

# Modeling and Analysis of Ladder-Network Transmission Lines with Capacitive and Inductive Lumped Elements

Ana Flávia G. Greco, Joaquim J. Barroso, José O. Rossi

Associated Plasma Laboratory, National Institute for Space Research (INPE), São José dos Campos, Brazil.  
Email: barroso@plasma.inpe.br

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## ABSTRACT

This paper examines the properties of wave propagation in transmission lines with periodic LC and CL cells, taking into account ohmic losses in resistors connected in series to lumped capacitors and inductors. First order time differential equations are derived for current and charge, thus allowing analysis of transient regimes of the lines being excited by a pulse of arbitrary shape. In particular we examine the propagation characteristics of periodic lines in which identical unit cells are repeated periodically and also discuss the interpretation of positive and negative phase velocities associated with the LC and CL topologies. Loss effects on the propagation bandwidths of both lines are also discussed, and it is shown that in the left-handed transmission line (CL configuration) the phase advance of the crest of the transmitted signal with respect to the source signal is due to the intrinsic dispersive nature of the CL line which, in contrast to the LC line, is highly dispersive at low propagation factors.

**Keywords:** Transmission Lines; Right-Handed (RH); Left-Handed (LH); Backward Wave; Metamaterials

## 1. Introduction

Characterized by negative constitutive parameters (electric permittivity and magnetic permeability) metamaterial [1-3] is an artificially produced material which exhibits properties not found in nature. Such materials are often referred to as LH (left-handed) media to express the fact that they permit propagation of electromagnetic waves with the electric field, the magnetic field and the propagation vector forming a left-handed triad, in contrast to conventional materials in which the triad of these three vectors follows the right-hand rule. However, showing losses and having operation limited to narrow frequency bands, it was soon found that such resonant structures, such as split-ring resonators, become difficult to implement for microwave applications [4,5]. But as early as 2002, there has been proposed [6-9] the use of transmission lines—structures that are not resonant, low-lossy and allows wide frequency range of operation—to implement and put into practice the concept of metamaterials [10-12].

From the above, this paper examines the properties of wave propagation in transmission lines with periodic CL cells (C and L denoting the series capacitance and shunt inductance in each cell). As will be discussed, such a CL

line exhibits a characteristic of left-handed (LH) medium in that the phase velocity is negative, namely, with direction opposite to the group velocity. In a comparative study, an analysis is made of its dual line LC (L and C denoting the shunt capacitance and series inductance in each cell). To this end, by assigning the charge in the capacitor and the current flowing in the inductor as state variables, first order time differential equations are derived for a general CL line and its dual configuration LC. The equations so derived allow the transient analysis of the lines being excited by a pulse of arbitrary shape. We also investigate the properties of propagation in the steady state sinusoidal regime and examine the dispersion characteristics of lossy lines by identifying the frequency bands and group delay in the frequency range 100 kHz to 1 MHz.

## 2. Circuit Equations

### 2.1. Line LC

Figure 1(a) shows a generalized LC line, where the assigned state variables are the loop current  $I_k$  in a generic section  $k$  and the corresponding charge  $Q_k$  stored in the capacitor. Of importance from the numerical standpoint, these variables were taken so that the line differential

equations are all first order. The line is further assumed to be lossy, where each inductor has an associated resistance  $r_{L,k}$ , each capacitor has a resistance  $r_{C,k}$  and the generator has a resistor  $r_s$ .

On the basis of Kirchhoff's laws of voltage and current, we then arrive at a generalized transmission-line equation that applies for any number  $N$  of sections  $k$ . For the LC line, the circuit equations are written for three categories of sections: starting, intermediate, and ending sections.

**Starting section:**

$$\begin{aligned} \frac{dI_k}{dt} &= \frac{V_s}{L_k} - \frac{r_s}{L_k}(I_k) - \frac{r_{L,k+1}}{L_k}(I_k) \\ &\quad - \frac{r_{C,k}}{L_k}(I_k - I_{k+1}) - \frac{Q_k - Q_{k+1}}{L_k C_k} \end{aligned} \quad (1)$$

$$\frac{dQ_k}{dt} = I_k$$

where  $k = 1$  and  $V_s$  is the input voltage, which may be a sinusoidal signal or a pulse of arbitrary shape (trapezoidal, triangular, etc.).

