

Effect of Temperature on the Capacitance of a Silicon Solar Cell in Static Regime

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Abstract

This work was based mainly on the capacitance of a silicon solar cell monofaciale in static conditions under monochromatic illumination and under the influence of temperature. Solving the continuity equation for excess minority carrier's charge density in the base of the solar cell resulting in the term of the temperature equation in the presence of an optical source. The latter has allowed us to obtain the photocurrent density and photovoltage. Subsequently we studied the capacitance situation of open circuit and short circuit depending on the temperature. Finally with the determination of capacitance under C_0 darkness.

Keywords

Capacitance, Photocurrent, Photo Voltage, Temperature

1. Introduction

Temperature is a very important parameter and cannot be neglected in the behavior of solar cells. Knowing that the total incident energy (100%), a small portion (13%) is reflected by the surface of the sensor and this small portion is extracted as electrical energy, therefore a large portion of energy is dissipated as heat [1]. This lead, under radiation at a relatively high operating temperature if not converted into electrical energy, is not drained. In this article, we will study the effect of temperature on the capacitance ability of a silicon solar cell monofaciale under monochromatic illumination under static conditions.

The study will be based primarily on the capacitance (open circuit in Short circuit, under dark), after studying the effect of temperature on the density of minority carriers, the photocurrent density, and photo voltage.

2. Theoretical Background

This solar cell is n + -p-p +(**Figure 1**) whose base is p-type. So the excess minority carrier's charge in the base is electrons. This translates into the greater mobility of electrons compared to that of [2] [3] [4] holes:

In the remainder of this study, the contribution of the issuer as well as the crystalline field that exists within the base has been neglected [5]. The solar cell is discussed in one dimension with the original axis taken at the junction. When the solar cell is illuminated, various phenomena occur there, governed by the following continuity equation [6]:

$$\frac{\partial^2 \delta(x)}{\partial x^2} - \frac{\delta(x)}{L^2} = -\frac{g(x)}{D}$$
(1)

$$g(x) = \phi(\lambda)\alpha(\lambda)(1 - R(\lambda))e^{(-\alpha(\lambda)x)}$$
(2)

$$\left(L(T)\right)^2 = \tau D(T) \tag{3}$$

$$D(T) = \mu \frac{k_b}{q} T \tag{4}$$

 $\delta(x)$ is denotes the density of electrons generated in the base at position x; (λ) is the flow monochromatic incident light; $\alpha(\lambda)$ is the absorption coefficient of monochromatic to the wavelength of material λ ; $R(\lambda)$ is the monochromatic reflectance of the material to the wavelength λ ; k_b is Boltzmann's constant; q is the elementary charge of an electron; τ is the lifetime of electrons in the base; D the diffusion coefficient of the electrons generated in the base [7]. It is a function of T; L is the diffusion length of the electrons generated in the database, it is linked to the flow of particles. It is a function of T; μ is the coefficient of electron mobility, g(x) is the generation rate of carriers.

Equation (1) is a differential of the second degree with second member whose general solution is:

$$\delta(x) = A \cosh\left(\frac{x}{L}\right) + B \sinh\left(\frac{x}{L}\right) - \frac{\phi(\lambda)\alpha(\lambda)L^2(1 - R(\lambda))e^{(-\alpha(\lambda)x)}}{D(L^2(\alpha(\lambda))^2 - 1)}$$
(5)





A and *B* are constants which are determined by the two following boundary conditions:

At the junction (x = 0) [8] [9]

$$\frac{\partial \delta(\omega, x, T)}{\partial x}\Big|_{x=0} = \frac{Sf}{D} \cdot \delta(\omega, 0, T)$$
(6)

At the rear side (x = H)

$$\frac{\partial \delta(x)}{\partial x}\Big|_{x=H} = -\frac{S_b}{D} \delta(x)\Big|_{x=H}$$
(7)

 S_f [8] and S_b [10] are respectively the junction recombination velocity at the junction and at the rear. Let us now study the excess minority carrier's density, whose profile is obtained from Equation (4).

3. Results and Discussion

3.1. Excess Minority Carrier's Density

Figure 1 and **Figure 2** are showed the excess minority carrier's charge density in open circuit and short-circuit condition depth in the database for different values of the temperature.

At Figure 1, we see that when the depth x increases, the carrier density decreases. And also for a low value of the depth (at about x < 0.005 cm) when the temperature increases, the carrier density decreases. By cons for x greater than about 0.005 cm, the carrier density is increased when the temperature is low. At Figure 2, we see that near 0 cm carrier density increases with depth and becomes maximum to values of x near close to 0 cm. From this maximum value,



Figure 2. density profile of excess minority carrier's charges of the solar cell short circuit depending on the depth in the base for different temperature values ($S_b = 2 \times 10^2$ cm/s, $S_f = 6 \times 10^6$ cm/s, $\lambda = 0.48$ µm.

the carrier density regresses when the depth x progresses. And also for a low value of the depth (at about the x is lower than 0.015 cm) when the temperature increases, the carrier's density decreases.

