Effect of deformation rate on the mechanical properties of arteries

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ABSTRACT

Pig aorta samples were tested uniaxially and equibiaxially at deformation rates from 10 to 200 %/s. Under uniaxial and biaxial testing, loading forces were reduced up to 20% when the deformation rate was increased from 10 to 200 %/s, which is the opposite to the behaviour seen in other biological tissues. A rate-dependent isotropic hyperelastic constitutive equation, derived from the Mooney-Rivlin model, was fitted to the experimental results (e.g. aorta specimens) using an inverse finite element technique. In the proposed model, one of the material parameters is a linear function of the deformation rate. The inverse relationship between stiffness and deformation rate raises doubts on the hypothesized relationship between intramural stress, arterial injury, and restenosis.

Keywords: Mechanical Properties; Artery; Uniaxial & Biaxial Testing; Deformation Rate; Viscoelasticity; Constitutive Model

1. INTRODUCTION

The knowledge of the viscoelastic properties is important to predict the biomechanical behaviour of soft tissues. To model their viscoelastic behaviour, first one performs appropriate mechanical tests to characterize deformation-rate effects, and then one selects a constitutive equation capable of representing those effects. Material parameter estimation is fundamental for posterior simulation of soft tissue at boundary conditions not selected in the experimental protocol.

The effect of deformation rate on the mechanical properties of soft biological tissues has been investigated, in particular for ligaments [1-7], tendons [4,7-9], spines [10-13], bones [14-17], liver [18], heart valves [19,20] and myocardium [21,22]. Most biological tissues stiffen with increasing deformation rate [4,7,9,11-13, 17]. This

time-dependent behavior has been described by viscoelastic constitutive models [6,22-26]. However, it was recently demonstrated that some biological tissues, such as liver, myocardium and skin, soften with increasing deformation rate [18,22]. Deformation rate effects of arteries, in particular thoracic aorta, were not included in previous studies.

Overstretch injury to the arterial wall during an angioplasty or stenting procedure has been shown to be correlated to the incidence of restenosis, i.e. in-growing tissue re-blocking the artery lumen [26,27]. Based on the hypothesis that lower deformation rate results in lower intramural stresses, slow balloon inflation has been proposed as a means to minimize vascular injury and reduce restenosis incidence [28]. Early studies did not conclude there was any difference in restenosis rates between conventional and slow balloon inflation [28-30], while some observed better immediate results [31,32]. In more recent studies, significantly lower restenosis rates were clinically observed with slow balloon inflation [33,34]. Slow stent deployment has also been proposed as a means to minimize arterial injury [35].

Finite element simulation of angioplasty and stenting can be used to optimize angioplasty procedure parameters, such as inflation pressure [36-40]. Optimization of inflation pressure rate requires accurate constitutive modeling of artery behavior including the effect of deformation rate. Numerous experimental studies have been performed to characterize the mechanical behaviour of arteries *in vitro* [41-44]. However, only a single deformation rate was used.

The objective of this study is thus to measure and model the effect of deformation rate on the tensile behavior of the arteries (e.g. pig aortas). In other words, the intention of this paper is to investigate experimentally the dependence of uniaxial and biaxial extensional stretching of arterial wall on the deformation-rate, and consequently to modeling the experimental data by means of an appropriate constitutive equation.



2. MATERIALS AND METHODS

2.1. Experimental Setup

Five thoracic aortas were harvested within the day of death of pigs from a local slaughterhouse and cleaned of remaining connective tissue. Then each artery was cut open along its length, and cut out in rectangular and cruciform-shaped specimens. The thickness of all specimens was measured with a vernier caliper (mean 2.4 mm, standard deviation 0.2 mm). Twelve samples were obtained from the five aortas: 8 rectangular samples (4 were cut in circumferential direction and 4 in axial direction) and 4 cruciform samples. Samples were stored in isotonic saline solution at 4° C for no longer than 8 hours prior to testing.

Rectangular and cruciform samples were used for uniaxial and biaxial testing, respectively. Rectangular samples were 40 mm long and 4 mm wide, but only 20 mm of the sample were free to deform. The distance between grips in cruciform samples was 40 mm. A nominal stretch ratio of 1.5 was applied to avoid permanent deformation on the tissue.

Uniaxial and biaxial testing was carried on a planar biaxial test bench (Bose Corporation, Minnetonka, MN) capable of applying a peak force of 200 N over a displacement range of 12 mm per actuator. A saline bath maintained at body temperature $(37^{\circ}C)$ was used. Samples were mounted in horizontal configuration with the help of grip clamps to the arms extending from the actuators over the top of the bath.

