

# Transmission Characteristics of Tuneable Optical Filters Using Optical Ring Resonator with PCF Resonance Loop

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## Abstract

A theoretical analysis of a tuneable optical filter is presented by proposing an optical ring resonator (ORR) using photonic crystal fiber (PCF) as the resonance loop. The influences of the characteristic parameters of the PCF on the filter response have been analyzed under steady-state condition of the ORR. It is shown that the tuneability of the filter is mainly achieved by changing the modulation frequency of the light signal applied to the resonator. The analyses have shown that the sharpness and the depth of the filter response are controlled by parameters such as amplitude modulation index of applied field, the coupling coefficient of the ORR, and hole-spacing and air-filling ratio of the PCF, respectively. When transmission coefficient of the loop approaches the coupling coefficient, the filter response enhances sharply with PCF parameters. The depth and the full-width half-maximum (FWHM) of the response strongly depends on the number of field circulations in the resonator loop. With the proposed tuneability scheme for optical filter, we achieved an FWHM of  $\sim 1.55$  nm. The obtained results may be utilized in designing optical add/drop filters used in WDM communication systems.

**Keywords:** Optical Ring Resonator, Photonic Crystal Fiber, Tuneable Optical Filter, Optical Fiber

## 1. Introduction

In last two decades, optical ring resonators (ORR) with different configurations, based on fiber and wave-guides [1-4], have been analyzed for different applications such as polarization sensing [5], bio-sensing [6], optical filters [7-9], dispersion compensation [10,11], optical triggering and optical integration/differentiation [2], optical bistability [12,13], add/drop multiplexer [14], optical switching [15-18], and various other applications [19-21]. In the early works, steady state [2,22] and dynamic responses of ORR built on fiber were analyzed [23,24] for applications in polarization sensing [7], FM deviation measurement of a laser diode (LD) [24], optical triggering, optical integration/differentiation and fiber dispersion compensation [2], and rotation sensing [22]. Recently, dynamic resonance characteristic of fiber ring resonator has been analyzed for gyro systems [25].

A basic structure of an ORR consists of a  $2 \times 2$  port directional coupler and a fiber or waveguide loop connecting one of the input ports to one of the output ports, making a ring resonator with a function similar to a

Fabry-Perot interferometer. To achieve the resonance effect in an ORR, the loop length could be of the order of few micrometers [26] to tens of meters [23]. Generally, the characteristics of ORR based optical filters are determined by their frequency responses which in turn depends on the characteristic parameters of the ORR. The characteristic parameters of an ORR, that influence filter response, are resonator loop length, coupling coefficient of the coupler, transmission parameters of the loop fiber, and modulation frequency of the circulating field intensity in the resonator [23,24].

In accordance with developed theory of the ORR [27], the objective of this study is to present an analysis of tuneability of optical filter using an optical ring resonator when excited by a sinewave-modulated laser diode. Based on previous ORR structures [23,24], where singlemode fiber (SMF) used as the resonator loop, in this paper, we assume to use a PCF resonance loop built on a PCF coupler [28,29]. The peculiar properties of PCFs have interesting effects on transmission function of the ORR. Two influential parameters of the PCF on transmission are airhole diameter ( $d$ ) and the air-hole spacing ( $\Lambda$ ) [30-32].

To quantify the influence of the characteristic parameters of the loop PCF on the characteristic response of the ORR, the ORR loop transmission is analyzed by using the formulation developed by Pandian *et al.* [24] and Seraji *et al.* [23]. The novelty of the PCF-based ORR compared with SMF-based ORR is that the tuneability of response of the former is opened to more characteristic parameters and can be more compact due to shorter length of the loop.

## 2. Analysis of ORR with PCF Resonance Loop

A basic structure of an ORR built in PCF loop with transmission coefficient  $\alpha$ , coupling coefficient  $\kappa$ , coupler insertion loss  $\gamma_0$ , and loop delay time  $\tau$  is shown schematically in **Figure 1** with port (1) and (2) as input  $E_{in}(t)$  and output  $E_{out}(t)$ , respectively.

