

Modeling a General Equation for Pool Boiling Heat Transfer

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ABSTRACT

It is recognized that the nucleate pool boiling data available in literature are mainly related to four known correlations, each differs from the other by a varying magnitude of constant coefficients, depending on restrictive experimental conditions. The present work is concerned in developing an empirically generalized correlation, which covers the entire range of nucleate boiling with a minimum possible deviation from experimental data. The least squares multiple regression technique is used to evaluate the best coefficient value used in the correlations. An empirical correlation that fits a broader scope of available data has been developed by a non-linear solution technique leading to the following equation:

 $Nu_{B} = R_{3} \left(\frac{\Delta P}{\rho_{L}}\right)^{-0.65} \left(Pe_{B}\right)^{0.82} \left(\frac{\rho_{V}}{\rho_{L}}\right)^{1.1} \left(Pr\right)^{R_{4}} \text{ where the coefficients } R_{1} \text{ and } R_{3} \text{ both represent the effect of surface-liquid}$

combination. They are assessed independently for the used surface material and liquid.

Keywords: Pool Boiling; Nucleate Boiling; Linear and Non-Linear Technique Methods; Heat Transfer

1. Introduction

Boiling is a complex process and an intensive work is needed for its understanding. Within the last decades, several nucleate boiling models were formulated and could be grouped in two main categories: a) Bubble Agitation Models and b) Macro/Micro Layer Evaporation Models.

The bubble agitation models are based on the principle of agitating the liquid, but they carry away little heat. The heat transfer is considered within the turbulent forced convection. The obtained empirical pool boiling heat transfer models employ dimensionless groups based on both fluid and solid properties while the main constant in the model is found to depend on the geometry of the heater. The models found in literature are useful within the range of database used in developing their derivation.

Bubble agitation mechanism together with Helmholtz-instability mechanism can be used either to explain the heat transfer at the low heat flux regime or to explain CHF (Critical Heat Flux). They cannot account for the continuity of the pool-boiling curve. On the other hand, the macro/micro layer evaporation reproduces the poolboiling curve from the nuclear boiling to transition boiling. The macro/micro layer models play an important role in high heat flux region. The liquid layer includes the micro layer underneath the bubble and the macro layer on the base of coalescence and dries out periodically [1]. Heramura & Katto [1] assume the liquid-vapor interface is stationary and the entire surface heat flux contributes to macro layer evaporation.

Several numerical models were proposed based on the macro layer theory among that of Maruyama *et al.* [2]. Zhao *et al.* [3] put forward a model for transient pool boiling heat transfer. The model employed is too high heating rate to be realized in practical experiments for a horizontal surface. He *et al.* [4] concluded that the macro layer model is more suitable for the high heat flux regime. Dhir [5] confirmed that numerical simulations are not a substitute for detailed experiments. The experimental results are needed to validate the simulations. Numerical simulations provide additional insights into the boiling phenomena.

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Within the late decade, many researchers worked on viewing the pool boiling in microgravity (in the absent of buoyancy) to understand the lower limit of forced convection. Several workers are Lee [6], Herman [7], Wan & Zhao [8], and Kubota *et al.* [9].

Ji *et al.* [10] enhanced the pool boiling heat transfer in microgravity by using porous coating heating surface at atmospheric pressure and slightly moderate superheats.

Other researchers [11,12] enhanced the pool boiling by using nanofluid (water mixed with extremely small amount of nanosized particles). They concluded the enhancement of the thermal conductivity and convection heat transfer capability of the suspended particles of nanometer in size for many volume fractions of nanofluids.

The results of workers [6-12] can be used as a guidance in formulating proper equations that can be used in design. The aim of the present work is to use bubble agitation models to obtain a generalized empirical correlation that gives the best possible representation of collected data. Pool of data is collected from literature for various liquids effects with different plain test surfaces. For this purpose, linear and non-linear programming techniques were used in the evaluation of the proposed correlation.