**Intermediate segment:**

$$\begin{aligned} \frac{dI_k}{dt} &= -\frac{r_{L,k}}{L_k}(I_k) - \frac{r_{C,k-1}}{L_k}(I_{k-1} - I_k) \\ &\quad - \frac{r_{C,k}}{L_k}(I_k - I_{k+1}) + \frac{Q_{k-1} - Q_k}{L_k C_{k-1}} - \frac{Q_k - Q_{k+1}}{L_k C_k} \end{aligned} \quad (2)$$

$$\frac{dQ_k}{dt} = I_k; \quad k = 2, 3, \dots, N-1$$

**Ending segment:**

$$\begin{aligned} \frac{dI_k}{dt} &= -\frac{r_{L,k}}{L_k}(I_k) - \frac{r_{C,k-1}}{L_k}(I_{k-1} - I_k) - \frac{r_{C,k}}{L_k}(I_k) \\ &\quad + \frac{Q_{k-1} - Q_k}{L_k C_{k-1}} - \frac{Q_k}{L_k C_k} \end{aligned} \quad (3)$$

$$\frac{dQ_k}{dt} = I_k; \quad k = N$$

**2.2. Line CL**

Unlike the previous case, in the CL topology (Figure 1(b)) the state-variable current flows in the shunt branch  $L$ - $r_L$  of each section, so that that the CL line is described by a pair of equations:

$$\begin{aligned} \frac{dI_k}{dt} &= \frac{V_s}{L_k} - \frac{r_s}{L_k} \sum_{j=1}^N I_j - \frac{r_L}{L_k}(I_k) - I_{s_k} - Q_{s_k} \\ \frac{dQ_k}{dt} &= I_k \end{aligned} \quad (4)$$

where  $I_{s_k}$  and  $Q_{s_k}$  are recursive functions defined as

$$I_{s_k} = I_{s_{k-1}} + \text{mult}[k] \sum_{j=k}^N I_j, \quad \text{with } I_{s_0} = 0 \quad (5)$$

$$Q_{s_k} = Q_{s_{k-1}} + \frac{\sum_{j=k}^N Q_j}{\text{den}[k]}, \quad \text{with } Q_{s_0} = 0$$

in which  $\text{mult}[k]$  is the  $k$ -th element of the array  $\{r_{c,j}/L_k, j=1, N\}$  and the denominator term  $\text{den}[k]$  is the  $k$ -th element of  $\{L_k C_j, j=1, N\}$  as detailed in the Appendix.

**3. Dispersion Relations**

**3.1. Line LC**

To analyze the steady-state sinusoidal regime of a periodic line with elements  $r_{C,k} = r_C, r_{L,k} = r_L, C_k = C$  and  $L_k = L$ , we consider the spatial and temporal variations in the form  $\exp[-i(\omega t - \beta n)]$ , where  $n = x/p$ , with  $x$  denoting the position of a node with respect to the source at  $x = 0$ ,  $p$  representing the hypothetical length of the unit cell;  $\beta$  is the propagation factor, a complex quantity which quantifies the attenuation and phase shift across each cell. Then the voltage and current equations for the LC line (Figure 1(a)) are

$$\begin{aligned} CU_k &= Q_k \\ U_k &= V_k - (I_k - I_{k+1})r_C = V_k - I_k(1 - e^{i\beta})r_C \\ Q'_k &= I_k - I_{k+1} \\ V_{k-1} - V_k &= LI'_k + r_L I_k \end{aligned} \quad (6)$$

where  $U_k$  is the voltage across the shunt capacitor of the  $k$ -th section. Noting that the time derivate of the currents  $I_k$  is  $I'_k = -i\omega I_k$ , from the assumption of time-harmonic variation, and  $I_{k\pm 1} = I_k e^{\pm i\beta}$ , from the periodic boundary conditions, Equation (6) takes the form

$$\begin{aligned} I_k(1 - e^{i\beta})(1 - i\omega r_C C) &= -i\omega C V_k \\ V_k(e^{-i\beta} - 1) &= I_k(r_L - i\omega L) \end{aligned} \quad (7)$$

