^{2}\left(0,T;H^{1}_{0}(0,L)\right) \cap H^{1}\left(0,T;H^{-1}(0,L)\right)} \\ &\leq C \frac{a\left(u_{h}, u_{h} + \mathcal{N} \frac{\partial u_{h}}{\partial t} \right)}{\left\| u_{h} + \mathcal{N} \frac{\partial u_{h}}{\partial t} \right\|_{L^{2}\left(0,T;H^{1}_{0}(0,L)\right)}} \\ &= C \frac{a\left(u_{h}, u_{h} + \mathcal{N} \frac{\partial u_{h}}{\partial t} \right)}{\left\| u_{h} + \mathcal{N} \frac{\partial u_{h}}{\partial t} \right\|_{X_{h}}} \end{aligned}$$
(10)
$$\leq C \sup_{v_{h} \in X_{h}, v_{h} \neq 0} \frac{a\left(u_{h}, v_{h} \right)}{\left\| v_{h} \right\|_{X_{h}}} \end{aligned}$$

where C > 0 is a constant depending on k_w . Furthermore, one has the Galerkin orthogonality:

$$a(u-u_h,v_h) = 0 \quad \forall v_h \in X_h \tag{11}$$

The following theorem is a lemma equivalent to the Céa's Lemma known for elliptic boundary value problems. This theorem has been proved in [7] in a more general case.

Theorem 3. Let *u* be the solution of the variational Problem (2) and $u_h \in X_h$ be the solution of the discretized variational Problem (9). Then, there exists C > 0 depending on k_w such that:

$$\|u - u_h\|_{X_h} \le C \inf_{w_h \in X_h} \|u - w_h\|_X$$

Proof. From the discrete stability condition (10) and the Galerkin orthogonality (11), for any $w_h \in X_h$,

$$\begin{split} \|u_{h} - w_{h}\|_{X_{h}} &\leq C \sup_{v_{h} \in X_{h}} \frac{a(u_{h} - w_{h}, v_{h})}{\|v_{h}\|_{L^{2}(0,T;H_{0}^{1}(0,L))}} \\ &\leq C \Biggl(\sup_{v_{h} \in X_{h}} \frac{a(u_{h} - u, v_{h})}{\|v_{h}\|_{L^{2}(0,T;H_{0}^{1}(0,L))}} + \sup_{v_{h} \in X_{h}} \frac{a(u - w_{h}, v_{h})}{\|v_{h}\|_{L^{2}(0,T;H_{0}^{1}(0,L))}} \Biggr) \\ &= C \sup_{v_{h} \in X_{h}} \frac{a(u - w_{h}, v_{h})}{\|v_{h}\|_{L^{2}(0,T;H_{0}^{1}(0,L))}} \\ &\leq C \sup_{v \in X} \frac{a(u - w_{h}, v)}{\|v\|_{L^{2}(0,T;H_{0}^{1}(0,L))}} \\ &\leq C \|u - w_{h}\|_{X} \,. \end{split}$$

Then, $\|u - u_h\|_{X_h} \leq \|u - w_h\|_{X_h} + \|w_h - u_h\|_{X_h} \leq (1 + C)\|u - w_h\|_X$ for any $w_h \in X_h$ where C > 0 is an arbitrary constant depending on k_w . Hence, $\|u - u_h\|_{X_h} \leq C \inf_{w_h \in X_h} \|u - w_h\|_X$.

The following theorem proves the energy error estimate when the solution is smooth. This theorem has been proved in [7] in a more general case.

Theorem 4. Let u be the unique solution of Problem (1) and u_h be the unique solution of the equivalent discretized variational Problem (9). Assume $u \in H^2((0,L)\times(0,T))$, then

$$\|u - u_h\|_{L^2(0,T;H^1_0(0,L))} \le Ch \|u\|_{H^2((0,L)\times(0,T))}.$$

Proof.

$$\|u - u_h\|_{L^2(0,T;H^1_0(0,L))} = \|u - u_h\|_X \le C \inf_{w_h \in X_h} \|u - w_h\|_X \le C \|u - \Pi_h u\|_X$$

where $\Pi_h u \in X_h$ is the unique solution of the variational problem,

$$a(\Pi_h u, v_h) = a(u, v_h) \quad \forall v_h \in X_h.$$

Note that for any $v \in H^1((0,L)\times(0,T))$ such that v(x,0)=0 and v(x,t)=0 for x=0 or x=L,

$$\|v\|_{H^{1}((0,L)\times(0,T))}^{2} = \int_{0}^{T} \int_{0}^{L} |v|^{2} + \left|\frac{\partial v}{\partial t}\right|^{2} + \left|\frac{\partial v}{\partial x}\right|^{2} dx dt \ge \|v\|_{L^{2}(0,T;H^{1}_{0}(0,L))}^{2}.$$