However, for *x* greater than about 0.015 cm, the temperature of excess minority carrier's density is increased. The decrease of excess minority carrier's density versus depth for the two figures is explained by the fact that there recombination volume of excess minority carrier. Regarding temperature has an excitatory function, which explains the behavior of the carrier density.

3.2. Photocurrent Density

The photocurrent resulting from the diffusion of the density of minority carriers at the junction and its expression is given by [6]-[11].

$$J_{ph}\left(S_{b}, S_{f}, \lambda, T\right) = qD \frac{\partial \delta\left(x, S_{b}, S_{f}, \lambda, T\right)}{\partial x} \bigg|_{x=0}$$
(8)

Figure 3 below shows the photocurrent density as a function of the recombination velocity for different values of the temperature.

We observe that for small values of the recombination velocity at the junction S_f (S_f less than 2 × 10² cm/s) corresponding to open circuit, the photocurrent density is almost zero because the carriers are stored at the junction. From this value up to almost 5 × 10⁵ cm/s, the photocurrent density increases with that of Sf because some carriers starts passing through the junction. High values S_f (S_f



Figure 3. Photocurrent density as a function of the junction recombination velocity for different temperature values ($S_b = 2 \times 10^2$ cm/s, $S_f = 6 \times 10^6$ cm/s, $\lambda = 0.48$ µm).

greater than 5×10^5 cm/s) corresponds to the short-circuit because the maximum of carriers across the junction. As between 2×10^2 and 5×10^5 cm/s when the temperature increases, the photocurrent density decreases. Therefore 5×10^5 cm/s is the short-circuit increasing the temperature causes one of the photocurrent density. Having made the study of the photocurrent density, are now considering the photo voltage.

3.3. Photo Voltage

The voltage harvested under monochromatic illumination across the solar cell and use the Boltzmann relation, the photo voltage can be expressed by [10]:

$$V_{ph}\left(S_{b}, S_{f}, \lambda, T\right) = V_{T} \ln\left(\frac{N_{b}}{n_{i}^{2}}\delta\left(0, S_{b}, S_{f}, \lambda, T\right) + 1\right)$$
(9)

where V_T is the thermal voltage, N_b doping rate, or the concentration of minority carriers charge.

$$V_T = \frac{k_b}{q}T\tag{10}$$

$$n_i^2 = A \cdot T^3 \exp\left(-\frac{E_g}{k_b \cdot T}\right) \tag{11}$$

A is a specific constant of the material; n_i is the density of intrinsic carrier temperature.

Figure 4 shows the photo voltage depending on the junction recombination velocity (Sf) for different temperature value.



Figure 4. Photo voltage profile as a function of the recombination rate for different temperature values ($S_b = 2 \times 10^2$ cm/s, $\lambda = 0$, 48 µm).

We note that for low values of the junction recombination velocity corresponds to the open circuit where the carriers are stored at the junction recombination velocity; the photo voltage is at its maximum. Then she starts to decline from certain values of junction recombination velocity, now some carriers are beginning to cross the junction. Photo voltage tends to invalidity for high junction recombination velocity values corresponding to short circuit when the maximum carriers pass the junction. Also when the temperature is increasing, the photo voltage is decreasing.

3.4. Dark Capacitance

The diffusion capacitance of the solar cell is known as the resulting ability of the load change during the diffusion process within the solar cell [12]-[19]. It is expressed by the following term:

$$C(S_b, S_f, \lambda, T) = \frac{\mathrm{d}Q(S_b, S_f, \lambda, T)}{\mathrm{d}V(S_b, S_f, \lambda, T)}$$
(12)

As

$$Q = q\delta(x=0) \tag{13}$$

was as follows:

$$C = q \frac{\mathrm{d}\delta(x=0)}{\mathrm{d}V} \tag{14}$$

Equation can be mentioned as follows:

$$C = q \frac{\mathrm{d}\delta(x=0)}{\mathrm{d}S_f} \frac{1}{\frac{\mathrm{d}V}{\mathrm{d}S_f}}$$
(15)

Taking into account the expression of the photo voltage and the excess minority carrier's density, we obtain the following expression:

$$C = \frac{q(n_0(T))^2}{N_b V_T} + \frac{q\delta(x=0, S_b, S_f, \lambda, T)}{V_T}$$
(16)