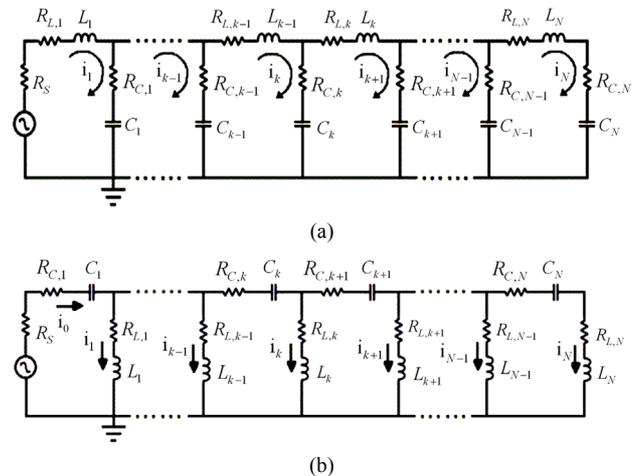


Figure 1. Line topologies (a) LC and (b) CL.

which leads to the dispersion relation

$$\frac{\tilde{\omega}^2 + i\tilde{\omega}a}{1 - i\tilde{\omega}b} = 4\sin^2(\beta/2) \quad (8)$$

with  $\tilde{\omega} = \omega/\omega_0$ ,  $\omega_0 = 1/\sqrt{LC}$ ,  $a = r_L/Z_0$ ,  $b = r_C/Z_0$ , and  $Z_0 = \sqrt{L/C}$ .

For a lossless LC line ( $a = 0$ ,  $b = 0$ ), Equation (8) reduces to

$$\tilde{\omega}^2 = 4\sin^2(\beta/2) \quad (9)$$

where the propagation factor  $\beta$  is real for  $\omega < 2\omega_0$  and becomes complex, when  $\omega > 2\omega_0$  as illustrated by the blue curves in **Figure 2**. In this plot, both the attenuation and phase shift per section are normalized to  $\pi$ ; for example at  $\omega/\omega_0 = 1$ , the phase shift per section is read as  $0.331\pi$ . Having no cutoff frequency, the RH periodic line exhibits the behavior of a low-pass filter. Propagation starts from zero frequency until the critical frequency  $\omega_c = 2\omega_0$  with no attenuation, as demonstrated by vertical branch of the dashed blue line ranging from  $\omega = 0$  to  $\omega_c$ . But for lossy lines, attention coexists at all frequencies since the attention curves bend rightwards. In addition for frequencies up to  $\omega_c \leq \omega_0$  the RH line behaves as nondispersive medium, in which the propagation factor  $\text{Re}\{\beta\}$  is a linear function of frequency.

### 3.2. Line CL

Analogously to the previous case, the voltage and current equations for the CL line (**Figure 1(b)**) are written as

$$\begin{aligned} V_k - V_{k+1} &= L(I'_k - I'_{k+1}) + r_L(I'_k - I'_{k+1}) = \frac{Q_{j+1}}{C} + r_C I_{j+1} \\ V_{k-1} - V_k &= L(I'_{k-1} - I'_k) + r_L(I'_{k-1} - I'_k) = \frac{Q_j}{C} + r_C I_j \end{aligned} \quad (10)$$

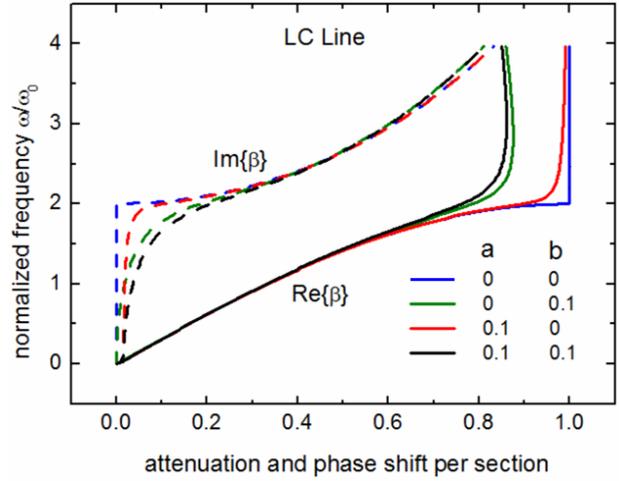
to give

$$\frac{1 - i\tilde{\omega}b}{\tilde{\omega}^2 + i\tilde{\omega}a} = 4\sin^2(\beta/2) \quad (11)$$