2. Theoretical Analysis of Bubble Agitation Models

The primary requirement for nucleation to occur or for a nucleus to subsist in a liquid is that the liquid should be superheated. There are two types of nuclei. One type is formed in a pure liquid; it can be either a high energy molecular group resulting from thermal fluctuations of liquid molecules, or a cavity resulting from a local pressure reduction such as that occurs in accelerated flow. The other type, formed on a foreign object can be either a cavity on the heating wall or suspended foreign material with a non-wetted surface. Rohsenow [13] assumes that the movement of bubbles at the instant of breaking away from the heating surface is of prime importance and obtained Equation (1) for heat transfer in the region of nucleation pool boiling.

$$\frac{C_L(T_s - T_{sat})}{H_{fg}} = C_{sf} \left(\frac{q}{\mu_L H_{fg}} \sqrt{\frac{\sigma g_c}{g(\rho_L - \rho_V)}} \right)^{0.33} \operatorname{Pr}^r \quad (1)$$

The recommended variation of r is within 0.8 to 2.0. Evaluation of C_{sf} from experimental results of many workers [14] prove to be a parameter which does not pick out only the nucleation ability of heating surface but contains the effect of physical properties of liquid.

Rohsenow [15] proposed the surface factor C_{sf} to prescribe the condition of heating surface in nucleate boiling. Various investigators [14] utilized this factor in their determination of empirical expressions. The surface factor is defined by Equation (1), generally known as Rohsenow empirical correlation.

Forster and Zuber [16] indicated that small bubbles grow rapidly and large ones slowly, but the degree of agitation in the surrounding liquid due to bubble growth remains the same. They derived the following empirical correlation:

$$\frac{qC_L\rho_L\sqrt{\pi\alpha_L}}{K_LH_{fg}\rho_V} \left(\frac{2\sigma}{\Delta P}\right)^{1/2} \left(\frac{\rho_L}{\Delta P}\right)^{1/4} = 0.0015 \operatorname{Re}_b^{0.62} \operatorname{Pr}_L^{1/3}$$
(2)

Equation (2) predicts the same heat transfer coefficient for a liquid boiling on any hot surface (for all heterogeneous cases only) or boiling in bulk (for all homogeneous cases only). Rohsenow's Equation (1) was developed and applied to the heterogeneous case only.

Forster and Greif [17] suggested a different approach by considering that the mechanism of high heat transfer rate, during nucleate boiling, is mainly due to the liquidvapor exchange. They obtained a dimensional empirical correlation, for the pool boiling heat flux q in water at 100 - 4763 kN/m², as shown in Equation (3).

$$q = (4.3)(10^{-5}) \frac{K_L T_{sat} \Delta P^2}{\sigma^{1/2} \left(H_{fg} \rho_V\right)^{3/2}} \left[\left(C_L T_{sat} \alpha^{1/2}\right)^{1/4} \left(\frac{\rho_L}{\rho_V}\right)^{5/8} \left(\frac{\mu_L C_L}{K_L}\right)^{1/3} \right]$$
(3)

This correlation is not as widely verified as that of Rohsenow.

Gupta and Varshney [18] obtained experimental data for boiling heat transfer, using distilled water, benzene and toluene as liquids over a heated horizontal cylinder made of stainless steel. Their data was correlated by the following dimensionless empirical correlation:

$$Nu_{B} = 1.39 \left(Pe_{B} \right)^{0.7} \left(\frac{\rho_{V}}{\rho_{L}} \right)^{0.21} \left(Pr \right)^{-0.21}$$
(4)

where Nu_B and Pe_B are the Nusselt and Peclet number of boiling respectively. Or it can be written as:

$$\frac{h}{K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}} = 1.39 \left(\frac{q\rho_L C_L}{\rho_V H_{fg} K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}}\right)^{0.7} \left(\frac{\rho_V}{\rho_L}\right)^{0.21} \left(\frac{\mu_L C_L}{K_L}\right)^{-0.21}$$
(5)

In order to derive a general correlation based on bub-

ble agitation phenomena to be more versatile than the correlations existing in literature, a search was made through published work in literature and found that the following four well known empirical correlations referred to in most publications:

- a) Rohsenow correlation (Equation(1))
- b) Forster and Zuber correlation (Equation (2))
- c) Forster and Greif correlation (Equation (3))
- d) Gupta and Varshney correlation (Equation (5)).

