

Observational Evidences for Extremely Strong Magnetic Fields in Solar Flares

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Abstract

New observational data related to the X1.1/2N solar flare of 17 July 2004 were investigated and compared with some old data for other powerful flares and non-flare regions. Observations were carried out with the Echelle spectrograph of the Kyiv University Astronomical Observatory. The Stokes $I \pm V$ profiles of several metallic lines with different effective Lande factors g_{eff} have been analyzed including the FeI 5434.5 line with very low magnetic sensitivity ($g_{\text{eff}} = -0.014$). The obvious evidences of the emissive Zeeman effect were found as in lines with great and middle Lande factors as in FeI 5434.5 line. On the basis of all analyzed data one can conclude that upper magnetic field limit in flares can reach 70 - 90 kG, *i.e.* about more order higher than the well-known magnetic fields in great sunspots. The possible physical nature of such superstrong fields is discussed.

Keywords: Sun, Solar Flares, Small-Scale Magnetic Fields, Superstrong Fields

1. Introduction

The upper magnetic field strength limit in the solar atmosphere is unknown at present. The point of view predominates that the strongest fields exist at great sunspots. Parker [1] said: "More recently we have come to appreciate that the sunspots is just the most conspicuous feature of a vast word of magnetic activity". According to observations, the field strength is here, as rule, 2100 - 2900 G and sometime 3500 - 4000 G. The record field of 6100 G in a sunspot was measured by Livingston et al [2].

Solar flares are very interesting and violent processes in solar active regions where the strongest fields could exist too. Each flare is a grandiose explosion in a wide range of height in solar atmosphere, with sharp increasing of temperature, gas pressure and ionization of plasma. Hot flare plasma outside magnetic flux tubes should press on walls of these tubes and may increase magnetic strength inside tubes. In addition, increasing of plasma ionization leads to amplification of electric currents and to magnetic field intensification, if structure of magnetic field is force-free.

Unfortunately, magnetic field measurements in flares are not quite simple as in sunspots. First of all, solar flares occur suddenly and develop rapidly. The most in-

teresting phases of a flare can be missed if observations carry out in non-continuous regime. Namely such non-continuous regime is usual for spectral observations, but namely spectral observations are the most suitable for magnetic field measurements in flares. As to instruments working in automatic regime like SOHO/MDI filter magnetograph, ones give mainly longitudinal magnetic field component—no field module. In addition, in case of the most powerful flares practically all spectral lines have more or less intensive flare emission. In some cases, this emission distorts magnetographic signal sufficiently and leads to the false field values, even false measured polarity [3-5].

As to spectral observations, some authors assumed the upper magnetic field strength limit in range of 50 - 100 kG [6,7]. In particular, Severny [6] assumed the 50 kG fields in small-scale (spatially unresolved) magnetic elements of sunspots. He wrote: "Thus, it is quite possible, that magnetic field inside sunspots has a fine structure, which is averaged during the observations and gives the fields of 10^3 gauss, whereas the true field in the separate elements of sunspots (perhaps, of different magnetic polarity and inclination) reaches of several tens of kilo-gauss; *i.e.* situation can be the same as in a case of granulation where the averaging of separate fields of granules (~ 100 gauss) leads to measured field on level

of a few gauss”.

Bruce [7] supposed that long wings of $H\alpha$ line in flares (till 8 \AA) could arise due to the Zeeman effect in 10 - 100 kG fields. Really, this is one of possible interpretation: some account to the line-profile broadening should give also temperature, turbulent velocity and electric fields. For these effects separation, the full Stokes parameters I , Q , U and V of polarized light are needed. In other hand, if magnetic field is subtelescopically tangled (for instance, in form of mixed polarity structures) the Stokes diagnostics could be useless: we can have essential broadening of Stokes I profile but practically zero polarization on different distances from line center.

More direct evidences to such “superstrong” fields in flares were obtained using observations in FeI lines with very low Lande factors [8-11]. It was shown that Stokes $I \pm V$ profiles of some such lines, e.g. FeI $^5F_1^5F_1$ $\lambda = 5123.723 \text{ \AA}$ and FeI $^5F_1^5D_0$ $\lambda = 5434.527 \text{ \AA}$, have in flares narrow and splitted emission peaks in their cores. Both named lines have the following Lande factors: theoretically, for LS coupling, $g_{LS} = 0.000$, but laboratory determined values $g_{Lab} = -0.013$ and -0.014 , respectively [12,13]. Observed splitting of emission peaks in $I \pm V$ profiles was $\Delta\lambda_V = 10 - 36 \text{ m\AA}$, whereas similar splitting in $I \pm Q$ profiles was close to zero ($\Delta\lambda_Q \leq 5 \text{ m\AA}$). Taking into account that non-flare profiles did not display similar effects, it was assumed the Zeeman effect manifestation in field of $B \approx 26 - 94 \text{ kG}$.

