

The Implications of the Sun's Dragging Effect on Gravitational Experiments

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Abstract

Experimental determinations of Newton's gravitational constant, Big G, have increased, in number and precision, during the last 30 years. There is, however, a persistent discrepancy between various authors. After examining some literature proposing that the differences in Big G might be a function of the length of the day along the years, this paper proposes an alternative hypothesis in which the periodicity of said variation is a function of the relative periodicity of the Sun-Earth distance. The hypothesis introduced here becomes a direct application of the Kerr Metric that describes a massive rotating star. The Kerr solution for the equations of the General Theory of Relativity of Albert Einstein fits well with this relative periodicity and adequately predicts the arrangement of the experimental G values reported by sixteen different laboratories. Also, the author explains how the Sun disturbs gravity on the surface of the Earth.

Keywords

Spinning Black Holes, Kerr's Metric, General Relativity, Torsion Measurements, Gravitational Constant G, Precession

1. Introduction

One of the oldest constants in Physics, the constant G (commonly called Big G) that appears in Newton's Law of Universal Gravitation, is the least precise value reported in official tables. In academia, this situation is justified by stating that the gravitational force is weak. The previous argument, however, is not justified if we consider the very precise experiments that had determined G in the last 30 years. Those results, however, show bar errors depicting a gap in error between scholars thus indicating that one, or more, systematic errors might be present. In 2015, Anderson, *et al.* [1] proposed a possible correlation between the observed

experimental variations of G and the length of the day, which is known to have a period of 5.9 years as shown by the 2013 work of Holme and Viron [2]. However, in 2015, M. Pitkin [3] applying a Bayesian model, concluded that the previous period of around six years is not admissible when a complete table reported by Schlamminger *et al.* [4], is included in the analysis. Pitkin ran a function with an age ranging from two to eight years (logically covering the six-year periodicity) and showed no strong correlation between data and hypothesis. This letter introduces a simple model to solve this apparent variability in the experimental determination of the universal gravitational constant.

The author assumes that gravitational interactions do not obey the superposition principle. That is, a third body can perturb the interaction between two bodies (or parts of an extended body). This idea is not new in physics. Hawking and Rayner express in [5] that per the Mach's principle, the "Local physical laws are determined by the large-scale structure of the universe". In addition, Pugach [6] reports extensive experimental work indicating that torsion balances are disturbed by the Sun's activity. The work reported in this paper will be limited to create the connection between a well-established theory and the experimental data available. Also, the author will make an evaluation of the possible perturbative parameters that disturbs the determination of Big G.

2. Quantitative Characterization

2.1. Model

Let's separate the experimental value of gravitational constant $G[G(\exp)]$ into the universal constant Big G and a perturbative effect $\delta G(r)$ that is a function of the distance to the perturbative source. In other words, $G(\exp) = G + \delta G(r)$. Considering that the Sun produces the next biggest gravitational field on Earth (only superseded by the Earth's field) we can redefine the perturbation parameter $\delta G(r)$ as expressed in Equation (1)

$$G_{\rm exp} = G + \left(\frac{a}{r_s}\right)^n {\rm d}G , \qquad (1)$$

where *a* is the average separation between the Earth and the Sun, r_s is the instantaneous Earth-Sun distance, and d*G* is a gravitational range that must be founded. A combination of theory and experimental data will help to set the number *n* that is expected to be an integer power (with more probable values of 1, 2, 3 or 4).

The ratio a/r_s was calculated using the monthly Earth-Sun distance reported by the web page [7]. A table with the monthly Earth-Sun distance along a 19-year interval was prepared to fit in as a function of time. Equation (2) describes that variation where the variable *t* represents the time in years,

$$\frac{a}{r_s} = \frac{1}{1 + e \cos\left(\frac{2\pi}{1.0}t - 3.2\right)}.$$
 (2)

Equation (2) was made using astronomical data, and consequently, its parameters are independent of our choice. Parameters in Equation (2) are a unit base of 1, an amplitude of 0.017 that is the Earth's orbital eccentricity *e*, a period of 1.0 year, and a negative initial phase of 3.2 rad.