For the sake of analysis, experimental data collected from many literatures [19-24] tabulated as heat flux (*q*), surface temperature (*T_s*), heat transfer coefficient (*h*), and coefficient h^* (equal to $h/q^{0.7}$). Moreover, physical and thermodynamic properties collected at the reported experimental conditions from literature [25-29] to be used in the analysis. The properties include, liquid thermal conductivity (*K_L*), liquid heat capacity (*C_L*), density of vapor (ρ_V), surface tension of liquid (σ), saturation temperature (*T_{sat}*), density of liquid (ρ_L), latent heat of vaporization (*H_{fg}*), and viscosity of liquid (μ_L).

The above stated correlations are of dimensionless form with the exception of the Forster and Greif correlation (Equation (4)). These equations can be represented by general equation as shown in **Appendix**. Many modifications to linear correlations have been tried to minimize the sum of squares of errors and to conclude some general correlations.

3. Results and Discussion

Boiling heat transfer studied earlier indicated that several variables are important in nucleate boiling such as pressure, fluid properties, surface condition, boiling temperature, kind and relative amount of impurities. The practical data showed that changes in magnitude of these properties and conditions could significantly affect pool boiling heat transfer.

A graphical analyses for 56 sets of literature data was used in studying the effect of heat flux, (q), and operating pressure (P), on boiling heat transfer coefficient, (h). **Figures 1-3** show the variation of heat transfer coefficient with heat transfer flux, (q). The lines in the figures are the best-fit lines of the reported data. **Figure 1** plotted for various liquids at different operating pressure and test surface. **Figure 2** corresponds to various liquid-surface combinations at constant atmospheric pressure. **Figure 3** reflects the behavior of various metal surfaces and operating pressures for the same liquid.

All the data can be represented by the empirical, Equation (6), with an average percentage error ranging from 0.012 to 11.8

$$h = h^* q^{0.7} (6)$$

The proportionality constant
$$h^{+}$$
 is proved to be a func-
tion of pressure and liquid-surface combination. Cichelli



Figure 1. Heat transfer coefficient (h) versus heat flux (q) at various pressures.



Figure 2. Heat transfer coefficient (h) versus heat flux (q) at atmospheric pressure.



Figure 3. Heat transfer coefficient (h) versus heat flux (q) for various metal surfaces.

and Bonilla [19] confirmed that the coefficient of heat transfer increases with absolute operating pressure in the nucleate boiling zone. They reported the following correlation:

$$h = 5400 \left(P\right)^{0.44} \tag{7}$$

Figure 4 shows the variation of the proportionality constant (h^*) as function of pressure (P) for a definite

liquid-surface combination. Equation (8) represents the relationship between h^* and P, that was obtained from best data fit of **Figure 4**.

$$h^* \alpha P^{0.25}$$
 (8)

The overall dependence of (h) on operating pressure (P) and heat flux (q) for different liquid-surface combinations is shown in **Figure 5**.

3.1. Linear Programming Analysis of Empirical Correlations

Equation (A.4) in **Appendix** used to test the validity of the published correlations. It is used to formulate a correlation that shows the best fit of the experimental data. All the data cited in the literature from [19-24] classified as eight liquids (Water, Benzene, Methanol, Carbon Tetrachloride, n-Butanol, Isopropanol, n-Amyl Alcohol, and n-Heptane) and four surfaces (Brass, Copper, Nichrome, and Stainless Steel) at different operating pressures grouped in 56 data sets.