Each sample was subjected to triangular wave form displacements of 1.5 stretch ratio of amplitude, and deformation rates of 10, 50, 100, 120, 140, 160, 180 and 200%, which correspond to frequencies of 0.1, 0.5, 1.0, 1.2, 1.4, 1.6, 1.8 and 2.0 Hz. Half of the samples were subjected to deformation rates in the following order: 160, 120, 50, 200, 140, 100, 10 and 180 %/s (randomly tested), and the other half were tested from 10 to 200 %/s in ascending order in order to be confident that the deformation rate effect observed do not depend on the testing procedure order. In vivo, the artery is constantly submitted to stresses while the tissue is inflated-deflated due to blood pressure. In order to mimic in vivo loading conditions during in vitro tests, pre-conditioning cycles are required. Each test lasted for ten cycles. The first five cycles were used to pre-condition the tissue in order to reach a steady-state behavior. The force-stretch data was

averaged over the last five cycles.

2.2. Statistical Analysis

The experimental data obtained in uniaxial and biaxial testing are represented in mostly all figures in this chapter as the mean value \pm the standard errors (i.e. Figure 3). The medians rather than the means are plotted only in Figure 6. As shown in Table 1, twenty loading-unloading cycles were used to obtain the circumferential and axial force-stretch mean curves per deformation rate and testing condition. Circumferential and axial forces at 1.5 stretch ratio from both uniaxial and biaxial tests were used to calculate the medians, ranges (percentile between 75% - 25%) and p values with the help of the statistical tests selected. A multiple comparison between forces belonging to all groups (i.e. experimental forces obtained at each deformation rate belong to one group) was performed to study the significance of the effect of deformation rate using the ANOVA Friedman statistic test for dependent variables. The Wilcoxon test was also used to compare all deformation rate forces by pairs (i.e. forces at the highest deformation rate, 200%/s, with forces at the lowest deformation rate, 10%/s). In order to identify if the effect of deformation rate is more significant than the effect of anisotropy, circumferential and axial forces at each deformation rate were also compared using the Wilcoxon test. A 0.05 level of significance (p = 0.05) were selected for the Friedman and Wilcoxon tests.

2.3. Inverse Modeling

An inverse modeling technique was used to adjust the parameters of the proposed constitutive model (described below in the results section). The reaction forces at the boundaries (grips) were predicted by finite element simulation of the experiment consisting of applying displacement conditions on a mesh of the same size and shape as the sample. In an optimization loop, the material properties were iteratively adjusted until the following objective function was minimized, which occurs when a set of force-displacement experimental data matches the values calculated by the model:

$$S(c) = S(c)_{Biaxial} + S(c)_{Uniaxial}$$
(1)

The objective function to be minimized is the sum of uniaxial and biaxial cruciform objective functions.

Table 1. Number of force-stretch experimental data averaged per deformation rate.

Testing condition	Specimen direction	Number of specimens	Number of cycles per specimen	Total number of cycles
Uniaxial	Circumferential	4	5	20
	Axial	4	5	20
Biaxial	-	4	5	20



Figure 1. Mesh and boundary conditions used in simultaneous simulation of uniaxial and biaxial testing.



Figure 2. Loading-unloading circumferential forcestretch cycles of a typical biaxial cruciform sample: Five steady-state cycles at a deformation rate of 100%/s.

$$S(c)_{Biaxial} = \sum_{i=1}^{N} \left[\left(F_x(d_x, d_y) - f_x(d_x, d_y, c) \right)^2 + \left(F_y(d_x, d_y) - f_y(d_x, d_y, c) \right)^2 \right]_i$$
(2)

where *c* is the vector of unknown material properties; d_x and d_y are the applied displacements; $F_x(d_x, d_y)$ and $F_y(d_x, d_y)$ are the experimentally measured reaction forces at the grips; and $f_x(d_x, d_y, c)$ and $f_y(d_x, d_y, c)$ are the reaction forces predicted by the finite element model in biaxial cruciform extension; *N* is the total number of data points gathered in the experiments.

$$S(c)_{Uniaxial} = \sum_{i=1}^{N} \left[\left(F(d) - f(d,c) \right)^{2} \right]_{i}$$
(3)



Figure 3. Sample to sample variability of thoracic aorta. Open diamonds and squares represent the mean force-stretch behavior at 10 and 200 %/s, respectively (n = 4): (a) circumferential, and (b) axial directions. Standard error bars are also shown.