Then,

$$\begin{split} \|u - u_h\|_{L^2(0,T;H^1_0(0,L))} &\leq C \|u - \Pi_h u\|_{L^2(0,T;H^1_0(0,L))} \\ &\leq C \|u - u_h\|_{H^1((0,L)\times(0,T))} \\ &\leq C \inf_{w_h \in X_h} \|u - v_h\|_{H^1((0,L)\times(0,T))} \\ &\leq Ch \|u\|_{H^2((0,L)\times(0,T))} \,. \end{split}$$

3.2. Predictor-Corrector Algorithm

The predictor-corrector method presented is based on the decomposition of the approximate solution $u_{hN}(x,t) = w_{hN}(x,t) + \sum_{k=1}^{N} \gamma_k(t) s_k(x)$ where u_{hN} is the solution of (9), s_k and γ_k are defined as in Theorem 2.

Predictor-Corrector Algorithm

1) Solve the problem: Find u_h such that:

$$\frac{\partial u_h}{\partial t} - k_w^2 \frac{\partial^2 u_h}{\partial x^2} = f \qquad \text{in} (0, L) \times (0, T]$$
$$u_h(0, t) = u_h(L, t) = 0 \qquad t \in [0, T]$$
$$u_h(x, 0) = 0 \qquad x \in (0, L)$$

2) For $N \in \mathbb{N}$ fixed, do:

a) For $k = 1, \dots, N$, compute $\gamma_k(t)$ and $\frac{d\gamma_k}{dt}(t)$ using Formula (7) or (8). b) If

$$X_{1h} := X_h + \operatorname{Span}\{s_k, k = 1, \cdots, N\},\$$

solve the problem: Find $w_{hN} \in X_{1h}$ such that:

$$\frac{\partial w_{hN}}{\partial t} - k_w^2 \frac{\partial^2 w_{hN}}{\partial x^2} = f - \sum_{k=1}^N \frac{\mathrm{d}\gamma_k}{\mathrm{d}t} s_k - k_w^2 \gamma_k \frac{\mathrm{d}^2 s_k}{\mathrm{d}x^2} \quad \text{in} (0, L) \times (0, T]$$

$$w_{hN} (0, t) = -\sum_{k=1}^N \gamma_{k\tau} (t) s_k (0) \qquad t \in [0, T]$$

$$w_{hN} (L, t) = -\sum_{k=1}^N \gamma_k (t) s_k (L) \qquad t \in [0, T]$$

$$w_{hN} (x, 0) = -\sum_{k=1}^N \gamma_k (0) s_k (x) \qquad x \in [0, L]$$

c) Compute $u_{hN}(x,t) = w_{hN}(x,t) + \sum_{k=1}^{N} \gamma_k(t) s_k(x)$.

3.3. Error Estimates

From [28], if $u \in \mathcal{C}^{0,\alpha}(0,T; H_0^1(0,L))$, $0 < \alpha < 1$, then for any $N \in \mathbb{N}$, $N \ge 1$, $\left\| u - \sum_{k=1}^N \gamma_k S_k \right\|_{L^2(0,T; H_0^1(0,L))} \le C N^{-\alpha} \ln N$ (12)

where γ_k and s_k are defined in Theorem 2 and C > 0 is an arbitrary constant.

The convergence result in the energy norm is given by the following theorem. **Theorem 5.** Let $u \in C^{0,\alpha}(0,T; H^1(0,L))$, $1/2 \le \alpha < 1$ be the solution of Problem (1) and let u_{hN} , $N \in \mathbb{N}$, $N \ge 1$ be the approximate solution of u using the Predictor-Corrector Algorithm 3.2 on a quasi-uniform and triangulation \mathcal{T}_h of the domain $(0,L) \times (0,T)$. Assume further that $N^{-\alpha} \ln N \le h$, then there is C > 0 such that:

$$\|u - u_h\|_{L^2(0,T;H^1_0(0,L))} \le Ch$$

Proof. Suppose $u(x,t) = w_N(x,t) + \sum_{k=1}^N \gamma_k(t) s_k(x)$ is the decomposition of the solution u as in Theorem 2. Set $u_{hN}(x,t) = w_{hN} + \sum_{k=1}^N \gamma_k(t) s_k(x)$ be the approximation of u using the Predictor-Corrector Algorithm 3.2.