The applicability of the four empirical correlations, Equations (1)-(3) and (5), in representing the data was test by using linear-programming; that by fixing some of



Figure 4. Proportionality constant (h^*) versus operating pressure (P).



Figure 5. Heat transfer coefficient (h) dependence on operating pressure and heat flux (q).

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the coefficients and evaluating the others and the average percentage error is used as test criteria for comparison purposes.

3.1.1. Rohsenow's Correlation

Equation (9) is a general expression for Rohsenow's correlation while the exact expression stated as in Equation (1).

$$\frac{C_L(T_s - T_{sat})}{H_{fg}} = C_{sf} \left(\frac{q}{\mu_L H_{fg}} \sqrt{\frac{\sigma}{\rho_L - \rho_V}}\right)^{R_1} \left(\frac{C_L \mu_L}{K_L}\right)^{R_4}$$
(9)

The coefficient C_{sf} reported, in the literature, to vary with each liquid-surface combination and it is independent of pressure [22]. The validity of Equation (9) was tested for the entire collected data by applying the least-squares method.

In the initial analysis of data, the pressure was assumed constant and the obtained results showed inconsistency in the calculated values of constants for various experimental conditions tested. The inconsistency in values of constants is most likely due to pressure effect, which was not considered as variable during the initial analysis of data. On the next try, a pressure parameter was introduced in an attempt to narrow the variation in the values of constants for different systems and to conclude general correlation. Four different expressions of pressure parameter cited from literature [16], and stated in Equation (10), was used and expected to have an effect on heat transfer in the region of pool boiling nucleation:

$$R_{3}\left(\frac{\Delta P}{\rho_{L}}\right)^{R_{2}}, R_{3}\left(\frac{\Delta P}{\rho_{V}}\right)^{R_{2}}, R_{3}\left(\frac{\Delta P}{\sigma}\right)^{R_{2}},$$
and
$$R_{3}\left(\frac{\Delta P}{\rho_{L}}\right)^{R_{2}}\left(\frac{\Delta P}{\sigma}\right)^{R_{5}}$$
(10)

Each of these pressure expressions selected to replace the coefficient C_{sf} in Equation (9) and the obtained results were compared and checked. The analysis found

that the pressure parameter $R_3 \left(\frac{\Delta P}{\rho_L}\right)^{\kappa_2}$ gives the lowest

minimum percentage error between the others forms and selected to replace C_{sf} and modify Equation (9) to the form showing in Equation (11).

$$\frac{C_L(T_s - T_{sat})}{H_{fg}} = R_3 \left(\frac{\Delta P}{\rho_L}\right)^{R_2} \left(\frac{q}{\mu_L H_{fg}} \sqrt{\frac{\sigma}{\rho_L - \rho_V}}\right)^{R_1} \left(\frac{C_L \cdot \mu_L}{K_L}\right)^{R_4}$$
(11)

To make the present work more general, the analysis was repeated for various liquids at certain test surface and different operating pressures, that by calculating all the coefficients (R_1 , R_2 , R_3 , and R_4) of Equation (11) for each set of data. The values of coefficients were selected, by the help of Equation (A.4), on the basis of using Equation (11) with the lowest average percentage error.

3.1.2. Gupta and Varshney Correlation

Equation (12) is a general expression for Gupta and Varshney empirical correlation, while Equation (5) represent the exact form.

$$\frac{h}{K_L}\sqrt{\frac{\sigma g_c}{\left(\rho_L - \rho_V\right)g}} = R_3 \left(\frac{q\rho_L C_L}{\rho_V H_{fg} K_L}\sqrt{\frac{\sigma g_c}{\left(\rho_L - \rho_V\right)g}}\right)^{R_1} \left(\frac{\rho_V}{\rho_L}\right)^{R_6} \left(\frac{\mu_L C_L}{K_L}\right)^{R_4}$$
(12)