Uniaxial extension predicted forces (f(d,c)) were fitted to the average of circumferential and axial reaction forces (F(d)); d is the applied displacement in uniaxial extension.

A combined uniaxial-biaxial mesh was used to simultaneously simulate uniaxial and biaxial testing conditions (Figure 1). Only one quarter of the rectangular and the cruciform sample area free to deform between the grips was meshed due to symmetry. The uniaxial and biaxial cruciform mesh-sections are linked together at the origin. Boundary conditions are shown in Figure 1. In some cases, a small stretch shift was applied to the experimental results to obtain better agreement between uniaxial and biaxial fits.

3. RESULTS

3.1. Experimental Results

Figure 2 shows cycle-to-cycle variability between cycles 6 and 10, for a typical sample tested biaxially at a deformation rate of 100%/s. Cycle-to-cycle variability was observed to be small for all cases.

Figure 3 shows the standard error bars (sample-tosample variability) of circumferential and axial forces at the lowest (10 %/s) and highest (200 %/s) deformation rates. Similar standard errors were obtained with the other deformation rates.

Figure 4 shows the effect of deformation rate on the circumferential behavior of thoracic aorta from uniaxial and biaxial testing. Lower forces were observed at higher than at lower deformation rates. In particular, a 20% lower force was observed at 1.5 stretch ratio with the lowest deformation rate (10%/s) compared to the highest deformation rate (200%/s). The same phenomenon was observed in the axial direction (**Figure 5**).



Figure 4. Effect of deformation rate on the force vs. stretch ratio curve of arteries: Circumferential force vs. stretch ratio curves in uniaxial and biaxial testing.



Figure 5. Effect of deformation rate on the force vs. stretch ratio curve of arteries: Axial force vs. stretch ratio curves in uniaxial and biaxial testing.

Figure 6 shows the median, percentiles (75% - 25%) and the ANOVA Friedman test p value obtained using the forces at 1.5 stretch ratio for each deformation rate. Significant differences (p < 0.05) were observed, with the highest significant differences under biaxial testing ($p \le 0.002$).

Table 2 shows Wilcoxon test p values of uniaxial and biaxial data of some of the deformation rate pairs analyzed (i.e. comparison of the force at the maximum deformation rate with the forces at other deformation rates). Marginally significant differences were observed in the forces of mostly all deformation rates. Similar p values were obtained in the other comparisons (i.e. 50%/s vs. all deformation rates).

In Figure 7, circumferential and axial forces at 1.5 stretch ratio of a typical sample are plotted against deformation rate. In this figure, one can observe that the effect of deformation rate (20 % difference in force between smallest and highest deformation rate) is approximately twice as large as the effect of anisotropy (difference in force between circumferential and axial directions). This supports the hypothesis of modelling deformation-rate effects with an isotropic model.

Table 3 shows p values calculated using the Wilcoxon test comparing the circumferential force to the axial force at 1.5 stretch ratio per each deformation rate.

3.2. Constitutive Modeling

Based on the experimental results, it was assumed that the arterial wall behaves as a pseudoelastic material [45], where the stress is independent of the loading path. Thus, we propose a new rate-dependant isotropic hyperelastic model, based on the Mooney-Rivlin model [46] given by the following strain energy density function:

$$W(I_1, I_2) = \sum_{i,j=0}^{\infty} a_{ij} (I_1 - 3)^i (I_2 - 3)^j$$
(4)

The Mooney-Rivlin model is a function of the 1^{st} and 2^{nd} invariants (I_1, I_2) of the Green-Cauchy tensor:

$$I_1 = \lambda_1^2 + \lambda_2^2 + \lambda_3^2$$
 (5)

$$I_{2} = \lambda_{1}^{2}\lambda_{2}^{2} + \lambda_{2}^{2}\lambda_{3}^{2} + \lambda_{3}^{2}\lambda_{1}^{2}$$
(6)

where λ_1 , λ_2 and λ_3 are the principal stretch ratios.

