

On Unparticle Searches through Photon-Photon Scattering

Trieu Quynh Trang, Ha Huy Bang, Nguyen Thu Huong, Sa Thi Lan Anh

Faculty of Physics, Hanoi University of Sciences, Vietnam National University, Hanoi, Vietnam

Email: c25tuan@gmail.com, hahuybang@hus.edu.vn

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Abstract

In this work, we study the effects of the spin-0 unparticle on $\gamma\gamma \rightarrow \gamma\gamma$ process. From the numerical results, we show that the cross section with unparticle effect should be about $10^{27} - 10^{30}$ times larger than the one that is confirmed by QED calculation. This could have important implications for unparticle searches and for the measurement of the photon-photon cross section.

Keywords

Unparticle Physics, Photon-Photon Scattering

1. Introduction

Photon-photon scattering is among the most important and carefully studied processes in particle physics [1]-[13]. In ref. [12] Liang and Czarnecki have shown a simple way of correctly computing the low-energy $\gamma\gamma$ scattering. Noteworthy, in ref. [13], we have pointed out the cross section with radion effect should be about 10^{20} times larger than the one without radion effects. It is well known that Georgi [14] made an interesting observation that a nontrivial scale invariance sector of scale dimension d_u might manifest itself at low energy as a nonintegral number d_u of invisible massless particles, dubbed unparticle u . If unparticles exist, their phenomenological implications should be discussed. In the literature, there have been many discussions which investigate various features of unparticle physics [15]-[24]. In some of these reseaches several unparticle production processes have been studied. Possible evidence for this scale invariant sector might be the signature of a missing energy. It can be tested experimentally by examining missing energy distributions. Other evidence for unparticles can be explored by studying its virtual effects. In this letter, we consider the phenomenology of unparticle signals in $\gamma\gamma$ scattering. The unparticle analysis of $\gamma\gamma$ scattering has been done first by Cakir and Ozansoy [25] and later by Chang, Cheung and Yuan [26]. This scattering is described by the Feynman diagrams presented in **Figure 1**.

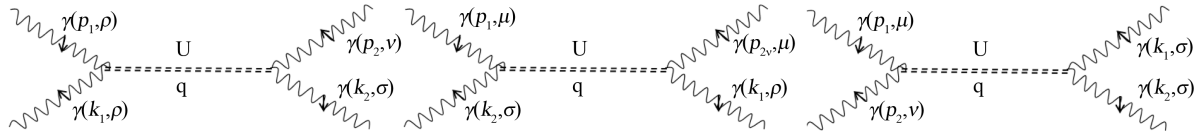


Figure 1. Feynman diagram for the gamma gamma scattering through a scalar unparticle.

The $\gamma\gamma u$ vertex is given by

$$4i \frac{\lambda_0}{\Lambda_u^{d_u}} (-p_1 p_2 g^{\mu\nu} + p_1^\nu p_2^\mu). \quad (1)$$

The spin-0 unparticle propagator is [27] [28]

$$\Delta_s(q^2) = \frac{A_{d_u}}{2 \sin(d_u \pi)} (-q^2)^{d_u-2}, \quad (2)$$

where

$$A_{d_u} = \frac{16\pi^2 \sqrt{\pi}}{(2\pi)^{2d_u}} \frac{\Gamma\left(d_u + \frac{1}{2}\right)}{\Gamma(d_u - 1) \Gamma(2d_u)}, \quad (3)$$

and

$$(-q^2)^{d_u-2} = \begin{cases} |q^2|^{d_u-2} & \text{if } q^2 \text{ is negative and real} \\ |q^2|^{d_u-2} e^{-id_u\pi} & \text{for positive } q^2 \text{ with an infinitesimal.} \end{cases}$$

The angular distribution is [26]

$$\begin{aligned} \frac{d\sigma_u}{d\cos\theta} = & \frac{\lambda_0^4}{16\pi\Lambda_u^{4d_u}} \frac{A_{d_u}^2}{\sin^2(d_u\pi)} S^{2d_u-1} \left\{ 1 + \left(\frac{1-\cos\theta}{2}\right)^{2d_u} + \left(\frac{1+\cos\theta}{2}\right)^{2d_u} \right. \\ & \left. + \cos(d_u\pi) \left[\left(\frac{1-\cos\theta}{2}\right)^{d_u} + \left(\frac{1+\cos\theta}{2}\right)^{d_u} + \left(\frac{1-\cos^2\theta}{4}\right)^{d_u} \right] \right\}. \end{aligned} \quad (4)$$

From (4), the total cross section is found to be [26]

$$\sigma_u = \frac{\lambda_0^4}{8\pi\Lambda_u^{4d_u}} \frac{A_{d_u}^2}{\sin^2(d_u\pi)} S^{2d_u-1} \left\{ 1 + \frac{2}{2d_u+1} + \frac{2\cos(d_u\pi)}{d_u+1} + \frac{\sqrt{\pi}}{2^{2d_u+1}} \frac{\Gamma(d_u+1)}{\Gamma\left(d_u+\frac{3}{2}\right)} \right\}. \quad (5)$$

We now turn to the numerical analysis of the total cross sections. The input parameters are $\lambda_0 = 1; \Lambda_u = 1.5$ TeV. The total cross section for the unparticle contributions for $d_u = 1.1, d_u = 1.2, \dots, d_u = 1.5$ at different energies is given in **Table 1**.