The pressure term is not included in the correlation of Gupta and Varshney as was the case of Rohsenow's empirical correlation, Equation (1). Equation (12) differs from Rohensow's Equation (9) by including the density ratio term and classifying other dimensionless term in well-known groups. A similar data analysis used for Equation (12) as it was with the case of Rohsenow's

Correlation. By substituting the various pressure forms of Equation (10) in place of coefficient R_3 in Equation (12), it is found that pressure expression $R_3 (\Delta P/\rho_L)^{R_2}$ gave the minimum average percentage error. This term is then chosen as the best fit expression for pressure and used in Equation (12) to obtain Equation (13).

$$\frac{h}{K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}} = R_3 \left(\frac{\Delta P}{\rho_L}\right)^{R_2} \left(\frac{q \rho_L C_L}{\rho_V H_{fg} K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}}\right)^{R_1} \left(\frac{\rho_V}{\rho_L}\right)^{R_6} \left(\frac{\mu_L C_L}{K_L}\right)^{R_4}$$
(13)

In order to make the present work more general, the analysis repeated for various liquids at specific test surface and different operating pressures then follow the same procedure as in case of Rohsenow's correlation. The analysis concluded that the modified Equation (13) provides a better data fit than that of Gupta and Varshney Equation (5).

3.1.3. Forster and Zuber Correlation

Equation (14) is a general expression for Forster and Zuber empirical correlation, while the exact form is given in Equation (2).

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$$\frac{qC_L\rho_L}{K_LH_{fg}\rho_V}\sqrt{\frac{\pi K_L}{\rho_L C_L}} = R_3 \left(\frac{2\sigma}{\Delta P}\right)^{R_5} \left(\frac{\rho_L}{\Delta P}\right)^{R_2} \left[\frac{\rho_L}{\mu_L} \left(\frac{C_L\rho_L(T_s - T_{sat})}{H_{fg}\rho_V}\sqrt{\frac{\pi K_L}{\rho_L C_L}}\right)^2\right]^{R_1} \left(\frac{\mu_L C_L}{K_L}\right)^{R_4}$$
(14)

The pressure term is taken care of in Forster and Zubers correlation, Equation (2). By applying a similar procedure as in previous cases, it was found that the validity of the above equation is restricted to specific experimental data near critical temperature difference. By checking the above results for Forster and Zuber empirical correlation, Equation (14), it is found that the overall average percentage errors is very high in predicting the published data under consideration.

3.1.4. Forster and Greif Empirical Correlation

Equation (15) is a general expression for Forster and Greif dimensional empirical correlation, while the exact form is given in Equation (3).

$$h = R_1 \left(T_s - T_{sat} \right)^{R_2} \left(\sigma \right)^{R_3} \left(\rho_L \right)^{R_4} \left(\Delta P \right)^{R_5} \left(K_L \right)^{R_6} \left(H_{fg} \right)^{R_7} \left(\rho_V \right)^{R_8} \left(C_L \right)^{R_9} \left(\mu_L \right)^{R_{10}}$$
(15)

Using the same sets of data analyzed previously gave a higher average percentage error as compared with the dimensionless empirical correlations as shown in **Table 1**.

The variation of some thermodynamic properties as function of pressure and temperature is not reported in literature and these properties were considered constant during the calculations, which may be the cause of the large average percentage error. Moreover, this empirical correlation dealt with fluid side effect of the problem and ignored the surface side effect on the nucleate boiling behavior. This might have added more error to the validity of this correlation.

The linear programming analysis recommended the use of the modified correlations of Rohsenow and Gupta & Varshney as they give closer prediction to the experimental data than in case of using Forster & Zuber and Forster & Greif modified correlations as showing in **Table 1**. Hence, the last two correlations excluded from any further analysis.