In its most common expression, only five parameters are selected (a_{10} , a_{01} , a_{11} , a_{20} , a_{30}) and the strain energy function reduces to:

$$W = a_{10}(I_1 - 3) + a_{01}(I_2 - 3) + a_{11}(I_1 - 3)(I_2 - 3) + a_{20}(I_1 - 3)^2 + a_{30}(I_1 - 3)^3$$
(7)

The underlying hypothesis for using the Mooney-Rivlin model is that one or some of the five parameters vary with deformation rate. This can be modeled by

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Figure 6. Deformation rate effect on (a, b) uniaxial and (c, d) biaxial forces at maximum stretch ratio ($\lambda = 1.5$): (a, c) circumferential and (b, d) axial direction.

Type of test	Uniaxi	al	Biaxia	ıl
Force	Circumferential	Axial	Circumferential	Axial
Deformation rate (%/s)			200	
10	0.07	0.07	0.07	0.07
50	0.07	0.07	0.07	0.07
100	0.14	0.07	0.07	0.07
120	0.07	0.07	0.07	0.07
140	0.07	0.07	0.07	0.07
160	0.07	0.07	0.07	0.07
180	0.37	0.07	0.72	0.42

 Table 2. Effect of deformation rate: Wilcoxon test p values.

Table 3. Effect of anisotropy (Circumferential vs. axial forces, n = 4): Wilcoxon test p values.

Deformation rate (%/s)	Uniaxial	Biaxial
10	0.47	0.47
50	0.14	0.47
100	0.27	0.47
120	0.14	0.47
140	0.27	0.47
160	0.14	0.47
180	0.14	0.47
200	0.14	0.47

	M-R Fitted Parameter Values (kPa)							
Deformation Rate (%/s)	a_{10}	a_{01}	a_{11}	a_{20}	<i>a</i> ₃₀	SUM/n x 10^{-03} (n = 24)		
10	9.25	7.57	5.77	6.94	4.43	2.69		
50	9.21	7.45	5.45	6.58	3.64	2.09		
100	8.90	7.23	5.51	6.66	3.86	1.76		
120	9.46	7.66	5.38	6.52	3.03	1.66		
140	9.96	8.07	5.30	6.45	2.36	3.45		
160	9.22	7.47	5.29	6.40	3.24	2.72		
180	9.07	7.32	5.34	6.49	2.95	2.80		
200	8.86	7.17	5.24	6.36	2.80	1.90		
Linear Regression	n Coefficients	& R^2		$(a_{ij} = 2$	$(A + B\dot{\lambda})$			
Α	9.32	7.60	5.69	6.85	4.27			
B x 10 ⁻⁰³	-0.70	-0.90	-2.40	-2.50	-8.20			
\mathbb{R}^2	0.01	0.04	0.80	0.78	0.65			

Table 4. First simulation results (material parameters) of uniaxial and biaxial cruciform tests.

Table 5. Second simulation results (material parameters) of uniaxial and biaxial cruciform tests.

M-R Fitted Parameter Values (kPa)							
Deformation Rate (%/s)	a_{10}	a_{01}	a_{11}	a_{20}	a_{30}	$SUM/n \ge 10^{-03} (n = 24)$	
10					4.99	2.64	
50					3.66	2.18	
100					3.74	1.88	
120	0.24	7 70	5 41	6.55	3.16	1.68	
140	9.24	1.19	5.41	6.55	2.75	3.56	
160						3.06	2.83
180							2.72
200	200				2.26	1.98	
Linear Regressio	Linear Regression Coefficients & R^2 $\left(a_{ij} = A + B\dot{\lambda}\right)$						
А	-	-	-	-	4.76		
B x 10 ⁻⁰³	-	-	-	-	-12.20		
R^2	-	-	-	-	0.88		



Figure 7. Comparison of the effect of deformation rate versus the effect of anisotropy. Typical sample forces at maximum stretch ratio (1.5).

making these parameters functions of the deformation rate. Those relationships are explored in this work.

3.3. Inverse Modeling Results

Initially, all five Mooney-Rivlin parameters were allowed to vary while being limited to a 10^{-4} kPa to 10^{4} kPa

Other fits were performed assuming that $a_{11} = a_{20} = 0$ to see if the model could be reduced to a simpler form

to see if the model could be reduced to a simpler form **(Table 6)**. First, parameters a_{10} , a_{01} and a_{30} were allowed to vary. The highest regression coefficient (absolute value of *B*) and the highest correlationcoefficient (R^2) were obtained for a_{30} . Thus parameters a_{10}

value range. Table 4 summarizes the fitted parameters

obtained for the uniaxial and biaxial cruciform force vs.

stretch ratio curve averaged for each deformation rate.

lowed to vary because this was the parameter with the

highest slope (absolute value of B) in the linear regres-

sions reported in Table 4. The remaining parameters

were maintained constant by using the average of the fitted values obtained for all deformation rates. The fit-

ted parameters obtained are summarized in Table 5. The

sum of least squares (last column in Tables 4 and 5)

were almost unchanged when the first four parameters

were kept constant, supporting the hypothesis that only

 a_{30} varies with deformation rate.