It is to be noticed that in refs. [10] [12] the authors have determined the differential and total cross section for the photon-photon scattering without radion and unparticle effects

$$\frac{d\sigma_0}{d\Omega} = \frac{139\alpha^4\omega^6}{(180\pi)^2 m^8} (1 + \cos^2\theta)^2, \quad (6)$$

$$\sigma_0 = \frac{973\alpha^4\omega^6}{10125\pi m^8}. \quad (7)$$

Next, in ref. [13] we have investigated the effect of the radion on photon-photon scattering. We obtained the total cross section as follows

$$\sigma_R = \frac{107}{15} \times \frac{|c|^4}{32\pi} \times \frac{S^3}{(S - m_\phi^2)^2}, \quad (8)$$

$$\text{with } c = -\frac{\alpha}{4\pi\Lambda_\phi} \left\{ a(b_2 + b_\gamma) - a_{12} \left[F_1(\tau_w) + \frac{4}{3} F_{1/2}(\tau_t) \right] \right\},$$

where $b_2 = \frac{19}{6}; b_\gamma = -\frac{41}{6}$ are the $SU(2)_L \otimes U(1)_\gamma$ β -function coefficients in the SM,

$$a_{12} = a + \frac{c}{\gamma}, \tau_w = \frac{4m_W^2}{q^2}, \tau_t = \frac{4m_t^2}{q^2}, q^2 = m_\phi^2.$$

The form factors are given by

$$F_{1/2}(\tau) = -2\tau [1 + (1 - \tau)f(\tau)],$$

$$F_1(\tau) = 2 + 3\tau + 3\tau(2 - \tau)f(\tau),$$

$$\text{with } f(\tau) = \begin{cases} \arcsin^2(1/\sqrt{\tau}), & \tau > 1, \\ -\frac{1}{4} \left[\ln \left(\frac{1 + \sqrt{1 - \tau}}{1 - \sqrt{1 - \tau}} \right) - i\pi \right]^2, & \tau < 1. \end{cases}$$

The important property of $F_{1/2}(\tau)$ is that, for $\tau > 1$, it very quickly saturates to $-4/3$, and to 0 for $\tau < 1$. $F_1(\tau)$ saturates quickly to 7 for $\tau > 1$, and to 0 for $\tau < 1$.

From this, we have found that the effects of the radion can be strong. Interestingly, we have shown that the cross section with radion effect should be about 10^{20} times larger than the one without radion and unparticle effects. Now, by the results just mentioned we give the numerical values of the ratio of the total cross section with unparticle effects σ_u of (5) to the σ_0 of (7) at different energies in **Table 2**.

So, direct computations have showed that the above cross section of (5) should be about $10^{27} - 10^{30}$ times larger than the one in (7).

Finally, we have obtained the ratio of the cross section with unparticle effects σ_u of (5) to the one with radion effects σ_R of (8) at different energies in **Table 3**. We take $\Lambda_\phi = 1.5 \text{ TeV}$; $m_\phi = 200 \text{ GeV}$ as input parameters.

Table 1. The total cross section with radion effects of the process $\gamma\gamma \rightarrow \gamma\gamma$ for $d_u = 1.1 - 1.5$ at different energies.

\sqrt{S} (GeV)	σ_u (fb)				
	$d_u = 1.1$	$d_u = 1.2$	$d_u = 1.3$	$d_u = 1.4$	$d_u = 1.5$
300	228.02	62.73	28.78	5.85	1.75
500	7.77×10^2	2.62×10^2	1.48×10^2	0.35×10^2	0.13×10^2
800	24.03×10^2	9.77×10^2	6.64×10^2	1.91×10^2	0.88×10^2
1000	4.11×10^3	1.83×10^3	1.36×10^3	0.42×10^3	2.16×10^3
3000	57.34×10^3	39.58×10^3	45.64×10^3	22.22×10^3	17.54×10^3

Table 2. The ratio of the total cross section with unparticle effects to one without radion and unparticle effects at different energies.

\sqrt{S} (GeV)	$\frac{\sigma_u}{\sigma_0} [10^{-27}]$				
	$d_u = 1.1$	$d_u = 1.2$	$d_u = 1.3$	$d_u = 1.4$	$d_u = 1.5$
300	31.35	8.63	3.96	0.76	0.24
500	106.85	36.03	20.30	4.83	1.86
800	330.44	134.43	91.38	26.22	12.20
1000	564.49	251.10	186.61	58.54	29.78
3000	7.88×10^3	5.44×10^3	6.27×10^3	3.05×10^3	2.41×10^3

Table 3. The ratio of the total cross section with unparticle effects to one with radion effects at different energies.

\sqrt{s} (GeV)	$\frac{\sigma_u}{\sigma_R} [10^{-6}]$				
	$d_u = 1.1$	$d_u = 1.2$	$d_u = 1.3$	$d_u = 1.4$	$d_u = 1.5$
300	58.31	16.04	7.36	1.43	0.45
500	163.92	55.27	31.14	7.41	2.85
800	246.21	100.16	68.08	19.53	9.09
1000	282.32	125.58	93.32	29.27	14.89
3000	471.01	325.12	374.94	182.52	144.03

It has already been shown that the total cross section σ_u is larger than the one σ_R by 6 - 8 order of magnitudes.

To conclude, in this letter we have studied the unparticle effects on gamma gamma scattering. From numerical results, we have found that the effects of the unparticle on the cross sections can be very strong. If the measurement is carried out at $\sqrt{s} = 300$ GeV - 3000 GeV, then the cross section for the photon scattering should be detectable. This could have important implications for unparticle searches at future colliders. Our work can be extended for other scatterings, for example, Bhabha scattering or $\gamma\gamma \rightarrow e^+e^-$ process.

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