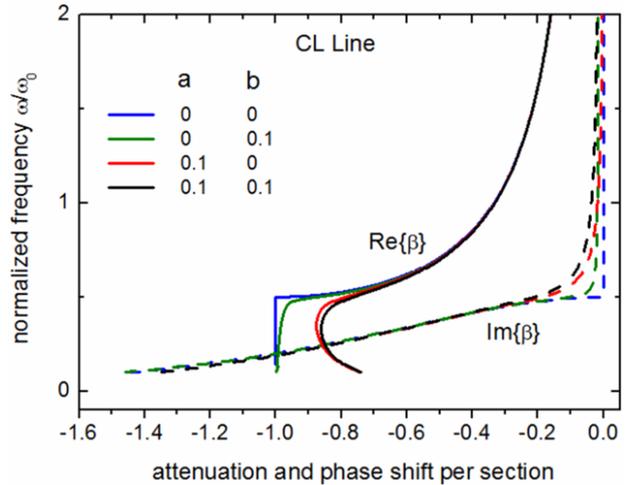
which for the lossless case reduces to

$$\tilde{\omega} = -\frac{1}{2\sin(\beta/2)} \quad (12)$$

The presence of the negative signal (coming from the ambiguity of  $\pm$  square root sign) is justified by examining Equation (12) in the limit  $\beta \ll 1$ , which leads us to  $\beta = \pm 1/\omega\sqrt{LC}$ , where  $\beta/p$  is the propagation constant of this CL line of concentrated elements. But from the theory of transmission lines [13], we know that  $i\beta$ , relates to the line immittances by  $\beta = i\sqrt{ZY}$ . The present case, in which  $Z = 1/(-i\omega C)$  and  $Y = 1/(-i\omega L)$ ,



**Figure 2. Dispersion/attenuation diagram for the LC line. In lossy lines ( $a, b \neq 0$ ), propagation and attenuation coexist at all frequencies. In the horizontal axis, the attenuation  $\text{Im}\{\beta\}$  (dashed curves), and phase shift per section  $\text{Re}\{\beta\}$  (solid curves) are both normalized to  $\pi$ .**



**Figure 3. Dispersion/attenuation diagram for the CL line.**

gives  $\beta = i\sqrt{\frac{i}{\omega C} \frac{i}{\omega L}} = \frac{-1}{\omega\sqrt{LC}}$  in the limit  $\beta \ll 1$ .

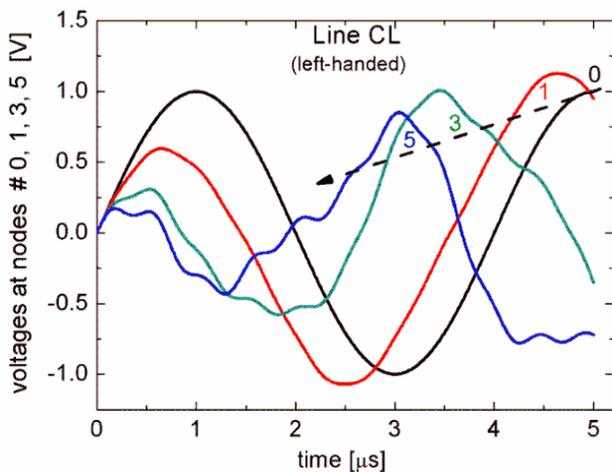
Thus, the dispersion relation of the CL line satisfies the conditions for excitation of backward wave, whose phase velocity is negative as illustrated in **Figure 3**. In the frequency range  $\omega > \omega_0/2$  for the lossless line (blue curves),  $\beta$  is purely real, representing propagation without attenuation in the system; when  $\omega < \omega_0/2$ ,  $\beta$  is imaginary, indicating attenuation without energy dissipation. It is interesting to notice that the underlying features of the propagation curve for the CL line, with the low-frequency branch bending downwards, have been experimentally observed in left-handed transmission lines yielding strong negative refraction index as large as  $-700$

at gigahertz frequencies [14].

