

# Coupling Effects of Depletion Interactions in Three-Sphere Systems with Different Size Ratio of Large- to Small-Sphere

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## ABSTRACT

The coupling effects of depletion interactions in three-sphere systems with different size ratio of large- to small-sphere are studied by Monte Carlo simulations in this paper. The numerical results show that this coupling effect is affected by the size ratio of large- to small-sphere: the larger the size ratio is, the larger the coupling effect will be.

**Keywords:** Coupling Effects of Depletion Interactions; Three-Sphere Systems; Size Ratio

## 1. Introduction

Recently, coupling effects of depletion interaction were reported in the three-sphere system or in the two-sphere system confined by two plates [1-3]. It was pointed out that when a large-sphere is suffered by more than one depletion force, these depletion forces will couple with each other and results in a strengthened depletion force. As a new finding, the coupling effect is a very interesting and significant topic for us to study, since the dynamic behavior of particles and the structure of the colloidal system will be affected by the strengthened depletion interactions. With the aid of the concept of excluded volume, Asakura and Oosawa (AO) suggested that, the mechanism about the depletion force is just related to geometrical factor, therefore it is very simple [4]. However, the coupling effect reminds us that it is not as simple as what they told. Obviously, further investigations on the coupling effects of the depletion forces are needed. For simplicity, in previous studies, the size ratio of large- to small-sphere of these systems is almost unchanged [5-14]. On the other hand, the coupling effect will be affected by the parameter of size ratio, from which the nucleation of cluster is also affected: the nucleation packing of cluster will be easier to go on if it increases with the increase of size ratio of large- to small-sphere. In this paper, through the three-sphere system schematically illustrated by **Figure 1(a)**, the relation between the coupling effect and size ratio of large- to small-sphere will

be taken into consideration. In this system, the large sphere B is placed between the two spheres A and C. Apparently it is suffered two opposite depletion forces from A and C at the same time. However, it was found that, in this three-sphere system, the depletion forces acted on B are larger than that of the two-sphere system described by **Figure 1(b)**, where the sphere C is taken out of from the three-sphere system. To describe this coupling effect of depletion interactions clearly, a parameter should be introduced. According to the definition of three-body interaction, the difference of the depletion forces in the three systems demonstrated by **Figures 1(a)-(c)** can be used to describe the coupling effect.

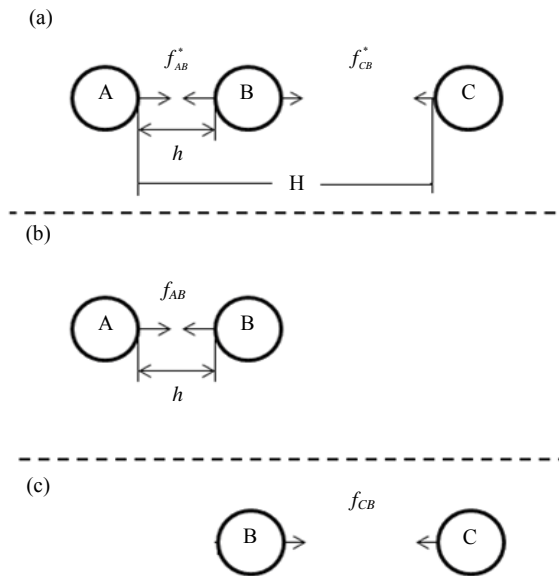
In this paper, the model and theory are presented in Section 2, and the numerical results and a simple discussion will be given in Section 3, finally a summary of our study will be given in Section 4.

## 2. Theory and Model

It is well known that the hard spheres mixture is characterized by the pair potential of

$$u(d) = \begin{cases} \infty & d < d_{ij} = (\sigma_i + \sigma_j)/2 \\ 0 & d > d_{ij} \end{cases} \quad (1)$$

where  $d$  is the distance between the two spheres in diameters  $\sigma_i$  and  $\sigma_j$ , respectively. The force exerted on the big sphere of radius  $R$  by a small sphere of radius  $r$



**Figure 1.** (a) Schematic diagram of the three-sphere system consisting of A, B and C; (b) is the two-sphere system where C is taken out of from (a); (c) is another two-sphere system where A is taken out of from (a). The dashed lines between (a) and (b), (b) and (c), are just used to separate the three systems.  $h$  is the separation of spheres A and B, and  $H$  the separation of A and C.