To analyze the field in the loop, the loss is assumed only to be due to loop bending, and other loss mechanisms such as splice (if any), confinement, and intrinsic losses of PCF are neglected for simplicity of analysis [30,32]. The loop transmission coefficient  $\alpha$  is defined as  $\alpha = \exp(-2\alpha_0 L)$ , where  $L$  is the loop length of the resonator (in m), and  $\alpha_0$  is the bending loss coefficient in dB/km due to bending radius  $R_{bent}$ , that is given by [33]:

$$\alpha_0 \cong \frac{1.57}{A_{eff} \beta \Lambda} \left[ \frac{1}{\sqrt{x}} \exp(-x) \right], \quad x = \frac{2}{3} \frac{R_{bent}}{\beta} \left( \frac{V_{PCF}}{\Lambda} \right)^3 \quad (1)$$

where  $A_{eff} = \pi \rho_{eff}^2$

is the effective core area and

$$V_{PCF} = k \rho_{eff} (n_{co}^2 - n_{eff}^2)^{1/2}$$

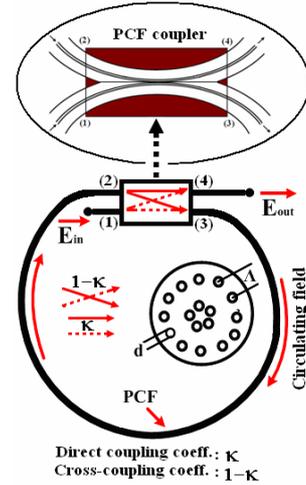
is the V-parameter of the PCF,  $\Lambda$  denotes the air-hole spacing,  $\beta = n_{eff} k$  is the propagation constant,  $k = 2\pi/\lambda$  is the wave number, and  $\lambda$  represents the wavelength in vacuum (For Neper to dB-scale conversion, the above expression should be multiplied by 8.686). The values of  $n_{eff}$ ,  $A_{eff}$  and  $V_{PCF}$  are determined by improved vectorial effective index method for different values of  $\Lambda$  and  $d/\Lambda$ , and the results are tabulated in **Table 1** [34].

The transmission coefficient  $\alpha$  of the ORR with a loop length of about 19 mm ( $R_{loop} = 3$  mm) made with four PCF structures of similar  $d/\Lambda$  and different  $\Lambda$  ( $=2.3, 4.0, 6.0,$  and  $8.0 \mu\text{m}$ ), is shown in **Figure 2**. The corresponding

transmission coefficients at  $1.55 \mu\text{m}$  wavelength are determined as  $\alpha = 0.96, 0.36, 0.04, 0.01$ , respectively. As  $\Lambda$  increases, the value of  $\alpha$  decreases. By doubling the  $\Lambda$  value,  $\alpha$  becomes 36 times higher.

The transmission coefficient is also affected by the air-filling ratio  $d/\Lambda$  with a constant  $\Lambda$ . By increasing the ratio  $d/\Lambda$ , the  $\alpha$  values will increase. In **Figure 3**, this case is depicted for  $d/\Lambda = 0.2, 0.3,$  and  $0.4$  at  $\Lambda = 4 \mu\text{m}$ .

By comparing **Figures 2** and **3**, we observe that the effect of  $\Lambda$  values on the ORR transmission is more than that of  $d/\Lambda$  variations. In general, with an input of a sinewave-modulated laser signal  $E_{in}(t)$  with an angular modulation frequency of  $\omega_m$ , the filtering characteristics of the ORR in **Figure 1**, after  $n$  number of field circulations in the resonator loop, can be expressed as (see Equation (2)) [24].