Naturally, such extraordinary conclusion needs an additional and careful verification on a base of new data

for other flares. Unfortunately, this is almost insuperable task for spectral observations. Enough bright flare emission in FeI 5123.723 and 5434.527 cores occur in very powerful solar flares only, of X1 class or more. More weak flares, e.g. of M or C importance, produce, as rule, ordinary Fraunhofer line profiles, without emission peaks in their cores. So, only unique spectral data are suitable to the problem, related to the suddenly occurred large and bright solar flares.

Some of such flares were observed on Echelle spectrograph of horizontal solar telescope of Kyiv University astronomical observatory [14]. About five decades from 1950s, this instrument is used for regularly carried out spectral observations in wide range of wavelength, from 3800 to 6600 \AA . Till 1975 these observations were made in Stokes I , but since 1975—in Stokes $I \pm V$ and $I \pm Q$ too. In this work, the new observational data related to the X1.1 solar flare are presented and analyzed.

2. Observational Data and Profiles of Lines

Flare of 17 July 2004 had occurred in active region 10649, at its tail part placed closely to spots of N polarity with magnetic strength of 2800 G and heliocentric distance $\mu = 0.88$. Observations of the flare were carried out with the Echelle spectrograph of the Kyiv University Astronomical Observatory [14]. From 8:01 to 8:16 UT six Echelle Zeeman spectrograms were obtained with a circular polarization analyzer, but at present study the first spectrogram only is analyzed containing the strongest flare emission (**Figure 1**), nearest to time of the X ray peak according to GOES data (7:56 UT). All

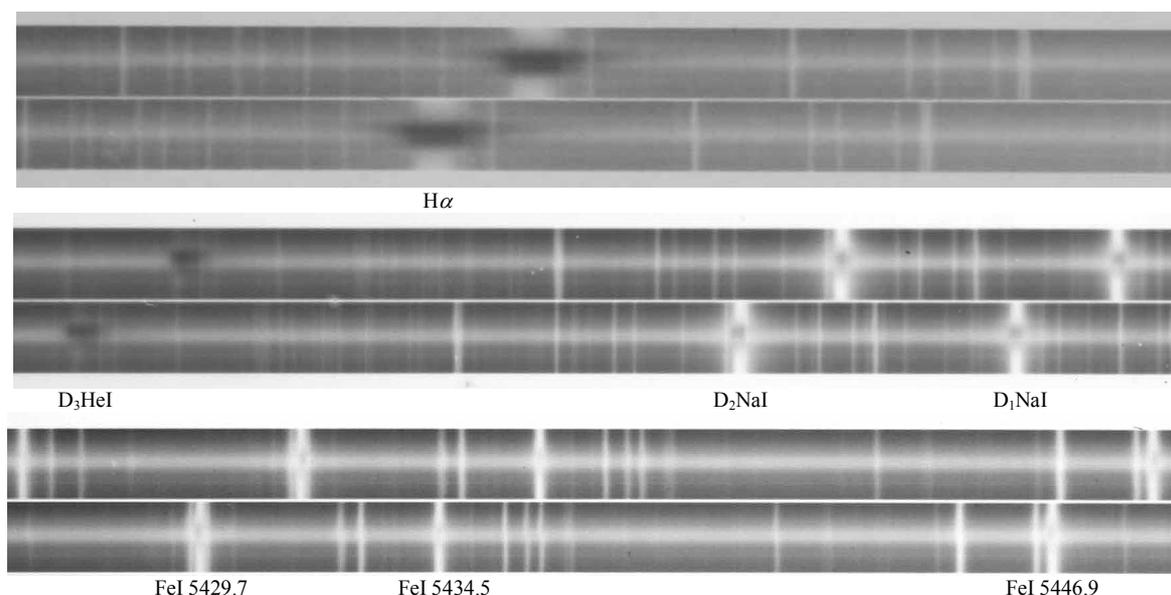


Figure 1. Fragments of the Echelle Zeeman spectrogram ($I + V$ and $I - V$ spectra) showing the flare emission (for 8:01 UT) in some spectral lines.

Table 1. List of selected spectral lines.

Number	Wavelength (Å)	Element and multiplet number	Equivalent width W (mÅ) Rowland table	Excitation potential (eV)	Effective Lande factor g_{eff}
1	5397.141	FeI-15	239	0.91	1.426
2	5434.534	FeI-15	184	1.01	-0.014
3	6301.508	FeI-816	127	3.65	1.669
4	6302.499	FeI-816	83	3.69	2.487

spectra were made with ORWO WP3 photoemulsion.