can then be written as  $k_B T \delta(d - d_{ij})$ . Consequently, the depletion force is the total force acted on the large sphere from the small spheres, and can be usually determined through the acceptance ratio method (ARM): if the potential and partition function of two systems are  $V_0, Q_0$  and  $V_1, Q_1$  respectively, where  $V_0$  and  $V_1$  are the external potentials corresponding to the two large spheres located at different positions, the free energy difference between these two systems is given by the following expression [6,11,15],

$$\beta \Delta F = -\ln \frac{Q_1}{Q_0} = -\ln \frac{\langle f[\beta(V_1 - V_0) + C] \rangle_0}{\langle f[-\beta(V_1 - V_0) - C] \rangle_1} \quad (2)$$

$$= -\ln \frac{N_{10}}{N_{01}}$$

where  $N_{10}$  is the number of samples drawn out from  $N$  simulated samples, which generated with the potential  $V_0$  where  $V_1$  is not infinite;  $N_{01}$  is the number of samples drawn out from  $N$  simulated samples, which generated with potential  $V_1$  where  $V_0$  is not infinite, and  $f(x) = [1 + \exp(x)]^{-1}$  is the Fermi function, and  $C$  is a constant which is usually set to a value of 0 for a hard sphere system.

We note that the depletion force can be determined by the differential of the free energy of the system obtained through Equation (2). Besides, it is important to emphasize that if there are many depletion forces acted on a

sphere at the same time, only the resultant force corresponding to the free energy will be determined by ARM. For the three-sphere system, a parameter is introduced to describe the coupling effect:

$$\Delta f = f - f_{AB} - f_{CB} \quad (3)$$

where  $f$  is the resultant force of the three-sphere system described by **Figure 1(a)**,  $f_{AB}$ ,  $f_{CB}$  are the two forces acted on large sphere B of the two-sphere systems described by **Figures 1(b)** and **(c)** respectively. In fact, in the three-sphere system, due to the presence of large spheres A and C, the two component forces acted on large sphere B are  $f_{AB}^*$ ,  $f_{CB}^*$ , and the relation between them is,

$$f = f_{AB}^* - f_{CB}^* \quad (4)$$

Obviously, if there is no coupling effect between  $f_{AB}$  and  $f_{CB}$ ,  $\Delta f = 0$ , otherwise  $\Delta f \neq 0$ , so  $\Delta f$  can be used to describe the coupling effect. For the sake of simplicity, we suppose that, in all the systems mentioned above, the positions of spheres A and C are fixed, only B moves from the position of the contact of A to the middle point of A and C, then the depletion forces  $f$ ,  $f_{AB}$ ,  $f_{CB}$  are determined by ARM through the systems described by **Figures 1(a)-(c)** respectively. The parameter  $h$  is the separation of A and B, and  $H$  is the separation of A and C.

### 3. Results and Discussions

In this section, only the unconfined three-sphere systems and the corresponding two two-sphere systems with size ratio of large- to small-sphere  $R/r = 3, 5$ , will be studied, respectively. In the simulations, the two or three large hard spheres are placed along  $z$  direction of the cell box of size  $L_x \times L_y \times L_z$ , but the small spheres are randomly distributed around the macro spheres to form a fluid; the number of micro-ions  $N$  is determined by the given volume fraction  $\eta$ , defined as  $\eta = \frac{NV_c}{(V - 3V_m)}$  or

$$\eta = \frac{NV_c}{(V - 2V_m)}$$

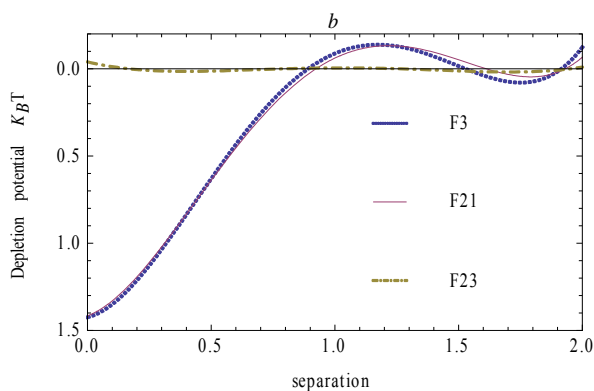
respectively, where  $V = L_x \times L_y \times L_z$  is the total volume of