**Figure 1.** Schematic diagram of an ORR connected to a laser diode at port 1.

**Table 1.** Calculated PCF parameters for different structures.

	$\Lambda$ ( $\mu\text{m}$ )				$d/\Lambda$		
	2.3	4	6	8	0.2	0.3	0.4
$n_{eff}$	1.4395	1.4450	1.4475	1.4485	1.4550	1.4434	1.4421
$\rho_{opt}$	1.292	2.412	3.672	4.922	2.412	2.516	2.533
$V_{eff}$	0.920	1.253	1.397	1.463	1.253	1.672	2.533
$A_{eff}$	371	145	204	305	145	151	131

$$\frac{E_{out}^{(n)}}{E_{in}} = -jA\sqrt{1+k_m \sin(\omega_m t)} \exp[-j\beta \cos(\omega_m t + \phi_{FM})] + jB \sum_{n=1}^{n=\infty} \left\{ C^{(n-1)} \sqrt{1+k_m \sin[\omega_m(t-n\tau)]} \exp\{j[-\beta \cos[\omega_m(t-n\tau) + \phi_{FM}] - n\omega\tau - n\pi/2]\} \right\} \quad (2)$$

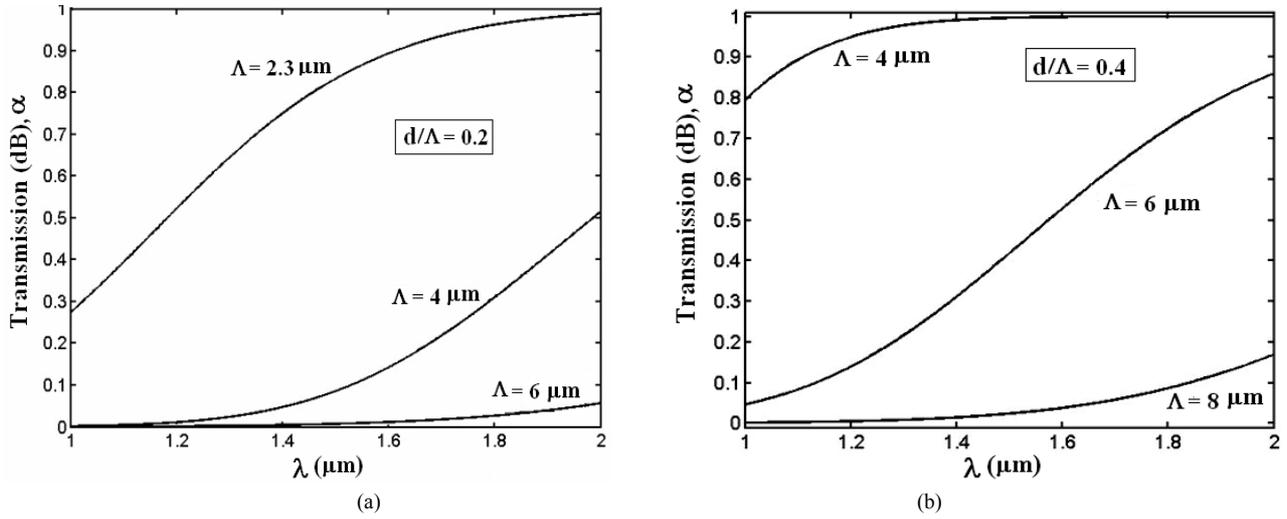


Figure 2. Transmission coefficient of the ORR in terms of wavelength for (a)  $d/\Lambda = 0.2$ , (b)  $d/\Lambda = 0.4$ , and different values of  $\Lambda$ .