### 4. Results

The following results refer to signals which propagate in the periodic LC and CL lines driven by a purely sinusoidal signal with 1-V amplitude.

**Figure 4** shows the time variation of signals propagating in a low loss ( $r_s = r_{LC} = r_{CL} = 0.001 \Omega$ ) CL line of twenty sections with  $L = 1 \mu\text{H}$ ,  $C = 1 \mu\text{F}$  and being excited by a signal of 250 kHz. Black, red, green and blue curves correspond to the nodes #0, #1, #3 and #5, where the signals are calculated. Note that the blue signal (node #5) reaches the positive peak ( $t \approx 3 \mu\text{s}$ ) before the other signals at the preceding positions {0, 1 and 3} closer to the exciting source. This is the phenomenon of phase advance in a CL line, and clearly apparent by the peak position at 3  $\mu\text{s}$  in the blue curve, before the other signals, which only reach their respective peaks after 3.5  $\mu\text{s}$ . Similar situation is displayed in **Figure 5** showing the time variation of signals propagating in a lossy line CL ( $r_s = r_{LC} = r_{CL} = 0.05 \Omega$ ) and excited by a signal of 200 kHz. We note that the signals distort themselves early in time ( $t < 2 \mu\text{s}$ ) so as to occur a pulse sequence of decreasing order {5, 3, 1, 0} at a time later to 8  $\mu\text{s}$ , when all signals turn out to follow sinusoidal variation. In **Figure 5**, the time difference between the crests of signals #5 and #0,  $t_0 = 11.21 \mu\text{s}$  and  $t_5 = 7.94 \mu\text{s}$ , namely,  $\Delta t = -3.27 \mu\text{s}$ , relates to the phase factor through  $\Delta t = (\beta/\omega)n$ , where  $n$  is the number of sections between two given nodes. From **Figure 3**, at  $\omega/\omega_0 = 200/159 = 1.26$  one reads  $\beta/\pi = -0.26$ ; then substituting  $n = 5$  and  $f = 200$  kHz (driving signal frequency) into  $\Delta t = [(\beta/\pi)/2f]n$  gives 3.25  $\mu\text{s}$ , in excellent agreement with the time difference  $\Delta t = -3.27 \mu\text{s}$  directly measured from **Figure 5**.

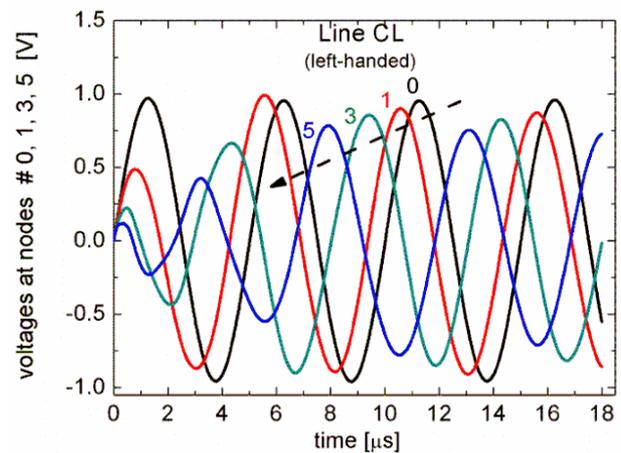


**Figure 4.** Voltage signals at nodes 0, 1, 3, and 5 in a low-loss, twenty-section CL line ( $r_s = r_{LC} = r_{CL} = 0.001 \Omega$ ,  $L = 1 \mu\text{H}$ ,  $C = 1 \mu\text{F}$ , resonance frequency  $f_0 = 159.0$  kHz) driven by a sinusoidal signal of 250 kHz.