From Theorem 4, (12) and the assumption $N^{-\alpha} \ln N \le h$, one has:

$$\|u - u_{hN}\|_{L^{2}(0,T;H^{1}_{0}(0,L))} = \|w_{N} - w_{hN}\|_{L^{2}(0,T;H^{1}_{0}(0,L))} \le CN^{-\alpha} \ln N \le Ch.$$

4. Numerical Experiments

The main purpose of this section is to present some computational results carried out with the Predictor-Corrector Algorithm 3.2. Even though the examples considered may not represent any real world situation, the difficulties due to singularities in the solution are well reflected. We consider two examples with known solutions. All computations are carried out using a code developped in python. The resulting linear systems are solved iteratively with a generalized minimal residual method (GMRES) and all the integrals are computed using a 15-point Legendre Gauss formula on the stifness, mass matrices and load vectors, and a 150-points Legendre Gauss formula for integrals used to compute the Fourier coefficients.

In all the computations, starting from a triangulation T_{h_0} , the refinement are perfomed by dividing the concerned lentgh by two ($h_{i+1} = h_i/2$). The order or rate of convergence v is computed using the formula:

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$$\nu = \log_2 \left(\frac{\left\| u - u_{h_i} \right\|_{L^2(0,T;H^1_0(0,L))}}{\left\| u - u_{h_{i+1}} \right\|_{L^2(0,T;H^1_0(0,L))}} \right).$$

In all the following examples, v_{exp} is the expected order of convergence. Figure 1 presents the type of mesh refinements used in the examples.

4.1. Example 1

In this example, we consider Problem (9) with the exact solution

$$u(x,t) = t^{1/2}x(x-1)$$
 and $f = \frac{\partial u}{\partial t} - \frac{1}{144}\frac{\partial^2 u}{\partial x^2}$, $T = 1$, $L = 1$.

The function has a line singularity at t = 0 and belongs to $C^{0,1/2}(0,1,H^1(0,1))$. **Table 1** presents the error estimates in the energy norm $\|\cdot\|_{L^2(0,T;H^1_0(0,L))}$ and the required number of Fourier coefficients *N* in order to obtain an optimal convergence



Figure 1. The uniform meshes: Refinements 1, 2 and 3

Τ	a	b]	le	1.	Error	estimates	in	the	energy	norm.
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au = h	2^{-1}	2^{-2}	2^{-3}	2^{-4}	2 ⁻⁵	V_{exp}
# Nodes	9	25	81	289	1089	
$\ u-u_h\ _{L^2(0,T;H^1_0(0,L))}$	$2.44e^{-1}$	$1.457e^{-1}$	9.068e ⁻²	5.836e ⁻²	3.855e ⁻²	
Order		0.745	0.684	0.636	0.598	0.25
N	75	681	4538	26570	144557	
$\ u-u_{hN}\ _{L^2(0,T;H^1_0(0,L))}$	$4.250e^{-2}$	2.036e ⁻²	$1.018e^{-2}$	5.096e ⁻³	$2.640e^{-3}$	
Order		1.062	0.999	0.998	1.004	1.00

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rate. Note that we need $N^{-1/2} \ln N \le h_j$ at each refinement $j = 1, \dots, 5$. A convergence of the classical FEM is observed and the convergence rate that is greater than 0.5, which is better than the expected convergence rate 0.25. However the convergence rate of the classical FEM is still not optimal. From the results obtained with the PCFEM, the optimal convergence rate is recovered.

Figures 2-4 present the plots of the exact solution, the approximate solution before adaptation and the aproximate solution using the Predictor-Corrector Algorithm 3.2 respectively at the fixed times t = 1/4, t = 1/2, t = 3/4 and t = 1. **Figures 5-7** present the 3-dimensional line countour plots of the exact solution, the approximate solution before adaptation and the approximate solution after using the predictor-corrector adaptation (PCFEM) with 75 lines. It is observed that the predictor-corrector solution has a better approximation with a least number of nodes in the refinement.

4.2. Example 2

In this example, we consider Problem (9) with the exact solution

$$u(x,t) = (t-1)^{3/4} x(x-1)$$
 and $f = \frac{\partial u}{\partial t} - \frac{1}{400} \frac{\partial^2 u}{\partial x^2}$, $T = 1$, $L = 1$

The solution has a line singularity at $t_0 = 1$ and is an element of $C^{0,3/4}(0,1,H^1(0,1))$ so the number of Fourier coefficients at each refinement on a triangulation of diameter *h* should satisfy $N^{-3/4} \ln N \le h$. Figure 8 shows the



Figure 2. The exact solution, the FEM solution and the PCFEM solution at Refinement 3 ($h = 2^{-3}$).