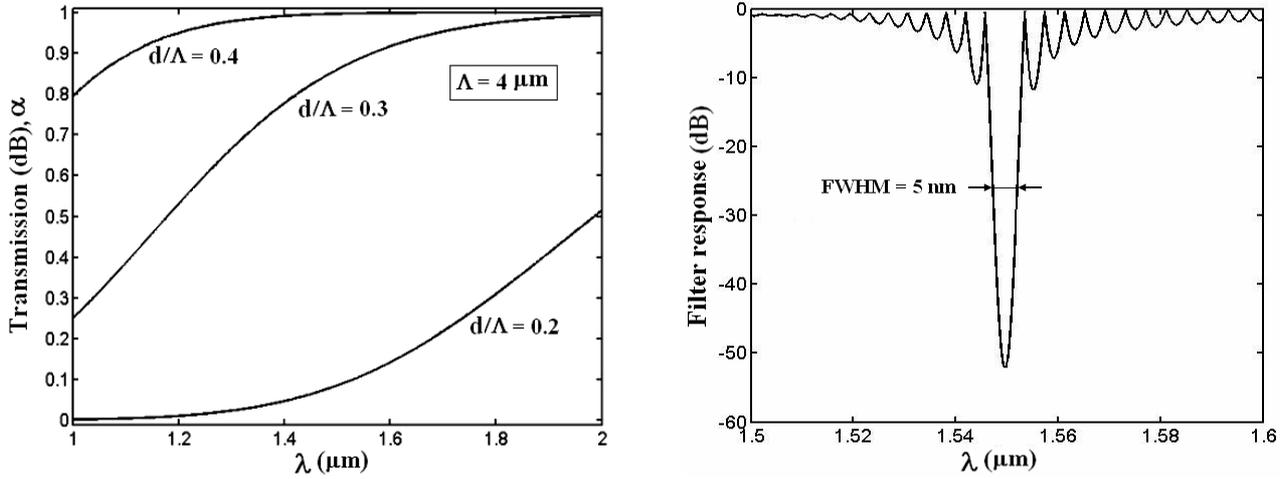


Figure 3. Transmission coefficient of the ORR in terms of wavelength for  $\Lambda = 4\mu\text{m}$  and different values of  $d/\Lambda$ .

Figure 4. Characteristic response of stop band filter based on ORR.

Where  $\beta = I_m k_f / f_m$  is frequency modulation index,  $k_f$  is the optical frequency deviation (Hz/mA),  $f_m$  is modulation frequency,  $I_m$  is the peak value of ac modulating current,  $k_m$  is the amplitude modulation index,  $\Phi_{FM}$  is angle between optical frequency deviation and amplitude modulating derive current of the given LD.

$$A = \sqrt{\kappa(1-\gamma)} \quad B = (1-\kappa)(1-\gamma)\alpha$$

and 
$$C = \sqrt{\alpha \kappa(1-\gamma)}$$

are constants and all other parameters have the same definitions as in Figure 1.

With a particular characteristic parameters given in Table 2, a stop band filter based on ORR is designed with a characteristic response plotted in Figure 4 at resonance wavelength of 1.55  $\mu\text{m}$ . The full-width at half maximum of the filter for  $t = 50\text{ris}$  is obtained as  $\sim 5\text{ nm}$ .

### 3. Effects of Laser Diode Parameters

To tune the filter based on the ORR, we can either change the parameters of the ORR such as  $\Lambda$   $d/\Lambda$  loop length,  $\kappa$  or/and parameters of input source such as  $f_m$  and  $k_m$  [24]. For instance, by increasing  $f_m$  from 95 GHz to 125 GHz, the resonance filter response shifts from 1.55  $\mu\text{m}$  to 1.615  $\mu\text{m}$ , respectively, as shown in Figure 5. The FWHM of the response slightly increases by increasing the modulation frequency of the input source.

The amplitude modulation index  $k_m$  also affects the filter response. By increasing its value, the effectiveness of filtration increases at the central wavelength, as shown in Figure 6(a). Another parameter of laser diode source influencing on the filter response is the modulation current  $I_m$  applied to the laser diode. Similar to effect of  $k_m$ ,  $I_m$

influences the depth of characteristic response, which increases by increase of  $I_m$ , as shown in **Figure 6(b)**.

In general, how the filter responds to the variation of  $I_m$  is shown in **Figure 7**. At about  $I_m = 40, 70,$  and  $90$  mA, the filter responses are at maximum values. Extreme low value occurs at about  $I_m = 65$  mA

#### 4. Effects of PCF Parameters

The transmission response of the filter strongly depends on the hole-spacing  $\Lambda$  of the PCF used in the loop of the resonator, as shown in **Figure 8**. For a given  $\Lambda$ , when  $d/\Lambda$  goes higher, the filter response goes lower and for a given  $d/\Lambda$ , higher the value of  $\Lambda$ , lower will be the filter response.