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No.	No. of Sets	Material of Test Surface	Equation (11)		Equation (13)		Equation (14)		Equation (15)	
			APE	CC	APE	CC	APE	CC	APE	CC
1	6	Nichrome	8.9	0.995	2.1	0.935	18.5	0.968	35.0	0.845
2	15	Copper	14.1	0.965	22.4	0.995	82.3	0.903	134.8	0.882
3	22	Brass	21.1	0.807	25.1	0.933	80.7	0.740	65.3	-0.04
4	13	Stainless Steel	15.6	0.858	16.2	0.553	78.4	0.586	74.0	0.228

Table 1. Linear programming results for Equation (11) to Equation (15).

Where APE = Average Percentage Errors; CC = Correlation Coefficient.

3.2. Non-Linear Programming Analysis of Empirical Correlations

The data were re-analyzed by non-linear programming methods by using the modified correlations of Rohsenow, Equation (11), and Gupta & Varshney Equation (13) in an attempt to improve the correlations for lower average percentage error.

By assume that:
$$\left[\frac{q}{\mu_L H_{fg}}\sqrt{\frac{\sigma}{(\rho_L - \rho_V)}}\right]$$
 in Equation
(4) is equal to X and $\left[\frac{q\rho_L C_L}{\rho_V H_{fg} K_L}\sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}}\right]$ in

Equation (13) is equal to Y, then use either of the following equations:

Binominal Expression:

$$\left[R_{3}(X \text{ or } Y)^{R_{1}} = (R_{7} + R_{8}(X \text{ or } Y))^{R_{9}}\right]$$
(16)

or non-linear Expression:

$$\left[R_{3}(X \text{ or } Y)^{R_{1}} = R_{7} + R_{8}(X \text{ or } Y)^{R_{9}}\right]$$
(17)

to replace the $R_3 X^{R_1}$ in Equation (11) or $R_3 Y^{R_1}$ in Equation (13) respectively. It is found from data fittings that a better representation can be obtained by using the binominal expression, Equation (16), with

 $(R_7 + R_8(X))^{R_9}$ in place of $R_3 X^{R_1}$ in Equation (11) and using the binominal expression, Equation (16), with $(R_7 + R_8(Y))^{R_9}$ in place of $R_3 Y^{R_1}$ in Equation (13).

Various forms of expressions for the constant R_7 was tried for both Rohsenow and Gupta & Varshney modi-

fied correlations. It was found that keeping R_7 as a constant and independent of other parameters yielded a better data representation. It was also concluded that the average percentage error would not be improved by using the non linear equations instead of the linear equation as showing in **Table 2**.

General Empirical Correlation

By using the best data fit for Equation (11) for different surfaces (Nichrome, Copper, Brass, and Stainless Steel), the variation of the powers of pressure expression term and Prandtl number in the equation were found to be approximately equal to 0.08 and 1.0 respectively leading to the following generalized Equation (18).

$$\frac{C_L (T_s - T_{sat})}{H_{fg}} = R_3 \left(\frac{\Delta P}{\rho_L}\right)^{0.08} \left[\frac{q}{\mu_L H_{fg}} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}}\right]^{R_l} \left(\frac{\mu_L C_L}{K_L}\right)^{1.0}$$
(18)

The coefficients R_1 and R_3 represent the effect of surface-liquid combination. They are assessed independently for each surface by the least-squares linear regression method and the results are stated in **Table 3**.

A similar analysis tried for Equation (13) and concluded that the powers of pressure term, Peclet number (Pe_B), and density ratio term (ρ_V/ρ_L) were relatively independent of surface—liquid combination as compared with the coefficient R_3 and the power of Prandtl number, R_4 . The best form of Equation (13) was tested for different data sets and concluded Equation (19).

$$\frac{h}{K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)}} = R_3 \left(\frac{\Delta P}{\rho_L}\right)^{-0.65} \left[\frac{q \rho_L C_L}{\rho_V H_{fg} K_L} \sqrt{\frac{\sigma g_c}{(\rho_L - \rho_V)g}}\right]^{0.82} \left(\frac{\rho_V}{\rho_L}\right)^{1.1} \left(\frac{\mu_L C_L}{K_L}\right)^{R_4}$$
(19)

A similar way was followed for Equation (19), to that of Equation (18), in finding the power R_4 and the coefficient R_3 and their best values are given in **Table 3**. The result of analysis of Equation (18) and Equation (19) listed in **Table 3** suggested the use of Equation (19) in preference to Equation (18).