Four spectral lines of FeI were used for magnetic field measurements in area of flare—two lines of 15th multiplet and two lines of multiplet No. 816 (**Table 1**). Lines Nos. 3 and 4 form at level of middle photosphere whereas the cores of lines Nos. 1 and 2—in range of upper photosphere and temperature minimum zone. Effective Lande factors g_{eff} listed in the table correspond to laboratory determined values from papers [12,13]. As it was pointed out above, theoretical effective Lande factor for 5434.534 line is $g_{\text{eff}} = 0.000$ for case of LS coupling. At first approximation, this line can be considered as Doppler one, whereas other three lines—as Zeeman sensitive lines. Only in case of very strong fields of $\sim 10^4$ G range, FeI 5434.5 should be considered as Zeeman line too [8].

Photometrical study with microphotometer MF-4 shown that lines Nos. 3 and 4 have ordinary Fraunhofer profiles in the flare, without any emission peaks at their cores (**Figure 2**). On this figure, $\Delta\lambda$ is distance (in mÅ) from center of telluric line $\text{O}_2 \lambda = 6302.000 \text{ \AA}$. Centers of lines FeI 6301.5 and FeI 6302.5 correspond to $\Delta\lambda$ about -500 and $+500$ mÅ, respectively; second telluric line O_2 has $\Delta\lambda \approx 750$ mÅ. Three photometrical sections Nos. -19,

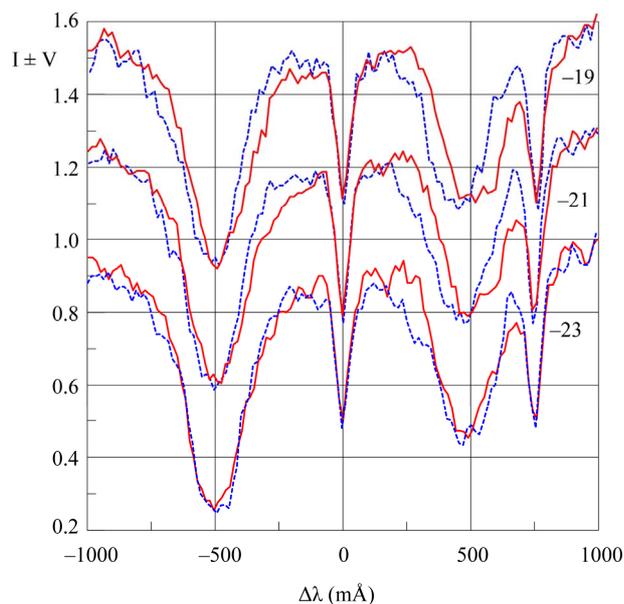


Figure 2. Observed Stokes $I \pm V$ profiles for the FeI 6301.5 and 6302.5 lines in the flare (see text).

–21 and –23 related to strongest flare emission are presented here. Space interval between nearest sections (e.g. –19 and –21) corresponds to 1 Mm on the Sun. Distribution of intensity for section No.–23 is original, whereas other distributions are shifted along ordinat axis on 0.3 and 0.6 in unit of continuum level.

One can see the lack of full Zeeman splitting in FeI 6302.5 line with the highest magnetic sensitivity ($g_{\text{eff}} = 2.487$) that indicates the moderate longitudinal magnetic field—of subkilogauss range (< 1 kG). The FeI 5397.1 line has more complicated and interesting profiles—with Fraunhofer wings and emission peaks in line core (**Figure 3**).

Second interesting effect: splitting of emission peaks appreciably exceeds the splitting of line wings. Similar effect for flares was found earlier by Lozitska and Lozitsky [3]. Notice, emission peaks in lines like FeI 5397.1 form at upper photosphere whereas Fraunhofer wings—below named level [15]. Thus, this second effect is evidence to the local magnetic field amplification at some level of the flare.

3. Measured Magnetic Fields

Magnetic fields in area of the flare were measured using

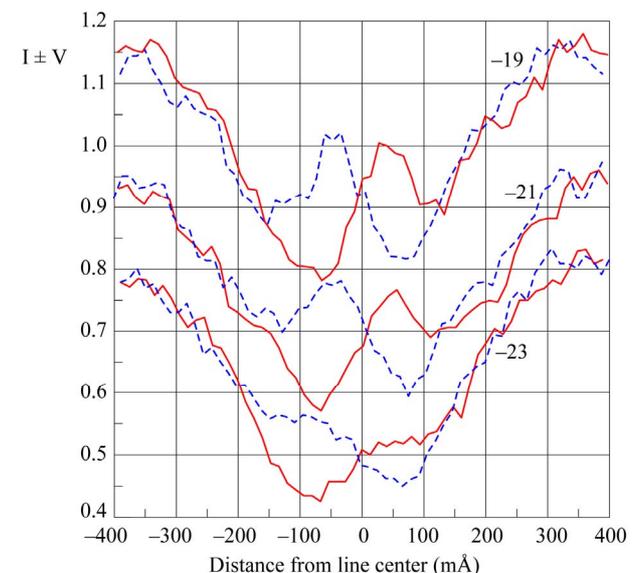


Figure 3. Observed Stokes $I \pm V$ profiles for the FeI 5397.1 line in the flare.