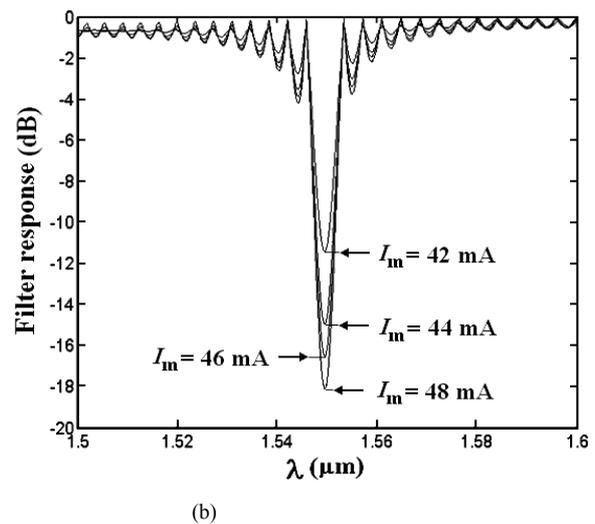
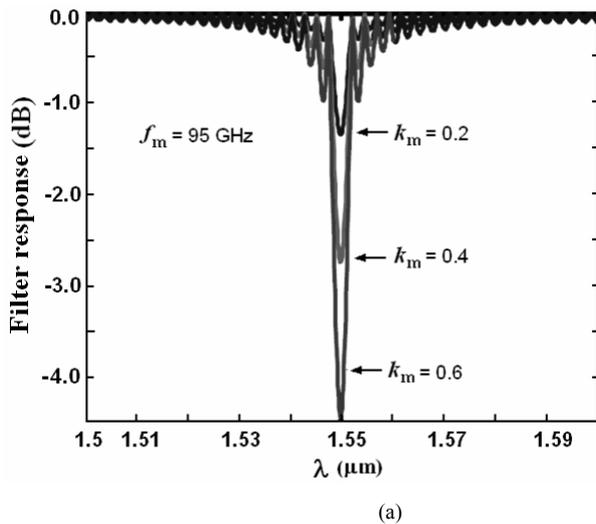


Figure 6. Effect of (a) amplitude modulation index  $k_m$  and (b) modulating current  $I_m$  on the filter response.

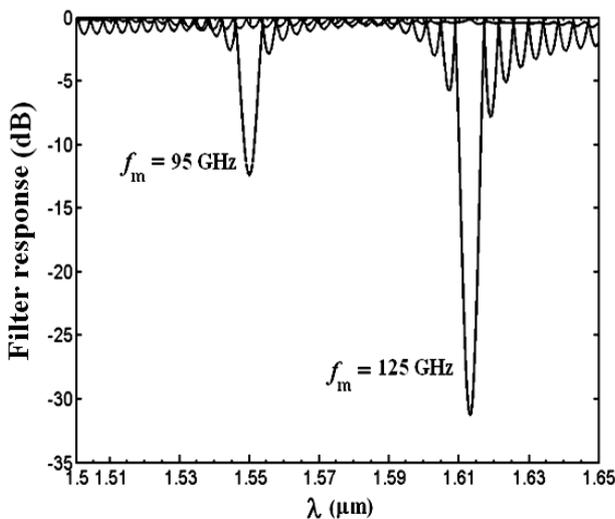


Figure 5. Effect of modulating frequency  $f_m$  on the filter response.

At a particular values of  $\Lambda$  and  $d/\Lambda$ , when the corresponding  $\alpha$  value tends to the value of  $\kappa$  the filter response reaches its maximum value.

#### 5. Effects of Characteristic Parameters of Coupler and Ring

The coupling coefficient of the resonator  $\kappa$  has direct effects on the filter response. Figure 4 is reproduced here by changing the value of  $d/\Lambda$  from 0.4 to 0.2, keeping  $\Lambda = 0.4$   $\mu\text{m}$ . When  $d/\Lambda$  decreases, the value of  $\kappa$  decreases for better filter response, as illustrated in **Figure 9**. Therefore, for an optimum filter response, the values of  $k_m$  and  $\kappa$  should be as large as possible, whereas  $\Lambda$  should be as small as possible. For the best condition,  $\Lambda$  should

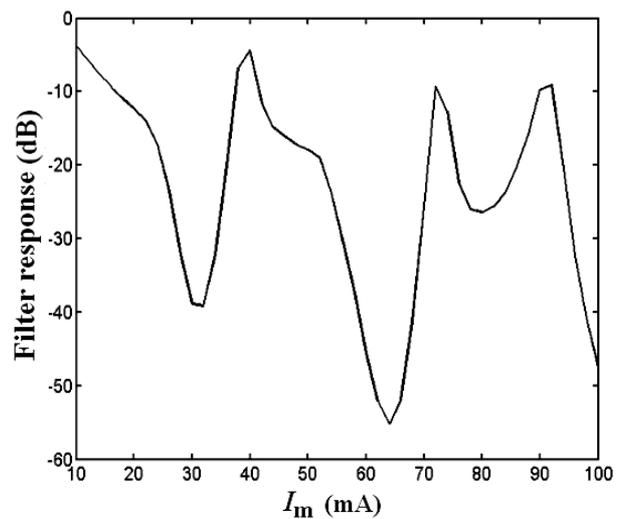


Figure 7. Variation of filter response versus modulating current.

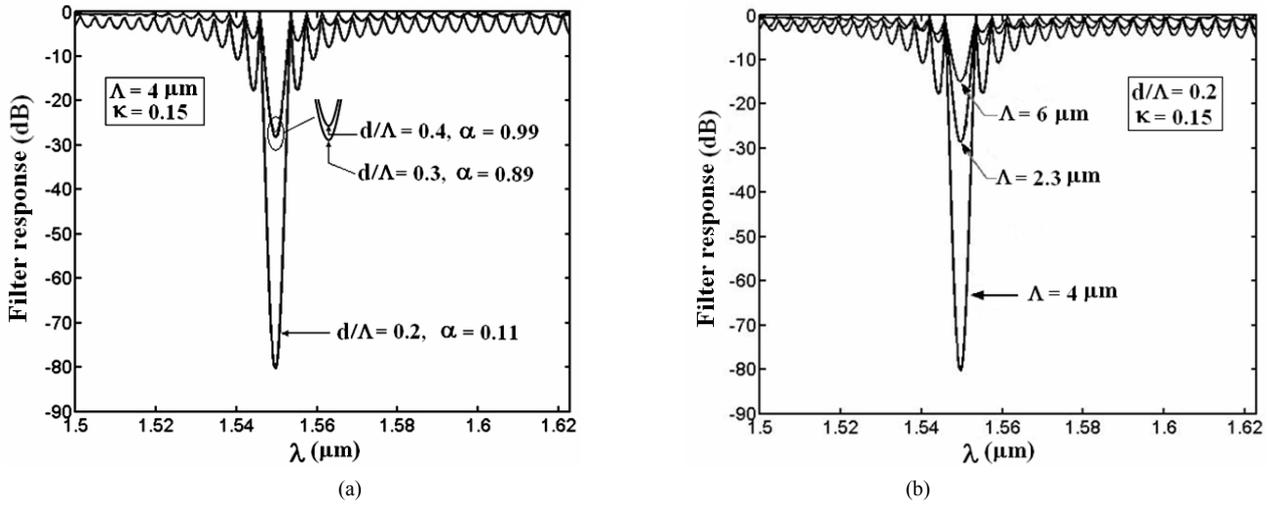


Figure 8. Effects of (a)  $d/\Lambda$  and (b)  $\Lambda$  on the filter response.