In a second simulation only parameter a_{30} was al-

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and a_{01} were assumed to be constant (i.e. $a_{10} = 12.59$ kPa, $a_{01} = 9.54$ kPa) which were calculated as the average value of the fitted results obtained for all deformation rates. Parameter a_{30} was fitted again (Table 7). The sum of least squares remained small although higher than in the Table 4. Repeated optimization with different initial parameter values consistently converged towards the same solution.

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rameter a_{30} (Table 7) are plotted against deformation rate in Figure 8. This parameter decreased from 11 to 8 kPa when the deformation rate increased from 10 to 200 %/s, respectively. A quadratic polynomial was also used to model the variation of a_{30} with deformation rate but its correlation coefficient was similar to that obtained with a linear relationship, i.e. 0.89 and 0.88 respectively. Therefore, the linear fit was preferred over the polynomial fit due to its simplicity.

Uniaxial and biaxial simulations fitted values of pa-

Table 6. Third simulation results (material parameters) of uniaxial and biaxial cruciform tests.

M-R fitted parameter values (kPa)							
Deformation rate (%/s)	a_{10}	a_{01}	a_{11}	<i>a</i> ₂₀	<i>a</i> ₃₀	$SUM/n \ge 10^{-03} (n = 24)$	
10	12.90	9.72			10.76	2.61	
50	12.35	9.31			9.82	2.68	
100	12.39	9.35			9.86	2.61	
120	12.73	9.65	0.00	0.00	8.91	3.04	
140	13.38	10.14	0.00	0.00	7.92	5.17	
160	12.37	9.36			9.12	3.40	
180	12.27	9.35			8.74	4.96	
200	12.35	9.41			8.17	4.09	
<i>Linear regression coefficients</i> & R^2 $(a_{ij} = A + B\dot{\lambda})$							
А	12.76	9.61	-	-	10.66		
B x 10 ⁻⁰³	-1.50	-0.60	-	-	-12.50		
R^2	0.07	0.02	-	-	0.73		

Table 7. Final simulation results (material parameters) of uniaxial and biaxial cruciform tests.

M-R fitted parameter values (kPa)												
Deformation rate (%/s)	a_{10}	a_{01}	a_{11}	<i>a</i> ₂₀	<i>a</i> ₃₀	$SUM/n \ge 10^{-03} (n = 24)$						
10					10.91	2.62						
50					9.45	2.78						
100		0.54	0.00		9.53	2.70						
120	12 50			0.00	8.87	3.07						
140	12.59	9.54			8.41	5.33						
160											8.76	3.50
180					8.38	5.13						
200					7.85	4.21						
Linear regression	Linear regression coefficients & R^2											
А	-	-	-	-	10.66							
B x 10 ⁻⁰³	-	-	-	-	-13.70							
R^2	-	-	-	-	0.88							



Figure 8. Parameter a_{30} dependency on deformation rate for both, uniaxial and biaxial extensions.

Table 8. Linear regression adjusted parameters.

Paran	neter	Units	Uniaxial & biaxial cruciform
A	1	kPa	10.66
Ê	•	kPa*s	-2.06





Figure 9. Mooney-Rivlin computed reaction forces versus stretch ratios in uniaxial and biaxial testing at 10 and 200 %/s. An increase in the deformation rate decreases the reaction forces at a particular stretch ratio.

The following linear equation is proposed for parameter a_{30} :

$$a_{30} = f\left(\dot{\lambda}\right) = A + \hat{B}\left(\frac{\dot{\lambda}}{\lambda_N}\right) \tag{8}$$

Iteration

(b) (a) Minimization of the Objective Function Mooney - Rivlin Fitted Parameter 1. E+07 1. E+01 Case A ·Case_B 1. E+00 Case_C Objective Function 1. E+05 1. E-01 a30 (Mpa) 1. E+03 1. E-02 1. E-03 Case / 1. E+01 1. E-04 Case_B Case C 1. E-01 1. E-05 4 5 6 10 0 4 5 6 7 8 9 10

Figure 10. Optimization of the Mooney-Rivlin parameters for both uniaxial and biaxial tests at a deformation rate of 100%. (a) objective function, and (b) parameter a_{30} values as a function of iteration number, for three different initial guesses for material parameter a_{30} : Case_A (a_{30} , = 0.1 MPa), Case_B (a_{30} , = 1.0 MPa), and Case_C (a_{30} , = 0.0001 MPa).