The applicability of Equaiton (19) was examined for different surfaces as showing in **Figures 6-9**. The equation found to fit well for all the data with the exception of Brass. The deviation in the results for Brass is due to

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No. of Sets	Material of Test Surface	Equation (11)						Equation (13)					
		Linear		Binomial		Non Linear		Linear		Binomial		Non Linear	
		APE	CC	APE	CC	APE	CC	APE	CC	APE	CC	APE	CC
6	Nichrome	8.9	0.995	13.2	0.994	15.0	0.996	2.1	0.935	2.6	0.925	2.5	0.995
15	Copper	14.1	0.965	29.3	0.955	17.6	0.980	22.4	0.995	31.7	0.999	32.9	0.991
22	Brass	21.1	0.807	59.5	0.875	98.0	0.576	25.1	0.933	45.9	0.968	45.7	0.967
13	Stainless Steel	15.6	0.858	17.0	0.665	25.9	0.796	16.2	0.553	18.7	0.989	20.2	0.987

Table 2. Comparison of linear, binominal, and non linear expressions versions of Equations (11) and (13).

Table 3. Linear programming results for Equations (18) and (19).

					Equation (18)		Equation (19)					
No.	No. of Sets	Test Surface Material	R_3	R_1	Average Percentage Error	Correlation Coefficient	R_3	R_4	Average Percentage Error	Correlation Coefficient		
1	6	Nichrome	0.0019	0.684	15.5	0.976	19145.12	-1.300	6.4	0.954		
2	15	Copper	0.0080	0.256	25.4	0.891	5438.43	-0.758	15.3	0.978		
3	22	Brass	0.0045	0.228	21.1	0.907	4383.20	-0.132	36.0	0.771		
4	13	Stainless Steel	0.0093	0.300	27.6	0.705	4944.28	-0.470	27.3	0.893		



Figure 6. Experimental data predictions using Equation (19) for Nichrome surface.



Figure 7. Experimental data predictions using Equation (19) for Copper surface.



Figure 8. Experimental data predictions using Equation (19) for Brass surface.



Figure 9. Experimental data prediction using Equation (19) for Stainless Steel surface.

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limited available data at very low pressure. Equation (20) represents the dimensionless form of Equation (19).

$$Nu_{B} = R_{3} \left(\frac{\Delta P}{\rho_{L}}\right)^{-0.65} \left(Pe_{B}\right)^{0.82} \left(\frac{\rho_{V}}{\rho_{L}}\right)^{1.1} \left(Pr\right)^{R_{4}}$$
(20)

The analysis concluded that Equation (20) is valid for the entire available data and represent a more generalized correlation than the correlations found in literature.

4. Conclusions

A graphical analysis concluded that the empirical Equation (6) is showing the effect of heat flux (q) and operating pressure (P) on the boiling heat transfer coefficient (h).

$$h = h^* q^{0.7} \tag{6}$$

where h^* is a function of pressure and for different liquid-surface combinations, it is found to vary with the pressure as follows:

$$h^* \alpha P^{0.25} \tag{8}$$

56 sets of literature data were tested on each of the four known correlations, Rohsenow, Forster & Zuber, Forster & Greif, and Gupta & Varshney, by using the linear and non-linear programming solution. The concluded results show that any of these correlations does not fit the entire data satisfactory. To improve their predictions, the correlations were modified, including additional parameters in an attempt to close up the deviation in the values of calculated parameters. The modified correlations of Rohsenow and Gupta & Varshney responded better to the applied modification than that of Foster & Zuber and Foster & Grief and they were considered for further analysis.