the cell box,  $V_c = \frac{4}{3} \pi r^3$  is for the volume of the micro-

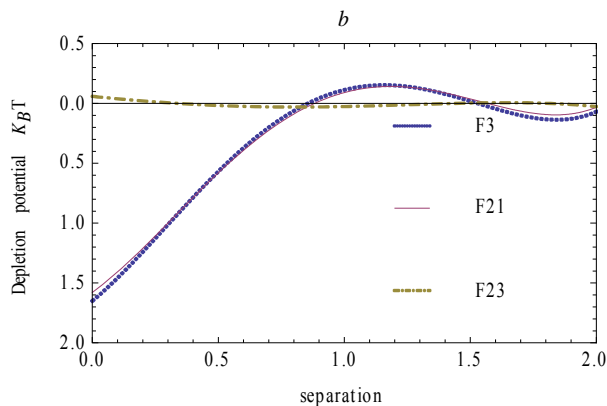
sphere,  $V_m = \frac{4}{3} \pi R^3$  denotes the volume of the macro-

sphere. As unconfined systems, the period boundary condition is applied to all the three directions of  $X, Y$  and  $Z$  in the Monte Carlo simulations. Then the configurations of the micro-spheres will be sampled according to the Metropolis algorithm [15] with the two macro-spheres A and C fixed while B is moved from the contact of A to the middle point of A and C. Each micro-sphere

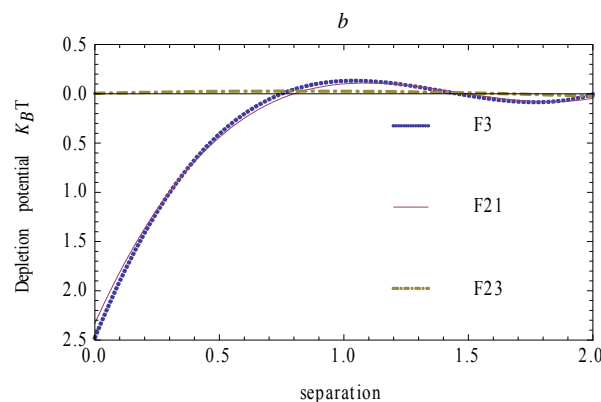
is orderly chosen involving a trial displacement. Except for an overlapping with the macro-spheres and the other micro-spheres, the new position of the micro sphere is randomly accepted. In our simulations,  $1.0 \times 10^5$  Monte Carlo steps (MCS) are used for the equilibrium of the system and other  $3.0 \times 10^5$  MCS to collect data, and the depletion potential is set as 0 while the two macro-spheres A and B are at contact, *i.e.*,  $h = 0$ . In this way, the depletion potentials, then the corresponding depletion forces in the unconfined, two- and three-sphere systems,  $f_{AB}$ ,  $f_{CB}$ , and  $f$ , are determined by ARM respectively. For the sake of simplicity, the cases corresponding the unconfined three-sphere systems with  $R/r = 3, 5$ , and volume fraction  $\eta = 0.12$  and  $0.23$ , were studied, respectively, and the corresponding depletion forces are shown in **Figures 2-5**, respectively. In these figures, the depletion force  $F$  is plotted as a function of  $h$ , and the units of  $F$  and  $h$  are  $\pi R \rho k_B T$  and  $2r$ , respectively, where  $\rho$  is the number density of the small sphere,  $k_B$  the Boltzmann constant,  $T$  the absolute temperature. In **Figures 2-5**, the short-dashed lines denote the depletion forces of the three-sphere system, the solid and long-dashed lines denote the two-sphere systems described by **Figures 1(b)** and **(c)**, respectively. From **Figures 2-5**, it's evident that, 1) no matter the volume fraction is large or small, the depletion force between spheres B and C is almost 0; 2) the depletion force of the three-sphere system is larger than that of the other two-sphere system (consisting of A and B). Furthermore, according to Equation (4), the corresponding depletion force difference of the three-sphere system and the two corresponding two-sphere systems  $\Delta f$  are also determined and there-sults for  $\eta = 0.12$  and  $0.23$  are shown in **Figures 6** and **7**, in which the dashed and solid lines are for the confined and unconfined systems respectively. From **Figures 6** and **7**, it is evident that  $\Delta f \neq 0$ , so the deple



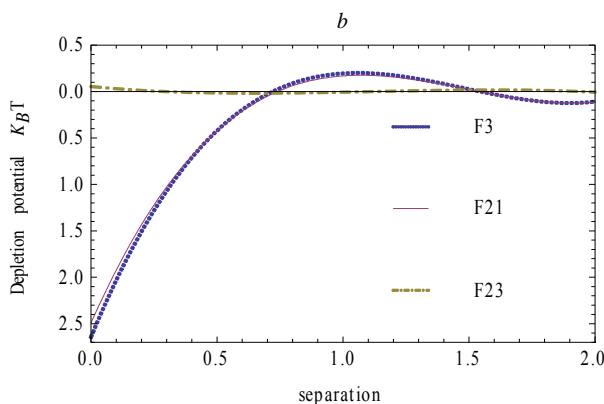
**Figure 2.** The depletion forces of the unconfined three-sphere system (described by the dotted lines) and of the corresponding two-sphere systems (by the solid and dot-dashed lines respectively). The volume fraction is  $\eta = 0.12$  and the size ratio is  $R/r = 3$ .



