

Crystallization in the Three-Component Systems Rb₂SO₄-MSO₄-H₂O (M = Mg, Co, Ni, Cu, Zn) at 298 K

Veronika Karadjova¹, Donka Stoilova^{2*}

¹Department of Inorganic Chemistry, University of Chemical Technology and Metallurgy, Sofia, Bulgaria; ²Institute of General and Inorganic Chemistry, Bulgarian Academy of Sciences, Sofia, Bulgaria. Email: *stoilova@svr.igic.bas.bg

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ABSTRACT

The crystallization in the three-component systems Rb_2SO_4 -MSO_4-H₂O (M = