

# Nonequilibrium Diffusion of Boron in SiC at Low Temperatures

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## ABSTRACT

*Nonequilibrium diffusion of Boron in 3C SiC was performed using a flow of carbon vacancies. The temperature of diffusion was 1150-1250°C and concentration of Boron in doped area reached about  $10^{19}$  to  $10^{20}$  cm<sup>-3</sup>. It is shown that after thermal annealing in vacuum the characteristics of fabricated structures are close to those of the structures made by the conventional technology.*

**Keywords:** Diffusion, Activation Energy, Silicone Carbide, Annealing, Vacancy

## 1. Introduction

Silicon carbide is widely accepted as an ideal semiconductor material for advanced high power and high frequency electronic devices with potential applications in spacecraft, aircraft, ships, land vehicles, communication systems, as well as power generating, conversion, and transmission systems. These radiation-hard, high performance devices can operate in harsh environment and at elevated temperature range beyond the capabilities of conventional silicon devices. A key process step in the fabrication of silicon carbide devices is controlled impurity diffusion for p-type-conversion of selected n-type areas in the active region of a device. The conventional diffusion method for silicon carbide requires a high temperature between 1600-2300°C. This high temperature process is complicated, expensive, and can generate more undesirable structural defects.

It is well known that mobile point defects play an important role in diffusion processes in semiconductors. Experiments on doped SiC samples show that diffusion is mediated by interstitials and vacancies [1,2]. To quantitatively assess the role of interstitials and vacancies, *ab initio* investigations of the relevant processes in 3C-SiC were performed, which showed that carbon vacancies and silicon interstitials in both p-type and intrinsic materials were the most abundant mobile defects. Thus the interaction between point defects and the diffusing impurity atoms can be used to lower the temperature at which the dopant can be introduced. In this paper, we report on the

diffusion parameters of boron diffusion in 3C-SiC in a flow of carbon vacancies at low temperature regime interval, 1100-1250°C, and present some electrophysical results.

## 2. Experimental Results and Discussion

The diffusion of boron in 3C-SiC in flow of Carbon vacancies was carried out at different diffusion temperatures between 1100 and 1250°C using the method described in [3,4]. In this low-temperature diffusion technique, a thin film of borosilicate glass is created on the surface of a silicon carbide sample. The sample is then annealed under special conditions in the air. Carbon atoms from the SiC surface interact with oxygen and leave the SiC surface. As a result a flow of carbon vacancies was created in the near-surface region. Thus diffusion of Boron was performed in nonequilibrium conditions in the flow of carbon vacancies from the surface into volume of crystal.

Initial experiments were conducted on n-type 3C-SiC/Si structures grown by chemical vapor deposition. The thickness of the SiC epitaxial layer was about 5-6 μm, and the free electron concentration in the n-type 3C-SiC/Si layers about  $10^{17}$ - $10^{18}$  cm<sup>-3</sup>. Under the conditions of this experiment, an anomalously high diffusion coefficient and solubility of B was observed. After diffusion at 1150-1250°C for 15-30 minutes, all the 3C-SiC/Si epitaxial films of the thickness 3-5 μm became p-type. Calculations showed that in this case the diffusion factor of

boron,  $D$ , was about  $10^{-10}$  cm<sup>2</sup>/sec, and the solubility- $N_B = 10^{18}$ - $10^{20}$  cm<sup>-3</sup>. It should be noted that in the case of ordinary diffusion, the magnitude of  $D_B \sim 10^{-10}$  cm<sup>2</sup>/sec can only be observed at a very high temperature  $\sim 2100$ - $2300$  °C.

We suggested two possible reasons for this phenomenon of an increase of the diffusion rate. The first possible reason is the interaction of B atoms with a flow of vacancies from the surface to the bulk. The second possible reason is the interaction of B atoms with strain-related growth defects in the 3C-SiC/Si films. It is well known that a large concentration of strain-related growth defects can occur at the 3C-SiC/Si interfaces due to the large lattice mismatch between 3C-SiC and Si. The large defect concentration at the 3C-SiC/Si interface can influence the diffusion processes near the surface in the case of thin films.

To reveal the role of growth-related defects, we conducted also diffusion studies in homoepitaxial 6H-SiC/6H-SiC structures, which had no growth defects from lattice mismatch. In these structures, type conversion from n-type to p-type was observed also at diffusion temperatures  $\geq 1150$  °C.

The distribution of B in the SiC/SiC samples was measured by spreading current method. In this method, the samples were polished by diamond paste at a gradient of 1-3 degrees to the surface of the SiC film. Then the distribution of specific resistance along the surface of the SiC film was measured to obtain the distribution profile of boron,  $N_B(x)$ .

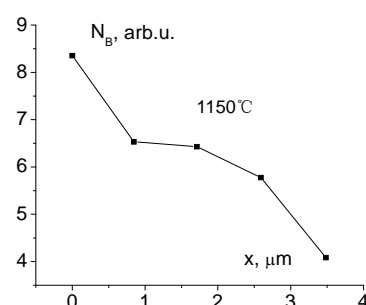
On the basis of our experimental data, shown in **Figures 1** and **2**, the following parameters of B diffusion in the presence of vacancy flow were evaluated:  $D_B$  about  $5.5 \times 10^{-11}$ - $5.0 \times 10^{-10}$  cm<sup>2</sup>/s in the temperature interval of 1150-1250 °C; activation energy of B diffusion,  $\epsilon_a \approx 0.9$ - $1.15$  eV. This activation energy of diffusion is typical for interstitial mechanism of diffusion in ordinary semiconductors.

Secondary ion mass spectroscopy (SIMS) is a direct method to confirm the penetration of Boron atoms into SiC in our diffusion experiments. SIMS measurements showed (see **Figure 3**) that both atomic  $B_1$  and molecular boron  $B_2$  were detected on the surface of SiC and also deep within the bulk [5], as well as the peaks of boron trimer,  $B_3$ . In **Table 1**, the depth distribution of boron and other elements and clusters measured by SIMS is presented. It is shown that atomic boron is present in the whole diffusion area. Relative concentration of carbon and silicon atoms is constant for all depths, concentration of oxygen atoms lowered and that of lattice cluster ions,  $C_2^-$ ,  $Si_2^-$ ,  $SiC^-$ , and  $SiC_2^-$ , increased.

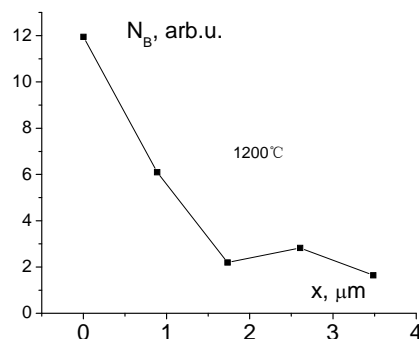
It is necessary to note that boron clusters are registered on an initial surface of SiC<B>, and  $B_2$  and  $B_3$  are not registered at the depth of about 0.024 mkm. Boron dimers

**Table 1. Depth distribution of Boron and other elements and clusters measured by SIMS**

Depth, $\mu\text{m}$	0,024	0,12	0,24	0,36
B	17	20	20	20
C	40	40	45	45
O	120	100	100	80
OH	25	30	25	20
$B_2$	0	5	5	10
$C_2$	35	30	25	35
Si	40	40	45	45
$O_2$	15	10	5	0
SiC	25	30	30	35
$SiC_2$	15	15	20	20
$Si_2$	20	25	25	25
$B_3$	0	0	5	10



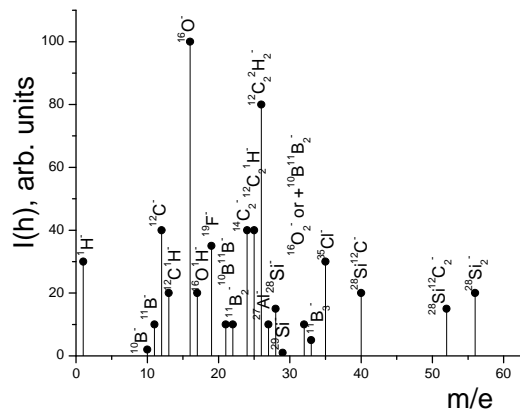
**Figure 1. Distribution of  $N_B$  versus  $x$  after diffusion in non-equilibrium conditions at 1150 °C.**



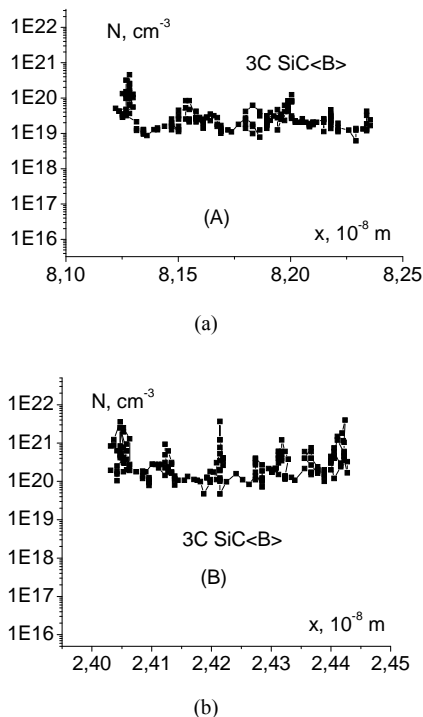
**Figure 2. Distribution of  $N_B$  versus  $x$  after diffusion in non-equilibrium conditions at 1250 °C.**

$B_2$  appear at the depth of 0.12 mkm, and boron trimmers  $B_3$  at 0.24 mkm.  $B_2$  and  $B_3$  intensity increased up to 10 at the depth of 0.36 mkm. To perform semiquantitative analysis of the relation between atomic and cluster boron, we need to take into account an electron affinity of B,  $B_2$ ,  $B_3$ .

The interesting feature of the spectra is the presence of peaks of boron trimer  $B_3$ . The  $B_2$  boron dimers in silicon are known from the literature, whereas the data for  $B_3$  are



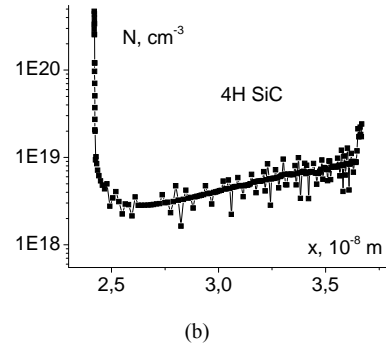
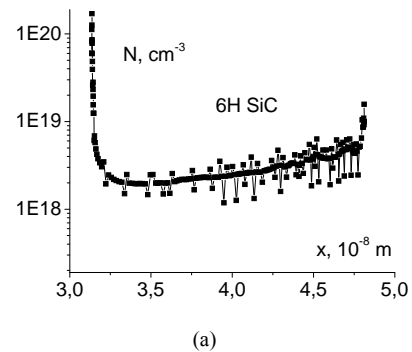
**Figure 3.** Mass spectra of negative ions from SiC<B> surface spattered by Cs<sup>+</sup> (E=2 keV at 300 C, time of spattering 1 hour).



**Figure 4.** Distribution of Boron in the near surface area obtained from CV data (3C SiC samples).

absent. In our opinion, this is due to nonequilibrium conditions of diffusion and modified surfaces of SiC. The clustering of boron obtained from SIMS data shows that the boron diffusion in SiC can be provided by interstitial boron diffusion mechanism with suppression of formation of BX complex of boron with lattice atoms C and Si.

To define the concentration of Boron in the near surface area, the Schottky barriers on doped layer of samples were fabricated and their characteristics measured. **Figures 4 and 5** shows the distribution of Boron obtained from CV data of Ni-p-3C SiC<B>, p-6H SiC<B>, and Ni-p-4H



**Figure 5.** Distribution of Boron in the near surface area obtained from CV data (6H and 4H SiC samples)

SiC<B> structures. As seen from figures, under diffusion of B in a flow of carbon vacancies, an anomalously high solubility of B in SiC is observed. The distribution of doped impurity near the surface for samples on the basis of 3C polytype of silicon carbide has a “quasiperiodical” character. We believe that “quasiperiodical” distribution of Boron is related to synergy processes, which can be observed due to interaction between impurity and point defects of crystal at high concentration [6].

## 2.1 Mechanism of Diffusion

The mechanism of a conventional high-temperature diffusion of a Boron in a silicon carbide is rather complex and there exist various points of view about this mechanism.

In the work [7] (published in 2001), the diffusion experiments with B-implanted 6H-SiC at temperatures between 1700°C and 2100°C were performed. After diffusion, B concentration profiles were measured by means of secondary ion mass spectrometry. Accurate modeling of B diffusion was achieved on the basis of a kick-out diffusion mechanism, which involves only point defects in the Si sublattice.

In the work [8] (published in 2003), the authors, on the basis of diffusion experiments at 2000-2200°C in 6H SiC, suggest that the boron impurity possesses two different solubilities depending on occupancy at either Si or C lattice sites. The solubility ( $6 - 9 \times 10^{19} \text{ cm}^{-3}$ ) in a Si va-

cancy is about one order of magnitude higher than that ( $1 - 10 \times 10^{18} \text{ cm}^{-3}$ ) in a C vacancy. According to [9], the concentration of carbon vacancies in SiC at high temperatures can be as high as  $3 - 5 \times 10^{17} \text{ cm}^{-3}$ , and, according to authors of [8], the carbon vacancies will contribute to the diffusion of boron in SiC. It can be speculated that the boron atoms migrate into SiC bulk by occupying substitutional carbon sites ( $B_C$ ) or forming associates of boron and carbon vacancies ( $B-V_C$ ). The diffusion of boron via carbon vacancies has a much higher diffusion rate due to the existence of high-density carbon vacancies in SiC, which lead to a boron profile with a fairly long tail.

In the work [10] (published in 2009), for a modeling of diffusion of Boron in SiC, it is also supposed, that the diffusion occurs both on carbon and silicon sublattices. Reported profiles of high-temperature ( $500^\circ\text{C}$ )-implanted boron ions diffused in 4H-silicon carbide at  $1200-1900^\circ\text{C}$  for 5-90 min were simulated through a “dual-sublattice” modeling, in which a different diffusivity is assigned for diffusion via each sublattice.

These approaches are related to conclusions of the work [11], where the authors claim that most mobile states of Boron are ( $B_C-V_C$ ) associates. These mobile associates dissociate substantially during diffusion-annealing process with formation of more stable ( $B_{Si}-V_C$ ) associates.

On the basis of abovementioned, for our experiments with a flow of carbon vacancies, we can speculate that the nonequilibrium concentration of  $V_C$  in the near surface area of sample is very high. Most likely, the diffusion in this area proceeds via mobile ( $B_C-V_C$ ) associates.

As to the low energy of activation about 0.9-1.15 eV, determined in our experiment, we can mention the following. In the work [12], *ab initio* theory of point defects and defect complexes in SiC was developed. It was shown that for interstitial-mediated boron diffusion in 3C-SiC, the migration and kick-out/kick-in barriers depended on the charge state and started from 0.7 eV and higher.

Certainly, the further researches are necessary for the specification of the mechanism of diffusion and large solubility of Boron in our conditions.

## 2.2 Applicability of This Technology

Since the low temperature diffusion is performed in a flow of carbon vacancies, the obtained p-SiC layers contain defects and their characteristics are worse in comparison with the structures fabricated by the conventional technology. Applicability of this technology is subject to availability of the efficient ways of removing of defects in the structures. Annealing experiments showed that annealing of above mentioned Schottky barriers at  $500$ ,  $700$ , and  $900^\circ\text{C}$  led to an improvement of IV and CV characteristics.

To study in more details the possibility of removing of defects in SiC<B> samples, the influence of thermal annealing on the electroluminescence spectra was investigated [13].

Subsequent thermal annealing of the p-n structures was made at temperatures between  $550$  and  $600^\circ\text{C}$  for 5 minutes. Electroluminescence spectra were measured at room temperature and reverse biases ( $15-20 \text{ V}$ ) were applied for p-n samples before and after annealing. Density of reverse current was about  $2-3 \text{ A/cm}^2$ .

Figures 6 and 7 shows the electroluminescence spectra of p-n junctions made by a low temperature diffusion of boron in 3C-SiC epilayer with various thicknesses (1 and 12 microns). As shown in figures, the electroluminescence spectra correlate strongly with the amount of defects in the epitaxial films.

The deep Boron level in SiC is related to ( $B_{Si}-V_C$ ) center [11], and, in our case of diffusion in a flow of carbon vacancies, we have to observe luminescence peaks related with  $B_{\text{deep}}$ : the peak at about  $810 \text{ nm}$  (nitrogen-deep- boron transitions) and the peak at about  $760 \text{ nm}$  ( $E_C-B_{\text{deep}}$  transitions) (see Figures 6-8). There is also the luminescence peak at  $630 \text{ nm}$  connected with carbon-silicon divacancy [14]. It is clear that all these defects were formed during the low temperature diffusion process.