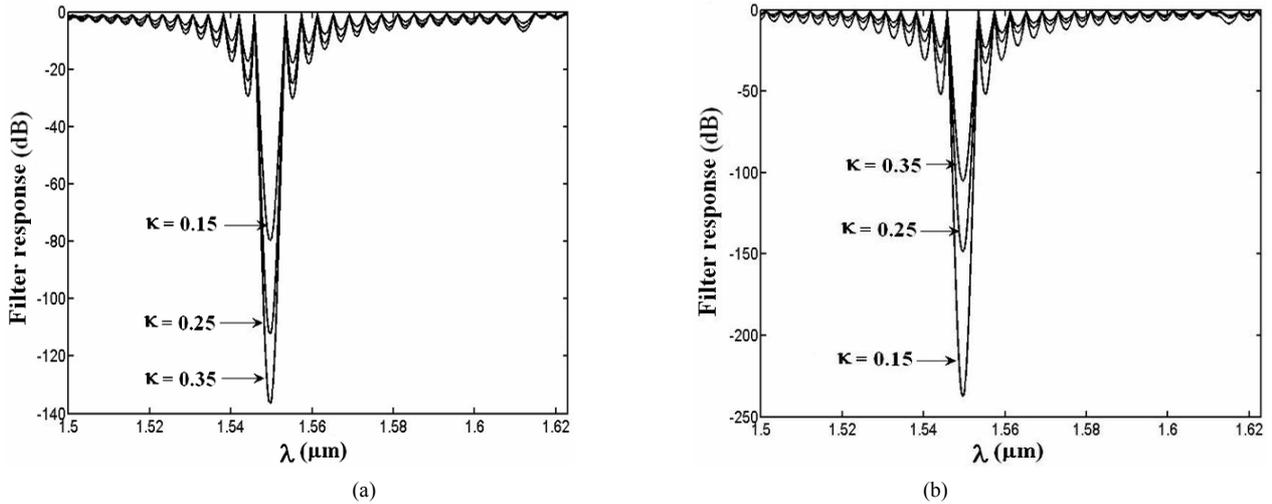


Figure 9. Effect of coupling coefficient  $\kappa$  on the filter response (a)  $d/\Lambda = 0.4$ ,  $\alpha = 0.99$  and (b)  $d/\Lambda = 0.2$  and  $\alpha = 0.11$ .

be chosen in such a way to operate the resonator at resonance state, *i.e.*, when  $\alpha = \kappa$  [35].

With the elapse of time, trapping of the field in the resonator loop stabilizes the filter response for narrower FWHM. In **Figure 10**, the filter response at  $\lambda = 1.55 \mu\text{m}$  is illustrated with the same PCF and source parameters values of **Figure 4**. The FWHM becomes narrower exponentially at central wavelength  $1.55 \mu\text{m}$ .

Typically, its value becomes 5 nm at circulation time of  $t = 50\tau$  ps and reaches to 1.55 nm at  $t = 150 \tau$  ps. That is, by tripling the loop delay time of the resonator, we obtain 70% reduction in value of the FWHM.