The total perturbative influence of the planets and the Moon in G(exp) are too low in comparison with the contribution of the Sun. So, they were ignored in the calculations. However, to get seven significant figures in Big G, it will be necessary to include at least Jupiter.

2.2. Direct Observations of G

Over the last 30 years, a good number of experimental points with very precise results defining Big G's value with six significant figures have been accumulated. **Figure 1** combines the experimental points of G(exp) from references [8]-[22] and our function (1) using n = 2. (The yellow spot representing the Sun was offset a little intentionally). The only two degrees of freedom that we have are the base value of the universal constant Big G and the value of the perturbative parameter dG to minimize the difference between our function and the experimental data as much as possible. A minimum quadratic regression was used to better fit the experimental points up to six significant figures. The values of the universal constant G and the perturbative amplitude dG for this purpose is $6.65471 \times 10^{-11} \text{ Nm}^2 \cdot \text{kg}^{-2}$ and $0.019372 \times 10^{-11} \text{ Nm}^2 \cdot \text{kg}^{-2}$. The last two values come from statistical data that is affected for more than one unnoticed effect; that means that the real values could be a little offset from those values that come from these experiments. This paper has the main intention to indicate the way to clean up that statistical error. Figure 1 shows the trajectory of the Earth (lower ellipse), the experimental points of Big G (closed circles), and our fitting model (upper ellipse). The biggest experimental G values occur around the perihelion of the orbit and the lower values occur around the aphelion.

2.3. Influence of Big G on Terrestrial Gravity

An additional correlation between the author's model and other parallel gravitational measurements connected indirectly with Big G is available. For example, the work done by Nicolas *et al.* [23] in Grasse, France shows the correlation. **Figure 2** has three parts that include: Nicolas *et al.* experimental points of absolute gravity, the paper model, and the experimental points of Big G [8]-[22] from 1998 to 2004.

The agreement between the experimental points and the two models is good enough to increase the credibility of the hypothesis introduced in this paper. (It is important to clarify that their model is based on the measurement of little g being dependent on Big G and not the other way around). In this graphical comparison, amplitudes have a different meaning. The matching of maxima and minima phase is another point in favor of the idea being introduced here.



Figure 1. Plots of our Equation (1) and the experimental points from [8]-[22]. Points with same color come from same lab in different attempts. The *x*-axis and *y*-axis are in units of the Astronomical Unit *a* and the *z*-axis shows the values of Big G times ten the eleventh power in units of Nm²·kg⁻².



Figure 2. Plots of our Equation (1) (yellow) and some G(exp) points from 1998 to 2004 (points with horizontal error bars). The points with only vertical error bar (red diamonds) come from micro-variations of little g in Grasse, France.

3. Experimental Approach

The possible variation of the gravitational interaction between any two masses with the distance to the Sun is something relatively easy to corroborate in comparison with the mathematical evaluation of that interaction. Besides that, the variation will not be easy to observe because those values could be many times smaller than the ambient noise if the apparatus is not carefully prepared.

The author's setting included three variations to the Cavendish's balance [24]. The main one is the introduction of two strings versus the one used by Cavendish. The nine-millimeter separation of the two strings works like a mechanical filter that reduces the effect of weak perturbations with low duration of time. The second variation was the introduction of a microscope because the microscopic displacement of the balance's iteration is stable enough to be observed. That happens because the change of the gravitational influence of the Sun is large enough in time. The microscopic motion output is digitally recorded on every trial. The observation was conducted every day approximately before sunrise to give the balance enough time to clean out the noise coming from the human activity of the previous day.

The third improvement on the balance was the use of four masses in the form of a cone. That is the most effective form to produce the greatest gravitational interactions. The spherical masses are not the most effective, as is assumed by many scholars, because the mass of a body increases with radius but the distance between centers of masses decreases when the radius is squared. On a cone, if the cone's vertex is keeping constant, the center of mass does not move if the size of the base is increased.

The experimental balance has a permanent state of vibration. The technique for this kind of situation is to accumulate enough data to find the average position and its standard deviation. In this case, that method is not efficient. A better method is to record the two points that are farthest from the equilibrium point daily and represent both points in the graph. Of course, the size of the error bar of every point will fit inside the two points that were recorded. **Figure 3** has the observations from December the 2nd to July the 4th. The negative slope on the graph is matching well with the results reported over the last 30 years.