Figure 3. The exact solution, the FEM solution and the PCFEM solution at Refinement 4 ($h = 2^{-4}$).



Figure 4. The exact solution, the FEM solution and the PCFEM solution at Refinement 5 ($h = 2^{-5}$).



Figure 5. 3D plots of the exact solution, the FEM solution and the PCFEM solution at Refinement 3.



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Figure 6. 3D plots of the exact solution, the FEM solution and the PCFEM solution at Refinement 4.



Figure 7. 3D plots of the exact solution, the FEM solution and the PCFEM solution at Refinement 5.



Figure 8. The graph of $\ln \left\| u - u_{h_j} \right\|_{Y_i}$ against $-\ln h_j$, $h_j = 2^{-j}$, $j = 1, \dots, 5$.

error estimates of the FEM and the PCFEM. The convergence rate of the classical FEM is not optimal while the one of the PCFEM is optimal.

5. Conclusions

This paper presents an adaptive predictor-corrector finite element method (PCFEM) for the one-dimensional heat equation with zero initial and boundary conditions. The method is based on a Fourier decomposition with a correction of the solution by a predictor-corrector algorithm. The PCFEM is shown to be efficient on quasi-uniform meshes and then can be used without a particular mesh treatment or a special interpolation error in the energy norm. In the predictor-corrector algorithm, the stiffness matrices are the same, only the load vector is modified to correct the solution, and then the complexity of the PCFEM algorithm and the classical FEM algorithm are of the same order.

The PCFEM developed can be applied to the heat equation with non-zero initial and boundary conditions, and the proof of the stability conditions and the Homogenization methods can be adapted using the proofs developed in [7] provided that the initial and boundary conditions are regular enough. It is worth saying that the method developed is another motivation to continue using standard FEM to solve partial differential equations.

Conflicts of Interest

The author declares no conflicts of interest regarding the publication of this paper.

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List of Notations

The following are some variables used in the paper.

Variable	Meaning
п	A natural number that is either 1 or 2
m	$m \in \mathbb{N}$ is a natural number
N	$N \in \mathbb{N}$ is a natural number
$k_{_W}$	$k_w \in \mathbb{R}$ is a constant
X	$x \in \mathbb{R}$ is a single variable
x	$\boldsymbol{x} = (x_1, \cdots, x_n) \in \mathbb{R}^n, n \ge 2$
t	$t \ge 0$ is a non-negative real variable
T	T > 0 is a positive real number
L	L > 0 is a positive real number
а	$a \in \mathbb{R}$ is a real number
Ь	$b \in \mathbb{R}$ is a real number
k	$k \in \mathbb{N}$ is a natural number
$a(\cdot, \cdot)$	A bilinear form
β	$\beta = (\beta_1, \cdots, \beta_n) \in \mathbb{N}^N$ is a multi-index
$oldsymbol{eta}_k$	$eta_{_k}=rac{k\pi}{L}k_{_w}$, $k\in\mathbb{N}$, $k\geq 1$
α	$lpha \in \mathbb{R}$ is a real number
$t_i, i = 0, 1, 2$	$t_i \ge 0$ is a real non-negative number
Ω	$\Omega \subset \mathbb{R}^n$ is a domain
$\partial \Omega$	the boundary of the domain $\ \Omega \subset \mathbb{R}^n$
V	A Hilbert space V
X	A vector space X
и	A scalar function $u \equiv u(x,t)$
V	A scalar function $v \equiv v(x,t)$
arphi	A scalar function $\varphi \equiv \varphi(x,t)$
\mathcal{N}	A linear operator \mathcal{N}
С	A letter used to design a space of continuous functions
$\mathcal{T}_{_{h}}$	A triangulation with diameter <i>h</i> >
X_{h}	A vector space of piecewise linear polynomials
$X_{_{1h}}$	A vector space of functions of two variables
u_h	A function in X_h

Continued	
С	C > 0 is an arbitrary constant
W_h	A function in X_h
Π_h	The projection of X to X_h
${\mathcal Y}_k$	A function $\gamma_k \equiv \gamma_k(t)$, $k \in \mathbb{N}$, $k \ge 1$
S_k	A function $s_k(x)\sin(k\pi x/L)$
V	The convergence rate
V_{exp}	The expected convergence rate