### 6. Filter Tuning Based on Modulating Frequency

With reference to **Figure 5**, the response wavelength of

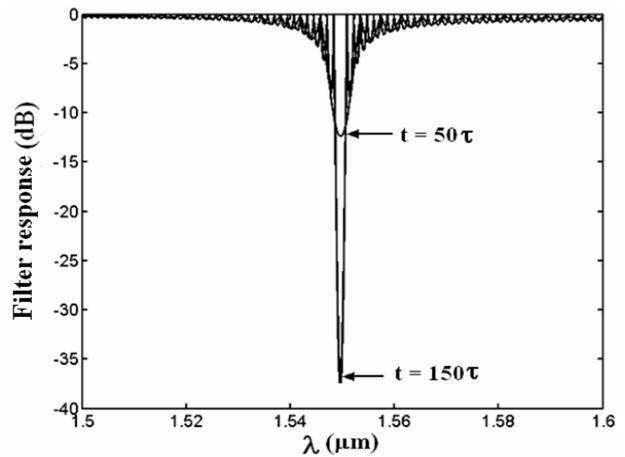
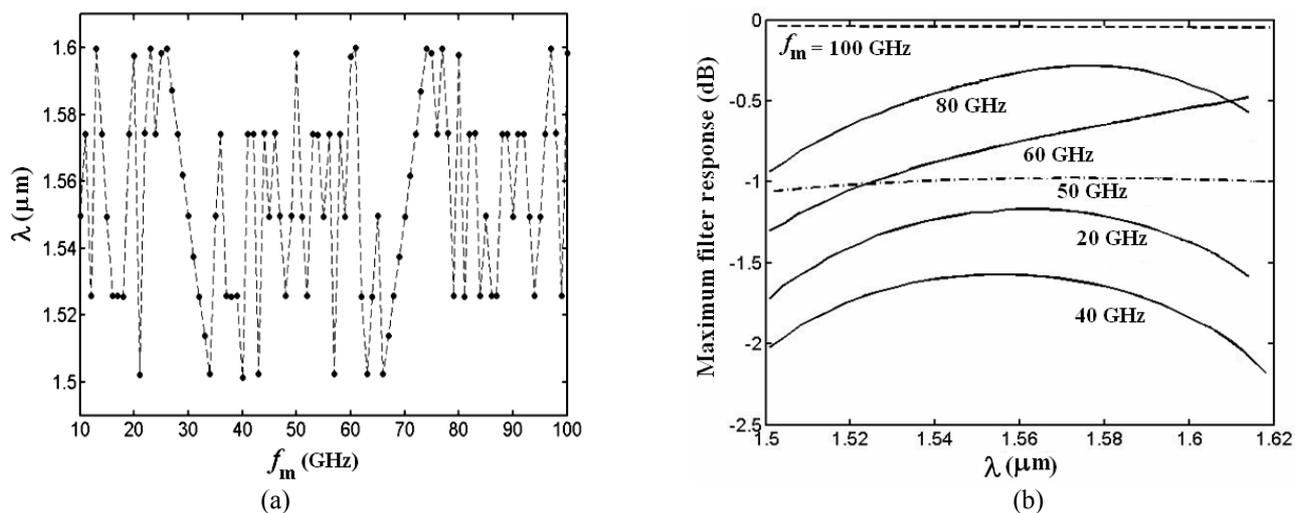


Figure 10. Narrowing of filter response in the resonator loop by increasing the loop delay time.



**Figure 11. (a) Maximum filter response obtained by the applied modulating frequencies and (b) maximum response versus wavelength for different modulating frequency.**

the filter may be shifted up or down by varying modulating frequency  $f_m$  applied to the laser diode source. **Figure 11(a)** illustrates the operating wavelengths in terms of modulating frequency where the values denoted by solid circles represent the maximum filter response obtained by the applied modulating frequencies. At some modulating frequencies such as 20, 40, 60, 80 GHz, there are more resonance peaks in the response. To realize the tuneability of the filter rendered by the modulating frequency, **Figure 11(b)** is depicted for the maximum response as a function of the operating wavelengths. For  $f_m$  with values as multiple of 50 GHz, the maximum response is almost independent of the wavelength. By using these curves, one can obtain maximum filter response at different modulating frequency.

## 7. Conclusions

In conclusion, we presented a tuneable optical filter based on an optical ring resonator with a loop made of photonic crystal fiber. It is shown that the filter response can be varied by modulation frequency of the input signal. The transmission amplitude of the filter can be optimized by parameters such as modulation signal amplitude, hole-spacing of the PCF, and coupling coefficient of the ORR. The filter response enhances sharply with PCF parameters, when transmission coefficient of the loop approaches the coupling coefficient. It is further illustrated that by signal field circulations in the ORR loop, the filter response stabilizes to a narrower FWHM.

With the proposed tuneability scheme for optical filter, we presented maximum filter response with respect to operating wavelengths and achieved an FWHM of  $\sim 1.55$

nm. The obtained results may be utilized in designing optical add/drop filters used in WDM communication systems.

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