two methods: a) by splitting of “center of gravity” of whole $I \pm V$ profiles and b) by relative shifts of limited intervals of these profiles. In first method, relative shift of the ‘centers of gravity’ of the $I + V$ and $I - V$ profiles is interpreted as a result of the Zeeman Effect and calibrated in values of magnetic field using well known formula

$$\Delta\lambda_H = 4.67 \times 10^{-13} g \lambda^2 B \quad (1)$$

where the Zeeman splitting $\Delta\lambda_H$ and wavelength λ are in Å, and field strength B in gauss (G). In a case of homogeneous field and small Zeeman splitting (exactly, if $\Delta\lambda_H/\Delta\lambda_D \ll 1$, where $\Delta\lambda_D$ is the Doppler half width of line), the longitudinal magnetic component B_{\parallel} can be obtained from such measurements. If magnetic field is inhomogeneous, such observations give likely the longitudinal magnetic flux [16].

Measurements with second method were calibrated analogously, using formula (1) too. For the same spectral line, these data reflect the magnetic field strength in narrower interval of height in solar atmosphere than ‘center of gravity’ data. It is useful to remember, that second method is similar methodically to classical magnetographic measurements [17] with fixed exit slits in line profile.

First method was applied to FeI 6301.5 and 6302.6 lines, and second—to FeI 5397.1 (Table 2). In this table, measurements in FeI 5397.1 for $|\Delta\lambda| = 200 - 300$ mÅ correspond to Fraunhofer wings of line, and for $|\Delta\lambda| = 0 - 200$ mÅ—to splitted emission peaks in line core. The error bars are about ± 40 G for FeI 6302.5, ± 60 G for FeI 6301.5 and ± 75 G for 5397.1. One can see that all three magnetosensitive lines (excluding FeI 5397.1 for $|\Delta\lambda| = 0 - 200$ mÅ) give similar results. In photometrical section No.-23 (which places closely to magnetic inversion line $B_{\parallel} \approx 0$) all lines give practically coincided results, whereas in sections Nos. -19 and -21 the measured field strengths are some differ. At these sections, the following inequality is obvious: $B_{5397.1} < B_{6301.5} < B_{6302.5}$. Taking into account that equivalent widths of lines W correlate as: $W_{5397.1} > W_{6301.5} > W_{6302.5}$, we can admit the negative vertical magnetic field gradient ($\partial B/\partial h < 0$) at photospheric layers of the flare.

Likely, such situation (*i.e.* $\partial B/\partial h < 0$) was in a limited range of photosphere. As it follows from splitting of

Table 2. Comparison of the magnetic fields measured in different lines and places of the flare.

Line	$B(-19)$, G	$B(-21)$, G	$B(-23)$, G
FeI 5397.131 $ \Delta\lambda = 200 - 300$ mÅ	520	340	170
FeI 5397.131 $ \Delta\lambda = 0 - 200$ mÅ	2070	2630	2860
FeI 6301.508	600	420	150
FeI 6302.499	720	550	80

emission peaks of FeI 5397.1 line ($\Delta\lambda = 0 - 200$ mÅ), the kilogauss fields ($B \approx 2.1 - 2.9$ kG) had existed in upper photosphere. Thus, we can expect that at least two layers were in area of flare: a) $\partial B/\partial h < 0$ in range of the middle photosphere, and b) $\partial B/\partial h > 0$ at upper level. Perhaps this indicates non-monotonous magnetic field distribution in area of the flare similar to a case described earlier by Lozitsky *et al.* [15] for another flare.

4. Polarized Effects in Core of FeI 5434.5 Line

Observed $I \pm V$ profiles of FeI 5434.5 line in flare of 17 July 2004 are presented on Figure 4. Similar to cases of FeI 6301.5, 6302.6 and FeI 5397.1 lines (see above Figures 2 and 3), distribution of intensity for section No.-23 is original, whereas other distributions are shifted along ordinat axis on 0.3 and 0.6 in unit of continuum level.