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where $\dot{\lambda}$ is the deformation rate in s⁻¹, λ_N is the nominal stretch ratio ($\lambda_N = 1.5$), A and \hat{B} are the linear regression coefficients. The strain rate was normalized by the nominal stretch ratio, in order to account for the biological variability in peak stretch of aorta samples. Table 8 summarizes the adjusted parameters of Equation 8.

Figure 9 shows the computed forces at the lowest (10 %/s) and highest (200 %/s) deformation rates. Circumferential and axial forces were averaged in order to clearly exemplify the model prediction of the arterial behavior when the deformation rate is increased.

Although an excellent fit is observed between the biaxial data and biaxial simulation, the Modified Mooney-Rivlin model does not predict well the uniaxial behavior, especially in the nonlinear part of the curve at high stretch ratios (> 1.45).

Minimization of the objective function and parameter a_{30} estimation, using the data gathered at 100%/s deformation rate, is shown in Figure 10. Figure 10a shows the value of the objective function plotted against the iteration number in the optimization loop, for three different initial guesses. The objective function was evaluated 50 times in iterations 0 (initial guess) to 9. Figure 10b illustrates optimization of material parameter a_{30} .

Figure 11 shows the stress and stretch ratio distributions (at 1.5 stretch ratio) during uniaxial and biaxial testing. Scalar bars of the results at 10, 100 and 200 %/s deformation rates are shown for their comparison: One can observe the overall reduction of the stresses when the deformation rate was increased. The highest non uniform stress distribution was obtained in biaxial testing.

Iteration



Figure 11. Deformed configurations at 1.50 nominal stretch ratio for uniaxial and cruciform mesh: Distribution of (a) σ_{xx} stress and (b) stretch ratio in the principal direction. The stress and stretch distributions at 10, 100, and 200 %/s are shown inside the tables.



Figure 12. Parameter a₃₀ dependency on deformation rate for uniaxial extension: (a) circumferential and (b) axial directions.

3.3.2. Uniaxial Fit

In order to improve uniaxial simulations, two objective functions were defined, one for each direction:

$$S(c) = \sum_{i=1}^{N} \left[\left(F_x(d_x) - f_x(d_x, c) \right)^2 \right]_i$$
(9)

$$S(c) = \sum_{i=1}^{N} \left[\left(F_{y}(d_{y}) - f_{y}(d_{y}, c) \right)^{2} \right]_{i}$$
(10)

These objective functions were used to better predict the uniaxial extension experimental data. Different material properties were obtained for the circumferential and for the axial directions.

Uniaxial simulations fitted values of parameter a_{30} per deformation rates are shown in Figure 12. This parameter was reduced roughly from 33 to 23 kPa when the deformation rate was increased, following the same

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Figure 13. Mooney-Rivlin computed reaction forces versus stretch ratios in uniaxial testing at a deformation rate of 100 %/s.

Table 9. Mooney-Rivlin fitted parameters in uniaxial extension.

Maanay Divilin	Parameter	rs (kPa)	
Wooney-Kiviiii	Circumferential	Axial	
a_{10}	9.56	9.56	
a_{01}	18.01	7.87	
a ₁₁	0.00	0.00	
a ₂₀	0.00	0.00	
A	33.64	30.62	
\hat{B}	-5.32	-4.40	

(a) StressDistribution (Mpa) 10%/s 100%/s 200%/s 0 201 0.192 0 171 0.196 0.187 0.167 0.191 0.182 0.163 0.186 0.178 0.158 0.181 0.173 0.154 0.176 0.168 0.150 0.171 0.163 0.146 0.165 0.158 0.142 0.160 0.153 0.137 0.155 0.149 0.133 trend in circumferential and axial directions. It was also observed that material parameter a_{30} changed linearly as a function of the deformation rate.