As is well known, the amount of growth-related defects in SiC/Si structures decreases with an increase of the thickness of the epitaxial layer. The thickness of this strained transition layer is about  $1.5-3 \text{ microns}$ . Due to this,

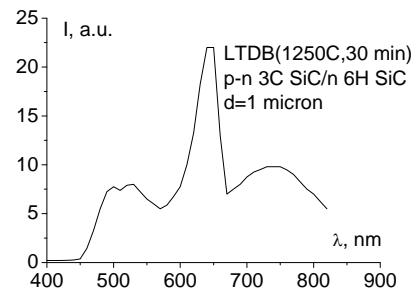


Figure 6. Spectrum of an electroluminescence of p-n junction made by low temperature diffusion of Boron in 3C SiC film (thickness-1 micron)

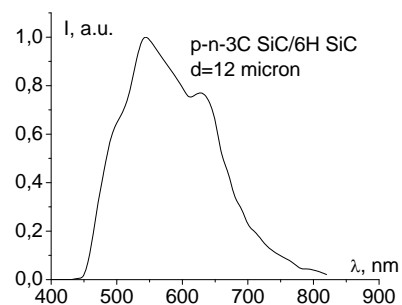


Figure 7. Spectrum of an electroluminescence of p-n junction made by low temperature diffusion of Boron in 3C SiC film (thickness-12 micron)

the peaks at 630 and 760 nm in a thin (1  $\mu\text{m}$ ) epilayer look more distinctly than in a thick epilayer (12  $\mu\text{m}$ ).

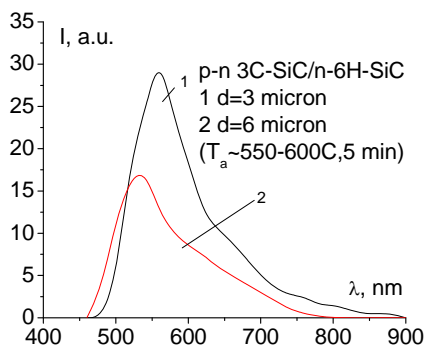
The optical absorption spectra of the p-n structures were also measured at room temperature. The optical absorption data were in good agreement with the electroluminescence data. To study the effect of annealing on the defect structure, we annealed the p-n junctions at the temperature about 550-600°C during 5 minutes. The resulting luminescence spectrum of an annealed sample is given in **Figure 8**. A reduction in the intensity of the luminescence peaks connected with deep Boron related defects (760 and 810 nm), and carbon-silicon divacancies (630 nm), entered during low temperature diffusion, was observed. For a comparison, an electroluminescence spectrum of the epitaxial p-n junction fabricated on 3C SiC by the conventional chemical vapor deposition (CVD) technology is shown in **Figure 9**. In this case, the crystalline quality of the sample is better with a lower concentration of growth defects. A luminescence peak is observed at about 530 nm corresponding to an energy gap about 2.35 eV for 3C-SiC polytype.

A comparison of **Figures 8 and 9** shows that the luminescence spectra of p-n junctions fabricated by low temperature diffusion followed by thermal annealing are close to the spectrum of epitaxial junction fabricated by the conventional CVD. At the same time, it is necessary to note that the peak at 490-500 nm is clearly observed for non-annealed samples and weakly visible on samples after annealing. We believe that this peak is due to transitions of hot electrons accelerated at applied reverse bias from minimum of conductivity band in L1 direction to the top of the valence band.

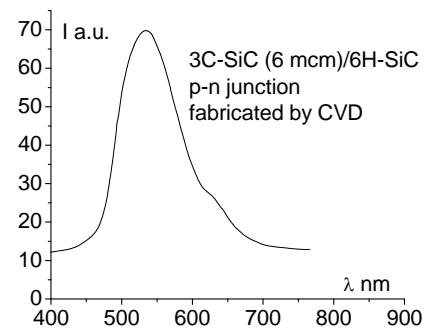
Thus it is shown that after thermal annealing in vacuum the characteristics of fabricated structures are close to those of the structures made by the conventional technology.

### 3. Conclusions

Nonequilibrium diffusion of Boron in 3C, 4H SiC was



**Figure 8.** A spectrum of electroluminescence of p-n junctions for SiC films with a thickness of 3 and 6 microns after an annealing.



**Figure 9.** A spectrum of an electroluminescence of p-n 3C SiC/6 H SiC junction fabricated by growth of p and n – SiC layers by CVD.

performed using a flow of Carbon vacancies. The temperature of diffusion was 1150-1250°C, and the concentration of Boron  $\sim 10^{19-20} \text{ cm}^{-3}$ . It is shown that thermal annealing in vacuum can be used to improve the characteristics of fabricated structures.

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