One can see the following:

- 1) Fraunhofer wings ($|\Delta\lambda| > 100$ mÅ) and narrow emission peaks in line core ($|\Delta\lambda| < 100$ mÅ); observed half-widths of these peaks are very small, about 50 mÅ;
- 2) Fraunhofer wings of $I + V$ and $I - V$ profiles are, in general, in good agreement—differences between these profiles are, as rule, less than 1%;
- 3) Differences between $I + V$ and $I - V$ profiles in line core are larger and reach 3% - 4%; this means presence of weak circular polarization here;
- 4) Locations of emission peaks in $I + V$ and $I - V$ profiles coincide, in general, in photometrical section No. -19, but ones are some different in sections Nos. -21 and

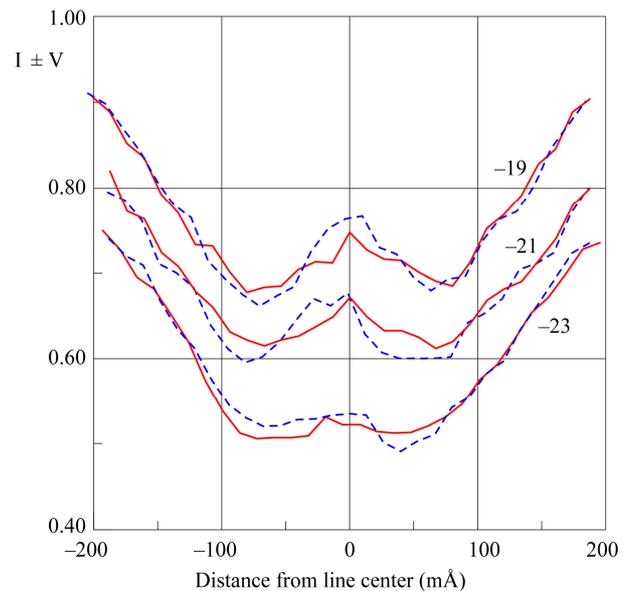


Figure 4. Observed Stokes $I \pm V$ profiles of FeI 5434.5 line in flare of 17 July 2004.

–23. At two last sections, pictures of $I \pm V$ profiles at their central parts looks as spectral splitting of tops of emission peaks. Value of this splitting reaches about 30 mÅ.

For estimation of noise level and possible error bars, similar photometrical $I \pm V$ profiles are needed observed in non-flare area too. Another way is suitable to this task: study of similar profiles of any strong telluric line; such line, naturally, should be considered as pure non-magnetic one.

For this purpose, FeI 5434.5 line was studied outside the flare—in area of weak magnetic field ($B_{\parallel} < 100$ G). In addition, non-magnetic effects were estimated by telluric H_2O $\lambda = 5919.644$ Å line observed on very large air mass, $M = 21.2$ (Figure 5).

From Figure 5 follows that observed line profiles outside flare are smooth, with small (less than 1% - 2%) differences between $I + V$ and $I - V$ profiles. From this point of view, observed effects in the flare reaching of 3% - 4% (Figure 3) likely are real.

5. FeI 5434.5 Line in Other Flares (Old Data)

Firstly similar features in core of FeI 5434.5 like Zeeman Effect were discovered in 2B-flare of 16 June 1989 [8]. Later, such spectral manifestations were found also in FeI 5123.7 line with effective Lande factor $g_{\text{Lab}} = -0.013$ [9]. Lozitsky and Staude [10] observed weak splitting of bisectors of $I + V$ and $I - V$ profiles in core of FeI 5576.1 line ($g_{\text{Lab}} = -0.012$) in solar flare of 5 November 2004; if this splitting to interpret as Zeeman effect, corresponding magnetic field is 10 - 15 kG. So, spectral features like

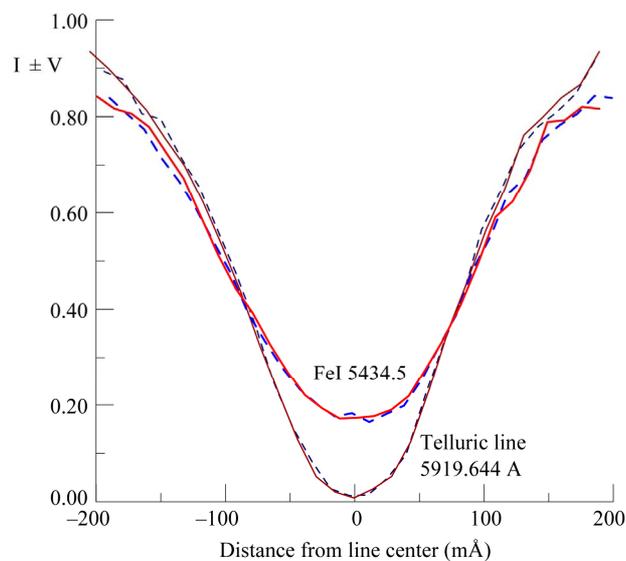


Figure 5. Observed Stokes $I \pm V$ profiles of FeI 5434.5 line outside flare and telluric line H_2O $\lambda = 5919.644$ Å (see text).