A bonus from **Figure 3** is that the curve shows the superposition of two permanent signals with an average repetition of around ten and six days respectively. Both periodicities can be explained if the nucleus of the Sun is a plasma torus that rotates with an angular momentum of value *J*. We can associate the ten days with a plasma torus precession of 20 days' period and the six days with its nutation. This idea is in correspondence with the visual information that the Sun's equator is rotating with a period of approximately 25 days.

The General Theory of Relativity (GTR) of Albert Einstein [25] is the necessary tool to understand this problem. The solution of the GTR found by Kerr [26], for a non-charged rotating body, discloses the consequences that can be used to explain all the problems analyzed in this paper. Kerr's Metric has the following solution,

$$c^{2} d\tau^{2} = \left(1 - \frac{\frac{2GM}{c^{2}}r}{r^{2} + \frac{J^{2}}{M^{2}c^{2}}\cos^{2}\theta}\right)c^{2} dt^{2} - A(r, M, J, \theta) dr^{2}$$
(3)
$$-B(r, M, J, \theta) d\theta^{2} - C(r, M, J, \theta) d\phi^{2} + D(r, M, J, \theta) c dt d\phi,$$

where the mass of the rotating body is M, and J is its angular momentum. The functions A, B, C, and D were introduced in (3) because these functions are not directly related to this paper. A consequence of the Kerr's Metric is a dragging effect on the spin of any gyroscope localized inside the gravity field of the rotating mass. It is interesting that the dragging effect cannot change the spin's value of the gyroscope (even if the spin is zero) but generates the spin precession. The author observed in his non-rotating horizontal bar the precession effect stated by Kerr's Metric as a variation of the azimuthal angle φ (see Figure 3) and the vertical angle θ (originally around 90 degrees; see Figure 4. It can be checked that Figure 3 and Figure 4 peaks coincide with each other in time.

The qualitative correspondence between theory and observation moves the author to do a quantitative analysis. Per the Kerr's Metric, a gyroscope spinning inside the field of a rotating mass has a precession at the frequency [27],

$$\boldsymbol{\Omega}_{p} = \frac{G}{c^{2} \rho^{5}} \Big[3 \big(\boldsymbol{J} \cdot \boldsymbol{\rho} \big) \boldsymbol{\rho} - \rho^{2} \boldsymbol{J} \Big], \qquad (4)$$

where *G* is the gravitational constant, J is the total angular moment of the rotation mass, *c* is the speed of light, and ρ is the distance to the dragging source. Starting from Equation (4) and the condition that the plane of rotation of the nucleus of the Sun almost coincides with the plane of the solar system allows us to calculate the Sun's plasma torus velocity with the necessary relativistic equation,

$$\sigma^{3}\Omega_{p}c\left(1-\beta^{2}\right)^{3/2}-GM\left(1+\frac{\alpha}{4}\right)R\beta=0,$$
(5)

where *R* is the Sun's plasma torus major radius, and *a* is the ratio of minor radius to the major radius (the value *a* comes close to 0.01 from the relativistic condition that the velocity of the plasma torus's faster points must be smaller than the speed of light), *M* is the mass of the Sun, and ρ is the Earth-Sun distance. The precession period of the Earth-axis of rotation is 26,000 years, and the core radius of the Sun is close to 280 million meters. Those numbers indicate that the speed of the Sun's plasma torus is around 0.986216c.

The Sun's static gravitational influence makes the Earth revolve around the Sun, and does not affect gravity on the Earth's surface. However, the dragging gravitational effect of the Sun's rotating nucleus is working here on Earth. We use the temporal part of Kerr's solution to calculate an equivalent gravity coming from the dragging effect. The resulting equation is

$$g_{\text{dragging}} \approx \frac{GMR^2}{r^4} \left(\frac{\beta}{1-\beta^2}\right)^2 \cos^2\theta \,.$$
 (6)



Figure 3. Azimuthal angle of the modified Cavendish's balance on the *y*-axis. According to the Newton's gravitational theory that angle should be constant and should not have any oscillating pattern if the distance between the attracting masses is kept constant.