The least squares multiple regression technique [30,31] is used to evaluate the best possible values of the constant coefficients in the correlation. The cumulative error squares were minimized by using an ordinary optimum seeking technique. Linear, binominal & non-linear correlations were tested in concluding the final correlation.

The use of non-linear solution technique did not improve correlations 11 and 13 that were concluded by the linear technique and hence Equation (20) gives the best representation of the entire tested data.

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Appendix

Sum of Squares of Errors

Equations (1) to (4) can be represented by a general equation:

$$y_i = \left(x_{i1}^{a1}\right)\left(x_{i2}^{a2}\right)\cdots\left(x_{ik}^{ak}\right)$$

where $i = 1, 2, 3, \dots, N$

$$y_i = \prod_{j=1}^k x_{ij}^{aj} \tag{A.1}$$

This equation represents a general form of those correlations and simplified by taking logarithms of both sides.

$$\ln y_i = a_1 \ln x_{i1} + a_2 \ln x_{i2} + a_3 \ln x_{i3} + \dots + a_k \ln x_{ik}$$

or

$$\ln y_i = \sum_{j=1}^k a_j \ln\left(x_{ij}\right) \tag{A.2}$$

For *N* number of data readings there will be *N* number of linear equations, while for the determination of *k* coefficients only *k* equations are required. A least squares multiple regression technique [30,31] was used to evaluate the best possible coefficient from raw data readings. The cumulative error squares minimized by an ordinary optimum seeking technique [31] resulting into k number of equations to provide *k* number of coefficients for the entire data. These equations mathematically represented in the form:

$$E_i = \ln y_i - \sum_{j=1}^k a_j \ln(x_{ij})$$
 (A.3)

The sum of squares of errors is expressed as:

$$E_{s} = \sum_{i=1}^{N} \left(\ln y_{i} - \sum_{j=1}^{k} a_{j} \left(x_{ij} \right) \right)^{2}$$
(A.4)

Nomenclature

 C_L —Heat capacity of liquid, J/kg·°C C_{sf} —Surface factor F_{μ} —Nucleation factor

g—Acceleration of gravity, m/s^2 g_c —Conversion ratio, kg m/kg·s² *h*—Heat transfer coefficient $W/m^2 \cdot C$ h^* —Proportionality constant H_{fg} —Latent heat of vaporation, J/kg K_L —Thermal conductivity, W/m °C Nu_B —Nusselt number for boiling = $(h/K_L)(\sigma g_c/(\rho_L - \rho_V)g)^{\prime}$ \dot{P} —Operating pressure (kN/m²) Pe_B —Peolet number for boiling = $(q/\rho_V H_{f_{e}}\alpha_L)(\sigma g_c/(\rho_L - \rho_V)g)^{\prime}$ Pr—Prandtl number *q*—Heat Flux, W/m^2s $R_1, R_2, R_3, \cdots, R_9$ —coefficients Re-Reynolds number Re_b —Reynolds number for bubbles = $(\alpha \wedge \pi)^2$

$$\frac{\rho_L}{\mu_L} \left(\frac{C_L \rho_L \Delta I \sqrt{\pi \alpha_L}}{H_{fg} \rho_V} \right)$$

Sr—Superheat ratio = $\frac{C_L (T_s - T_{sat})}{H_{fg}}$

 T_s —Test surface temperature, °C

 T_{sat} —Saturation temperature, °C

Greek Letters

- α_L Thermal diffusivity, m²/s, $K_L/(r_L C_L)$
- ΔP Pressure difference corresponds to $(T_s T_{sat})$, kN/m²
- μ_L Viscosity of liquid, kg/m·s
- ρ_V Density of vapor, kg/m³
- ρ_L Density of liquid, kg/m³
- σ Surface tension, kg/s²

Subscripts

- *b* Refers to bubble property
- *B* Refers to boiling condition
- L Refers to liquid condition
- sf Refers to surface factor
- *s* Refers to surface condition
- sat Refers to saturation condition
- v Refers to vapor condition