Zeeman effect were found in all three well known “non-magnetic” lines—FeI 5123.7, 5434.5 and 5576.1.

An example of such manifestations is given on Figure 6 for flare of 16 June 1989 [11]. One can see well visible splitting of emission peaks in core of FeI 5434.5 line; its value is close to 30 mÅ. There is characteristic splitting also bisectors of $I + V$ and $I - V$ profiles (thin lines on Figure 6): they are practically unsplit in Fraunhofer wings ($|\Delta\lambda| > 100$ mÅ), but essentially splitted (≈ 12 mÅ)—in line core.

Calculations showed [9] that observed $I \pm V$ profiles like presented here on Figure 6 can occur owing to multi-component magnetic field structure (Figure 7). Namely, it is possible to present the observed profile as a sum of Fraunhofer profile 1, unpolarized flare emission 2, first polarized emission 3 and second polarized emission 4; summary profile is presented here as 5. Half-width of unpolarized emission is about 55% relatively Fraunhofer profile 1, flare emission 3 – $\approx 40\%$ and flare emission 4% - 20% only. Magnetic fields was assumed zero for unpolarized emission, 3.5 kG and 65 kG for first and second polarized emission, respectively.

So, these calculations can reflect the following interesting effect: half-width of line profiles are the less, the more is magnetic field strength in subtelescopic fine structures. Earlier narrow spectral lines ($\approx 30\%$ - 50% in unit of undisturbed profile) in subtelescopic magnetic elements with strong field were assumed by Lozitsky [18] and also Lozitsky and Tsap [19] on a base of non-flare magnetic field observations.

Naturally, calculated line profiles presented on Figure

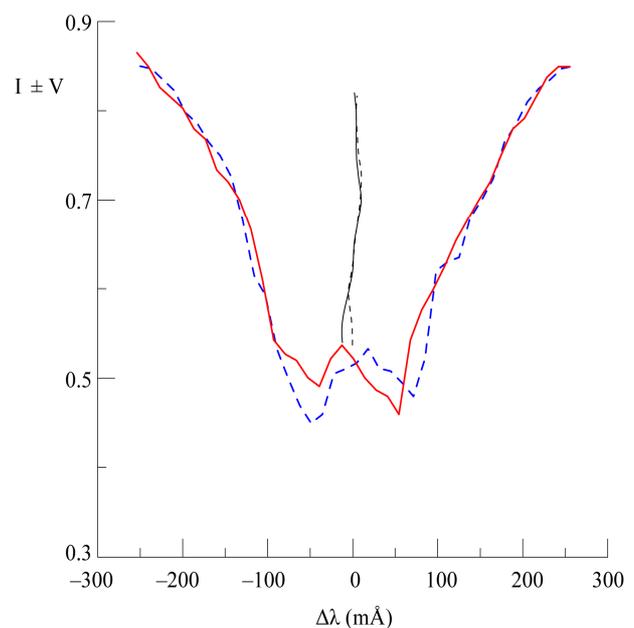


Figure 6. Observed Stokes $I \pm V$ profiles of FeI 5434.5 in solar flare of 16 June 1989 (Lozitsky, 2009).

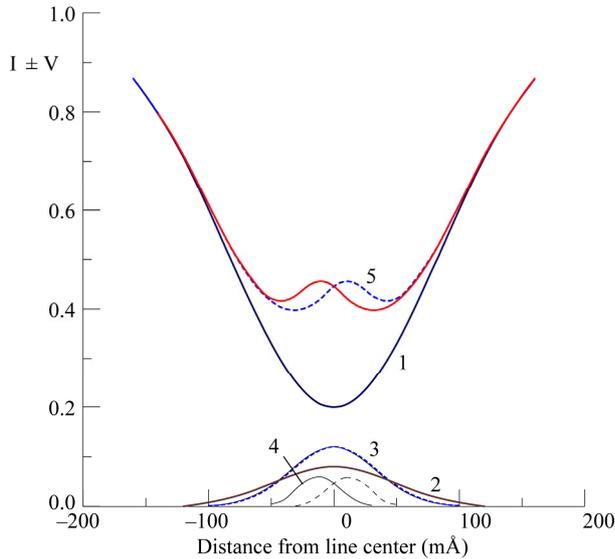


Figure 7. Calculated $I \pm V$ profiles of FeI 5434.5 line (see text).