Table 9 summarizes Mooney-Rivlin fitted parameters obtained just for uniaxial tests. The values of a_{30} obtained per deformation rate were almost 4 times greater than a_{30} values calculated with the combined uniaxial-biaxial simulation. Parameters a_{10} and a_{01} remained almost unchanged in both the uniaxial and combined uniaxial-biaxial simulations, with exception of parameter a_{01} in circumferential direction (i.e. 18 kPa).

The computed circumferential and axial forces obtained using the Mooney-Rivlin fitted parameters from **Table 9** in uniaxial tests, at a deformation rate of 100%/s, are shown in **Figure 13**. Uniaxial behaviour in both directions was well predicted.

Figure 14 shows the stress and stretch ratio distribution of rectangular samples. The reduction of the stresses as the deformation rate increases is observed when the stress distribution at each deformation rate is compared (scalar bars at 10, 100 and 200 %/s).

4. DISCUSSION

Contrary to the conventional understanding of biological

(b)		Stretch Ratio Distributi							
1	1	10%/s	1	00%/s	2	200%/s			
		1.514 1.503 1.493 1.482 1.471		1.516 1.507 1.498 1.489 1.480		1.516 1.507 1.498 1.489 1.480			
		1.461 1.450 1.439 1.429 1.418		1.471 1.462 1.454 1.445 1.436		1.471 1.462 1.453 1.444 1.435			

Figure 14. Deformed configurations at 1.50 nominal stretch ratio for uniaxial mesh using the material parameters of **Table 6.3** (Circumferential direction): distribution of (a) σ_{xx} stress and (b) stretch ratio in the principal direction. The stress and stretch distributions at 10, 100, and 200 %/s are shown inside the tables.

tissue behaviour, it was discovered in this study that the stiffness of thoracic aorta decreases with deformation rate, which was confirmed by the use of two statistical methods. Overall significant differences ($p \le 0.02$) in the mechanical forces of uniaxial and biaxial experiments were found using the ANOVA Friedman statistic test. In addition, the Wilcoxon test help us to observe which particular pair of deformation rates likely had a tendency to be significantly different (i.e. almost all deformation rates were different than 200%/s, p = 0.07). An increase

in the sample size for further studies will certainly reduce the Wilcoxon test p values obtained.

Giles *et al.* [22] observed a similar behaviour in load controlled planar biaxial laboratory tests of myocardium and skin samples. Hu and Desai [18] studied the variation of the elastic modulus of pig liver tissue at different deformation rates under compression tests, and found that at higher strain rates the liver has an apparently lower resistance (i.e. softer) to deformation than at lower deformation rates. However, under uniaxial displace-

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ment controlled tests, Pioletti et al. [7] found that an increase in the deformation rate results in a stiffer material response. In their work, they submitted ligament samples to deformation rates within 0.1 - 40%/s in ascending order. In the present study, thoracic aortas were tested over a 10 - 200%/s range of deformation rates. In both, uniaxial and biaxial tests, the peak force on the arterial samples was around 20% smaller at a deformation rate of 200%/s than at 10%/s (Figures 3 and 4). These results were observed on the average of the five steady-state loading cycles, after five-preconditioning cycles for each deformation rate. However, similar results were also observed on the first loading cycle at each deformation rate, eliminating the possibility of this behavior being due to pre-conditioning. The experimental procedure was validated by repeating the same experiments on latex and nitrile samples. As expected, the effect of deformation rate was opposite in these rubbers to what was observed in the arterial samples. An experiment without any sample was performed to measure the combined effects of inertia, friction, and water drag force, which were found to be negligible, i.e. less than 10 % of the effect of deformation rate.

Sample variability between circumferential and axial force vs. stretch ratio curves was found to be lower than the effect of deformation rate, supporting the use of an isotropic constitutive equation to predict the behavior of the arterial wall. Isotropic models have also been used to model the artery in other studies [43,47-50]. Forces instead of stresses were reported here due to the non-homogeneous stress and strain distributions found in speci- mens subjected to biaxial tests, where the highest and lowest stresses are found near the curved boundaries and near the center of the sample, respectively [51]. Moreover, researchers who are focused (like us) on experimental rather than theoretical work will appreciate to report the data as raw as possible.

The main constituents of arterial tissues are collagen and elastin. The elastic response of arteries is largely due to elastin because elastin is about 1000 times more extensible than collagen. The proportion of elastin to the other arterial constituents is the highest in the aorta. Elastin chains are crosslinked together to form rubberlike, elastic fibers [52]. Each elastin molecule uncoils into a more extended conformation when the fiber is stretched and recoils spontaneously as soon as the stretching force is relaxed. Its elasticity is based on changes in the entropy of the molecular chains, while the material is deformed. An imposed strain increases the order in the molecular network and thus decreases its entropy. According to the thermodynamic laws, the network would try to recover its former shape, increasing their entropy.