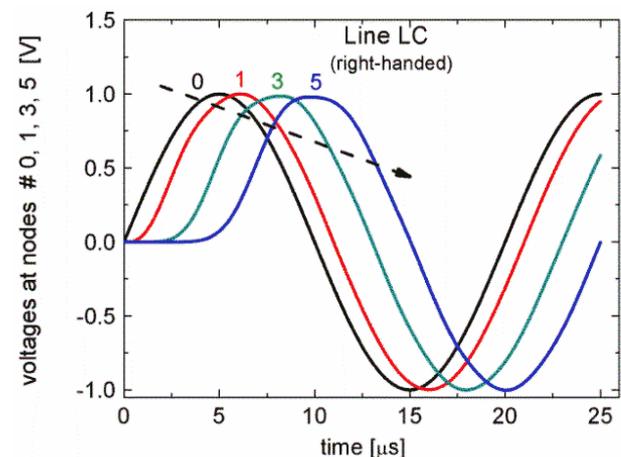
**Figure 6** shows the time variation of signals propagating in a low-loss ( $r_s = r_{LC} = r_{CL} = 0.001 \Omega$ ) LC line with twenty sections, with immittances  $L = 1 \mu\text{H}$ ,  $C = 1 \mu\text{F}$  and excited by a sinusoidal signal of 50 kHz. Black, red, green and blue curves correspond to the nodes #0, #1, #3 and #5. It appears that the phase velocity is positive, and the phase delay in the low-loss LC line is clearly indicated by the relative position of the peaks, where the blue signal, for example, reaches a peak in the first 10  $\mu\text{s}$  after the signals 0, 1 and 3 have reached their respective peaks. In a lossy LC line ( $r_s = r_{LC} = r_{CL} = 0.05 \Omega$ ) excited by a 20 kHz signal (**Figure 7**) it is noted that the same blue pulse (node #5) has the smallest amplitude in the sequence of pulses {0, 1, 3, 5}.

### 5. Conclusions

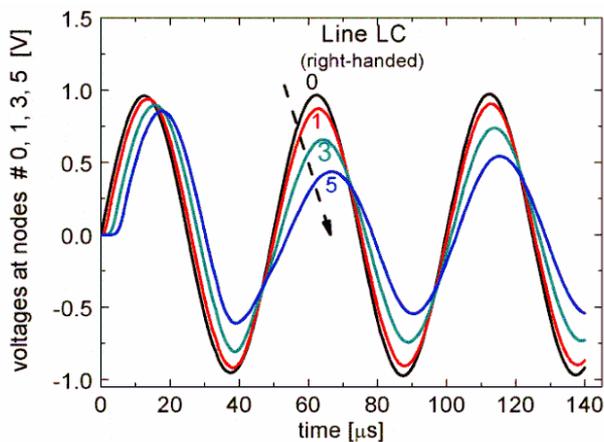
A system of time-varying ordinary differential equations



**Figure 5.** Voltage signals at nodes 0, 1, 3, and 5 in a lossy CL line of twenty sections ( $r_s = r_L = r_C = 0.05 \Omega$ ,  $L = 1 \mu\text{H}$ ,  $C = 1 \mu\text{F}$ ,  $f_0 = 159.0$ ) driven by a 200 kHz sinusoidal signal.



**Figure 6.** Voltage signals at nodes 0, 1, 3, and 5 in a low-loss LC line excited by a sinusoidal signal of 50 kHz.



**Figure 7. Voltage signals at nodes 0, 1, 3, and 5 in a LC lossy line excited by a sinusoidal signal of 20 kHz.**

for lumped-element transmission lines has been derived and solved with given initial conditions for the distributions of charge of the capacitors and the fluxes through the inductors.

We have investigated propagation phenomena in transmission lines consisting of LC and CL cells. By taking the inductor current and the charge stored in the capacitor as state variables, a system of first order equations has been derived for each topology of time (LC and CL) and numerically solved using the *Mathematica* software [15]. Considering both low-loss and lossy lines, in which identical inductors and capacitors are repeated periodically in each section, we have examined the transient and saturation processes which occur in the signals when excited by a sinusoidal source, thus clearly demonstrating the phase advance of signals that propagate in the CL line.