7 are very idealized—they do not present all fine details of observed profiles (Figure 6). Some of these details could have instrumental nature, but other—solar one. More detailed study shown that three types of the Zeeman-like perturbation occur in core of FeI 5434.5 line in area of solar flares [9]:

1) *NE* (“normal Zeeman”) type has appreciable relative shift of emission peaks in line core, but good agreement of $I \pm V$ profiles in line wings; namely such effect is presented on Figure 6 in line core, $|\Delta\lambda| < 100$ mÅ ;

2) *CE* (“crossover-emission”) type has one-peak emission in $I + V$ profile (or $I - V$), but two-peak emission in $I - V$ (or $I + V$); one can see similar picture on Figure 4 for photometrical sections No. –21. In the majority of cases, one-peak emission places in the middle position between two-peak one;

3) *A* (“absorption”) type is observed in Fraunhofer profiles at their central parts; such effect one can see from bisectors splitting on Figure 6. Similar features were observed also in pure Fraunhofer profiles outside flares.

In favour of Zeeman nature of named effects indicate also observations of $I \pm Q$ profiles of FeI 5434.5 line [9]. It was shown that profiles $I + Q$ well coincide with $I - Q$ (in the same positions of flares) and did not display any trustworthy (more than 3 mÅ) splitting of emission peaks in line core.

Preliminary statistics of *NE* effect is given on Figure 8 [9]. On this histogram, $\Delta\lambda_B^*$ is splitting of emission peaks in line core relatively line wings. It was assumed that $\Delta\lambda_B^* > 0$, if sign of emission peaks is the same as for other magnetosensitive lines with $g_{\text{eff}} > 0$, and $\Delta\lambda_B^* < 0$ in the opposite case. One can see that distribution for

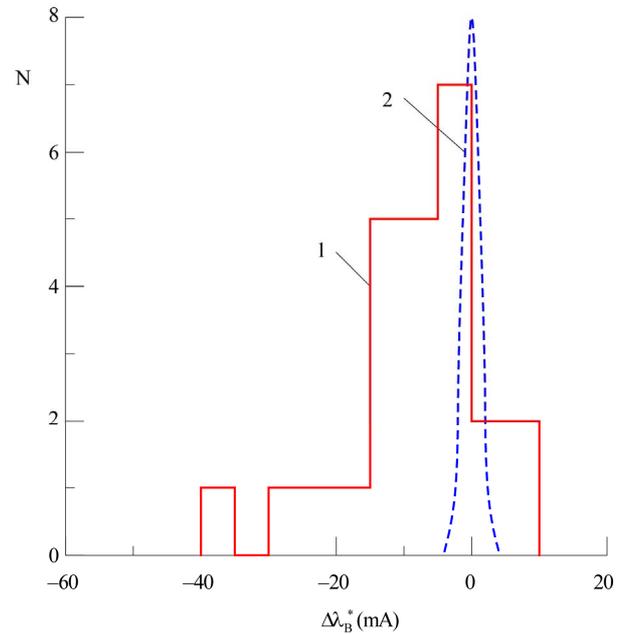


Figure 8. Distribution of splitting $\Delta\lambda_B^*$ in core of FeI 5434.5 line observed in flares (1) and non-flare region (2) according to [9]; N —number of cases.

flares (1) is much wider than for a non-flare region (2). In first approximation, distribution (2) can be considered as instrumental, caused by noise fluctuations on spectrograms. Majority of these fluctuations are less than 3 mÅ, whereas the values of $\Delta\lambda_B^*$ for flares (1) are much wide—in range from –40 mÅ to +10 mÅ.

Second important peculiarity: cases of $\Delta\lambda_B^* < 0$ for flares (1) predominate over $\Delta\lambda_B^* > 0$. This means that our observations confirm the negative sign of effective Lande factor of FeI 5434.5 line (remember, it is –0.014). But, naturally, such “superfine” effect could be discovered only in case if “superstrong” magnetic fields ($\sim 10^4$ G) produced the line splitting.

In fact, from formula (1) follows that magnetic field strength B for FeI 5434.5 line can be calculated via

$$B = 5.18 \times 10^6 \Delta\lambda_H \quad (2)$$

where B is in Gauss (G), and $\Delta\lambda_H$ in Å. If we assume $\Delta\lambda_B^* = 2\Delta\lambda_H$ (really, this is true for longitudinal magnetic field only), then for $\Delta\lambda_B^* = 30$ mÅ we have $\Delta\lambda_H = 15$ mÅ, and $B \approx 78$ kG. The strongest magnetic fields in flares from distribution (1) reach of $B \approx 90$ kG ($\Delta\lambda_B^* = 36$ mÅ).

6. Discussion and Conclusion

6.1. Pashen-Back Effect

This effect takes place in strong fields when the following condition is met [20]:

$$m_j g \mu_B B \gg \Delta E_{ik} \quad (3)$$

where m_j is the magnetic quantum number, g is the Lande factor, μ_B is the Bohr magneton, B is the magnetic field strength, and ΔE_{ik} is the energy difference in the atom between the terms of multiplet structure.