One possible explanation for the observed phenomenon is that when elastin is stretched at high rates it would attain a highly oriented conformation and as a result the cross-links would not be able to bear the load due to slippage. However, as the stretch occurs at lower rates, the elastin molecules would have more time ava- ilable to adjust in order to prevent slippage of the cross-links and as a result would be able to bear a higher load. This effect can be viewed as an energy dissipation effect where at high rates of stretching the energy is dissipated rather by the associative flow and the slippage between cross-links. At lower rates the energy dissipation is less. This effect is similar to the stretching behaviour of branched polymers when compared to that of their linear counterparts [53,54]. Here elastin at high extensional rates would behave more like a linear polymer where the linear polymers can slip one past the other to exhibit a reduced load mainly due to high oriented conformation imposed by the high stretching rate. On the other hand, elastin at low extensional rates would behave like a branched polymer where the presence of branches prevent cross-linking slippage and as a result bear higher loads.

Moreover, Trepat *et al.* [55] subjected human airway smooth muscle cells to a transient stretch-unstretch maneuver with zero residual macroscale strain, observing that the cell promptly fluidizes and then slowly resolidifies. Therefore, is it possible that changes in the alignment and configuration of adjacent fibers within the extracellular matrix occurs as soon as one fiber begins to shear over the other (i.e. fluid-like behavior), making softer the overall response in the tissue when stretching the arterial sample at the highest deformation rate (minimum relaxation in the tissue due to short testing time). At lower deformation rates, the fibers have more time to move back to their original configuration, thus restraining their ability to flow (stiffer response).

The constitutive equation selected here is capable of representing the nonlinear elastic behavior of the artery, including the effect of deformation rate, with only four parameters. It produces a unique curve fitting solution to experimental results for each deformation rate. More importantly, it was observed that Mooney Rivlin parameters a_{10} and a_{01} were not a function of the deformation rate. This could imply that they are related to the elastic behavior of the artery. A linear relationship was obtained between parameter a_{30} and the deformation rate, where the partially elastic contribution is given by coefficient \hat{B} .

As shown in **Figure 9**, the model better fits the biaxial data than the uniaxial data. Using higher weights for the uniaxial least square differences, it was possible to obtain a better fit for uniaxial data, at the expense of a deterioration of the fit for biaxial data. The use of material parameters a_{11} and a_{20} did not improve the fits significantly enough to justify increasing the complexity of

the model with these two extra parameters. The use of anisotropic constitutive equations in future studies might reduce this discrepancy between predicted and experimental forces in uniaxial direction.

Clinical results have shown that slow balloon inflation might reduce restenosis. It was hypothesized that the slower deformation rate gives as a result lower intramural stresses and lower arterial injury. This hypothesis is contradictory to our results, which would indicate that a lower inflation rate would result in higher intramural stresses. Other mechanisms might play a role in leading to lower restenosis rates, such as endothelium denudation or artery injury being dependent on stretch ratio rather than on intramural stress.

5. CONCLUSIONS

In the present study, the effect of deformation rate on the mechanical behavior of arteries in vitro under uniaxial and biaxial extensions was investigated. It was found that the loading force at a stretch ratio of 1.5 is reduced by 20 % when the deformation rate is increased from 10 to 200 %/s, implying that the stiffness of arteries decreases with deformation rate. This behavior might be a consequence of the faster fluidization and small resolidification that occurs in the cell at higher deformation rates. This effect of deformation rate was observed to be higher than the effect of anisotropy. The development of an isotropic hyperelastic rate-dependent constitutive model, derived from the Mooney-Rivlin hyperelastic model, is capable of representing this behavior. In its proposed form, the model has only 4 parameters, only one of which varies with deformation rate. The small number of material parameters and simple formulation increase the applicability of this model to numerical simulations. The inverse relationship between stiffness and deformation rate raises doubts on the hypothesized relationship between intramural stress, arterial injury, and restenosis.

Future work might be helpful in order to account the entire loading history of the arterial wall by performing additional characterizations such as relaxation time experiments, and with the aim of a viscoelastic model. The ratio between maximum stretch with balloon and inflation time might be computed as well to observe if the same phenomenon appears.

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