Arising from the periodicity of the lumped elements, a lossless CL line exhibits a sharp cutoff frequency at  $\text{Re}\{\beta\} = -\pi$  which originates a low-frequency band gap. Loss effects, however, make the gap disappear, such that the propagation curve turns out to be continuous and the bandwidth is enlarged, but at the expenses of high losses. On a complementary way, loss effects on LC lines decrease the passband, as demonstrated by the propagation/attenuation diagram in which energy dissipation starts to become significant for frequencies below the critical frequency  $\omega_c = 2\omega_0$  of the lossless LC line.

Finally we note that certain application of transmission lines demand a large number of cells, such as soliton excitation in nonlinear transmission lines typically requiring 500 cells, which may constitute a cumbersome task the drawing and wiring of 500 LC sections in SPICE-type circuit simulators [16]. On the other hand, a very large number of LC cells and their variant topologies are easily dealt with the mathematical formulation embodied in the system of differential equations.

## 6. Acknowledgements

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**Appendix**

We present in this appendix the equations implemented in the Mathematica software to the solution of differential equations that describe the temporal behavior and propagation of voltage and current signals in the LC (right-handed) and CL (left-handed) periodic lines.

**LC line: equation for the starting section**

*eqi*

$$= \text{Table} \left[ \left\{ i[k]'[t] = \frac{Vs[t]}{Ls[k]} - \frac{Rs}{Ls[k]} * i[k][t] - \frac{Rc[k]}{Ls[k]} * (i[k][t] - i[k+1][t]) - fN * \frac{(q[k][t] - q[k+1][t])}{Ls[k] * Cs[k]} - \frac{Rl[k]}{Ls[k]} * i[k][t], q[k]'[t] \right\}, \{k, 1, 1\} \right]$$

**LC line: equation for the intermediate section**

*eqs*

$$= \text{Table} \left[ \left\{ i[k]'[t] = \frac{Rc[k-1]}{Ls[k]} * (i[k-1][t] - i[k][t]) + fN * \frac{(q[k-1][t] - q[k][t])}{Ls[k] * Cs[k-1]} - \frac{Rc[k]}{Ls[k]} * (i[k][t] - i[k+1][t]) - fN * \frac{(q[k][t] - q[k+1][t])}{Ls[k] * Cs[k]} - \frac{Rl[k]}{Ls[k]} * i[k][t], q[k]'[t] \right\}, \{k, 2, NN-1\} \right];$$

**LC line: equation for the ending section**

*eqf*

$$= \text{Table} \left[ \left\{ i[k]'[t] = \frac{Rc[k-1]}{Ls[k]} * (i[k-1][t] - i[k][t]) + fN * \frac{(q[k-1][t] - q[k][t])}{Ls[k] * Cs[k-1]} - \frac{Rc[k]}{Ls[k]} * (i[k][t]) - fN * \frac{q[k][t]}{Ls[k] * Cs[k]} - \frac{Rl[k]}{Ls[k]} * (i[k][t]), q[k]'[t] \right\}, \{k, NN, NN\} \right]$$

**CL line: recursive functions for current and charge**

$$\text{denq} = \text{Table} [Ls[k]Cs[j], \{j, 1, NN\}];$$

$$\text{multi} = \text{Table} \left[ \frac{Rc[j]}{Ls[k]}, \{j, 1, NN\} \right];$$

$$Qs[k_] := Qs[k-1] + fN * \frac{\sum_{j=k}^{NN} q[j][t]}{\text{denq}[[k]]}$$

$$Is[k_] := Is[k-1] + \text{multi}[[k]] * \sum_{j=k}^{NN} i[j][t]$$

**CL line: general equations for currents and charges**

*eqs*

$$= \text{Table} \left[ \left\{ i[k]'[t] = \frac{Vs[t]}{Ls[k]} - \frac{Rs}{Ls[k]} * \sum_{j=1}^{NN} i[j][t] - \frac{Rl[k]}{Ls[k]} * i[k][t] - Is[k] - Qs[k], q[k]'[t] \right\}, \{k, 1, NN\} \right]$$

The initial conditions used are  $I_k(0) = 0$  and  $Q_k(0) = 0$ ; notice that the circuit elements are indexed to any values  $r_L, r_C, r_s, L_k, C_k$  in each section  $k$ .