The FeI $^5F_1^5D_0$ $\lambda = 5434.527$ Å line belongs to multiplet No. 15 [21]. Its Lande factors for lower (l) and upper (u) terms are following: $g_l = -0.014$, $g_u = 0/0$ that gives for effective Lande factor the above pointed value $g_{\text{eff}} = -0.014$. For lower splitted term with excitation potential $EP = 1.01$ eV, the nearest term is 5F_2 ; it has $EP = 0.99$ eV and $g = 0.5$. Then we can assume $\Delta E_{ik} = 0.02$ eV. Substituting \approx for \gg in Equation (3) yields an expression to estimate the “threshold” field B_p :

$$B_p \approx \Delta E_{ik} / m_j g \mu_B \quad (4)$$

Substituting in the numerical values of $m_j = 2$, $g = 0.5$ and $\mu_B = 0.58 \times 10^{-8}$ eV/G, we obtain $B_p \approx 3.4 \times 10^6$ G = 3.4 MG. Thus, one cannot expect any significant Paschen-Back effect in case of FeI 5434.5 line for field strength of $\sim 10^4$ G. This means that our application of formula (2) and magnetic field estimations are correct.

6.2. Possible MHD Model

Obviously, such very strong fields can not occur in the simplest case of untwisted magnetic flux tube. In fact, if magnetic flux tube is homogeneous and untwisted, the condition

$$p + B^2 / (8\pi) = p_{\text{ex}} \quad (5)$$

should be met, where p and p_{ex} are the gas pressures inside and outside the tube, respectively. If we assume $p_{\text{ex}} \sim 10^4$ dyn cm $^{-2}$ for upper photosphere level where the emission peaks of lines like FeI 5434.5 arise and also $p = 0$, we obtain $B \approx 500$ G only.

More strong fields should exist in twisted force-free magnetic structures. It was shown that a theoretical interpretation of “superstrong” magnetic field phenomena can be offered within the frame of a linear force-free model [22]. This model describes by Bessel’s functions J_0 and J_1 of zero and first orders; transversal magnetic field distribution is given on **Figure 9**. This model has a multipolar periphery and magnetic field till $\sim 10^4$ G with discrete values near the tube axis.

From the qualitative point of view, such a multipolar periphery could give a spectral signature similar to a strong ‘mixed polarity’ background field. The theoretical values of discrete strengths near the tube axis are following for photospheric level: 2.5 kG, 6.3 kG, 8.3 kG, 10.0 kG, 11.5 kG, 12.8 kG, etc. Thus, many modes of magnetic field structures can exist for the same gas pressure p_{ex} .

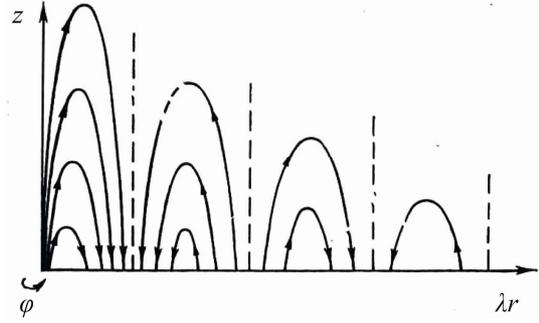


Figure 9. Transversal magnetic field distribution in linear force-free model by Soloviev and Lozitsky [22]. Here z and r are axis and radial coordinates, respectively. Entire lines with arrows indicate magnetic field lines.

However, it is doubtful that this model is suitable for interpretation above-named results with magnetic fields of 70 - 90 kG. The essence of the problem is in necessity of very large number of discrete layers with opposite magnetic polarity inside proper structure. Such layers are marked on **Figure 9** via stroke vertical lines. For instance, magnetic field strength of 2.5 kG (see above) can arise in a structure with one layer only. Further, field strength of 6.3 kG needs two layers, field of 8.3 kG - three layers, etc. It is clear, for field of 70 - 90 kG very many discrete layers are needed inside one small-scale (perhaps, subtelescopic) structure! Such many-lamellar magnetic field structure seems as unlikely. In other hand, namely such structures give a strong spectral camouflaging of the superstrong fields if these structures are spatially unresolved: in this case, we should observe essential widening of the magnetosensitive spectral line but practically zero Zeeman polarization at its wings. This assumption agrees with many-component structure of flare emission (see above **Figure 7**) and with presence of relatively strong unpolarized emission and weak-polarized (*i.e.* weak-split) one.

Naturally, named difficulty of theoretical interpretation does not refute the empirical data partly presented above on **Figures 4, 6** and **8**. The trustworthy polarized effects were discovered in several solar flares and these effects can not occur in kG fields—they need more strong fields, $\sim 10^4$ G. Probable existence of such superstrong fields presents a very important problem for modern solar physics.

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