**Figure 3.** The depletion forces of the unconfined three-sphere system (described by the dotted lines) and of the corresponding two-sphere systems (by the solid and dot-dashed lines respectively). The volume fraction is  $\eta = 0.12$  and the size ratio is  $R/r = 5$ .

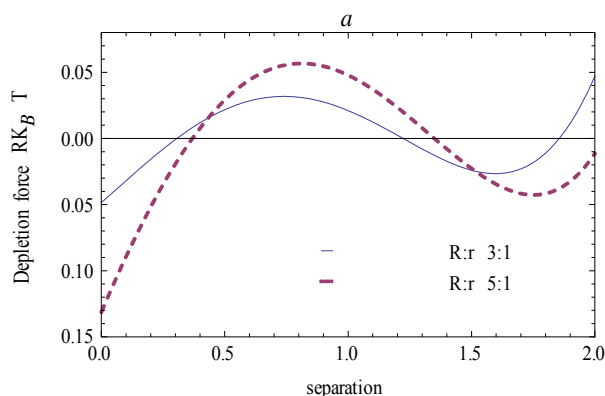


**Figure 4.** The depletion forces of the unconfined three-sphere system (described by the dotted lines) and of the corresponding two-sphere systems (by the solid and dot-dashed lines respectively). The volume fraction is  $\eta = 0.23$  and the size ratio is  $R/r = 3$ .

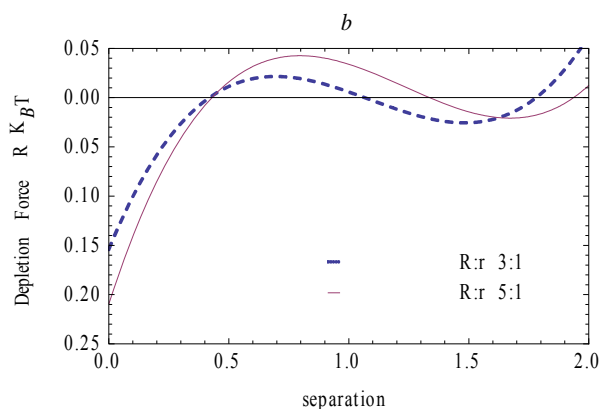


**Figure 5.** The depletion forces of the unconfined three-sphere system (described by the dotted lines) and of the corresponding two-sphere systems (by the solid and dot-dashed lines respectively). The volume fraction is  $\eta = 0.23$  and the size ratio is  $R/r = 5$ .

tion interactions acted on the sphere B are coupled with each other. In addition, from **Figures 6** and **7** we also find that, 1)  $\Delta f$  is enlarged when the volume fraction of the system increases from 0.12 to 0.23; 2)  $\Delta f$  is also enlarged when the size ratio of large- to small- sphere of the system of  $R/r$  is increased from 3 to 5. In other words, the larger the volume fraction is, the larger the coupling effect of the depletion interactions will be; the larger the size ratio is, the larger the coupling effect of the depletion interactions will be. This is very important for us to get the physical viewpoint of nucleation. Supposing in a binary colloidal system, free large-spheres are continually pushed together by small-spheres, a cluster consisting of the packed spheres is therefore built up. Furthermore, with the increasing size of the cluster, it will suffer more and more strengthened depletion interactions from the small-spheres, because both the depletion interactions and the coupling effect of these depletion interactions will increase with the increase of the size ratio of the large- or small sphere to cluster. As a result, the cluster will grow larger and larger, even a nucleation will turn



**Figure 6.** The difference of depletion forces of the three-sphere systems. The volume fraction is  $\eta = 0.12$  and the size ratio is  $R/r = 5, 3$  respectively.



**Figure 7.** The difference of depletion forces of the three-sphere systems. The volume fraction is  $\eta = 0.23$  and the size ratio is  $R/r = 5, 3$  respectively.

out. In other words, the coupling effect of depletion interactions is helpful for the nucleation packing.

## 4. Conclusion

In conclusion, we have studied the depletion interactions in the unconfined three-sphere systems with different size ratio through Monte Carlo simulations. It is found that this coupling effect is clearly and completely described by the parameter of the depletion force difference of the three-sphere system and the two corresponding two-sphere systems. It is also found that the coupling effect of depletion force is affected by both the size ratio and the volume fraction of the system.

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