So

$$c = c_1 + c_2 \tag{17}$$

where

$$c_{1} = \frac{q\left(n_{0}\left(T\right)\right)^{2}}{N_{b}V_{T}}$$
(18)

is the intrinsic capacitance (ability to darkness) and depends on the nature of the material through the intrinsic concentration is the ability of the solar cell due to the illumination. The following two figures show the ability of the solar cell in open circuit condition and short circuit as a function of junction recombination velocity for different values of the temperature. n_0 , through doping N_b and V_T thermal potential.

$$c_2 = \frac{q\delta\left(x=0, S_b, S_f, \lambda, T\right)}{V_T}$$
(19)

where C_0 is the solar cell capacitance under dark. Equation 19 can be rewritten in the form

$$C(S_b, S_f, \lambda, T) = C_0 \exp\left(\frac{V_{ph}}{V_x}\right)$$
(20)

With the logarithm function, we rewrite equation 20 as:

$$\ln C - \ln C_0 = \frac{V_{ph}}{V_T} \tag{21}$$

Equation (21) can be expressed as follows:

$$\ln C\left(S_{b}, S_{f}, \lambda, T\right) = \frac{1}{V_{T}} \cdot V_{ph}\left(S_{b}, S_{f}, \lambda, T\right) + \ln C_{0}$$
(22)

At Figure 5 we observe that for small values of the junction recombination velocity S_i concordant to open circuit, which holders are stored; capacity is at its peak. At this level the capacity remains virtually constant. Also this time the increasing temperature causes a drop in capacity. We note in Figure 6 that the capacity regresses when the junction recombination velocity rises. This can be explained by the fact that the capacitance is the measure of the ability of the capacitor to store the charges on its plates; in other words, it is a measure of its power storage. Whereas when the junction recombination velocity is low (Figure 5) holders are almost immobile, the maximum capacity and decreases with temperature excites carriers where increased thermal agitation. Where junction



Figure 5. Capacitance of the solar cell open circuit according to the junction recombination velocity for different temperature values ($S_b = 2 \times 10^2$ cm/s, $\lambda = 0.48$ µm).



Figure 6. Capacitance of the solar cell shorted according to the junction recombination velocity for different temperature values.

recombination velocity is high (Figure 6) the holders are movable in reduced capacitance.

The following figure describes the ability of the solar cell in open-circuit condition in function of temperature for different wavelength values.

We find open circuit at the level of **Figure 7** that the capacity decreases with increasing temperature. And also decreases for shorter wavelengths. The following figure shows the capacity of the solar cell in a situation of short-circuit in function of temperature for different values of the wavelength.

At **Figure 8** where we have the ability to short-circuit situation; we note that the moment the temperature is increased, the capacity is decreasing. The following figure shows the natural logarithm of the capacitance C with different values of the natural logarithm of the capacity under darkness C_0 according to the photo voltage for different values of the temperature. The logarithm of the capacitance curves versus photo voltage for various temperature values are given in **Figure 9**.

From this figure we have determined for each temperature value in its corresponding capacitance C_0 darkness. The temperature value is a straight line of slop $1/V_T$

 C_0 these values are given by **Table 1**.

The study of this table shows that when the temperature rises, the darkness under C_0 capacity also increases.

4. Conclusions

A theoretical study of a solar cell with monofaciale under static regime under



Figure 7. Capacitance of the solar cell open-circuit in function of temperature for small and different values of the wavelength ($S_b = 2 \times 10^2$ cm/s, $S_f = 0$ cm/s).



Figure 8. Capacitance of the solar cell by short-circuit function of temperature for different low and the wavelength values ($S_b = 2 \times 10^2$ cm/s, $S_f = 6 \times 10^6$ cm/s).

multispectral illumination is presented. From different studies, we could establish that the temperature favors the presence of photo generated electrons in the base by inhibiting their diffusion and their conduction. The diffusion capacitance of the solar cell increases with the temperature. Finally, the dark capacitance of the solar cell was obtained for any temperature value on the capacitance



Figure 9. Logarithm of the capacitance versus photo voltage for various temperature values ($S_b = 2 \times 10^2$ cm/s, $\lambda = 0.48$ µm).

Table 1. C_0 (Capacitance darkness) values for the function of temperature.

Т	200 K	303 K	333 K	400 K
C_0	$8.54 \times 10^{-22} \text{ F}{\cdot}\text{cm}^{-2}$	$1.01 \times 10^{-12} \mathrm{F}{\cdot}\mathrm{cm}^{-2}$	$5.25 \times 10^{-11} \mathrm{F}{\cdot}\mathrm{cm}^{-2}$	$5.21 \times 10^{-8} \text{ F}{\cdot}\text{cm}^{-2}$

(open circuit in Short circuit, under dark). Based on the C-V characteristics, a graphical method has been proposed for the determination of dark capacitances C_0 .

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