Figure 4. Variation of the microscopic angle between the horizontal plane and the vertical direction. According to Newton, that angle should stay constant at 90 degrees. The quadratic increment of the angle in the initial 60 days is associated with the increment in length of the strings that keep in place the balance. In other words, there is always an initial string elongating effect on this experimental setup.

The author, correlating its experimental points with the result from [8]-[22], deduces that the Sun's plasma torus plane of rotation makes an angle of 80 degrees with the zenith of the solar system (nutation cover only 2 degrees). Applying that angle in Equation (6), the gravity on Earth should oscillate annually between 100×10^{-8} m/s² and 92×10^{-8} m/s². The difference between those numbers $(8 \times 10^{-8} \text{ m/s}^2)$ is close to the 9×10^{-8} m/s² annual variation of *g* in France, reported by Nicolas *et al.* [23]. By now, the reader will know that the power of the inverse of the distance in our initial hypothesis becomes n = 3 per Equation (4). That inverse power is the same that appears in the tidal equation and that maybe is not a numerical coincidence. The analysis of that coincidence will also be discussed in a following paper.

Figure 1 only has 20 points affected without control of the Sun's dragging effect. On the other hand, the author report here 185 well characterized points taken in the first six months of 2017. An increment in accuracy could be done by mixing both data. The axis of **Figure 4** were converted with the functions,

$$x = 0.985 + 0.030/182 (N_{day} - 3)$$
 and $G = 6.6695 + 0.0009 \exp(0.0008\varphi)$, (7)

where N_{day} represent the number of the day from January first to July fourth, x is a parameter without units that represent the distance to the Sun (x = r/a), and φ is the balance's microscopic deflection. The author selects the points with minimum deflection. The idea was to minimize as much as possible the influence of the precession and the nutation effects. It is, for those points in Equation (8), the angle θ was close to 90 degrees. Figure 5 showed all the points under this condition.





The Fit function from Wolfram Mathematica 11.1 reported the numbers,

$$G_{\exp} = G + \frac{dG}{x^{3}} = \left[\left(6.6114 \pm 0.0031 \right) + \frac{\left(0.0616 \pm 0.0031 \right) \sqrt{1 + 9\cos^{2}\left(\theta\right)}}{x^{3}} \right] \times 10^{-11} \text{ N} \cdot \text{m}^{2} \cdot \text{kg}^{-2}.$$
(8)

4. Conclusions

The most influential labs in the world always turn toward the values of the fundamental constants in physics. Not just for trustworthiness but also for safe comparison with the reliable values from other labs. The constant's errors may not overlap; however, the scientific community expects those errors to be close to the same significant figure on every constant. That has not been happening to the gravitational constant Big G in the last 30 years. Some differences cover three orders of magnitude, far away from every error. The explanations proposed by now to solve this contradiction have not been established without reasonable doubt.

The author proposed that the experimental values of Big G will depend inversely on some power of the distance to the Sun, and show an experimental work to test the hypothesis. Its experimental confirmation justifies the search for any theory that interconnects all the information available.

The author founded the mathematical explanation that covers all the points on the determination of Big G in the General Theory of Relativity under the assumption that the core of the Sun has the form and condition of a relativistic rotating plasma torus. The Sun's core is massive enough to produce visible dragging effects here on Earth assuming its spin suffers precession and nutation.

The experimental work introduced here clearly shows how others failed to null the dragging effects from the Sun by making, half of the time, the gravitational force to be right-handed and, the other half of the time, left-handed because those effects are not symmetric in space and cannot be neutralized. The remaining difference showed up in every result of Big G and if displayed on an annual graph, will create a pattern as an inverse relationship with the distance to the Sun. Modern gravitational theory dictates that the connection between the experimental value of Big G and the distance to the Sun, follows a negative third power law. One of the questions that arises from this research, "Is the dragging effect changing along the years, forcing in that way the variation of the length of the day"? This question serves as an exigency for the opening of labs to make diary controls of the dragging variation. The relativistic nature of this effect could potentially be dramatically unstable. Furthermore, while one might suggest the duration of the day could be an explanation it could also be the case that the inverse result is true. The author will show experimental evidence on a new paper that this dragging function is responsible for the Sun's tidal effect, meaning they are not correlative but rather one in the same.

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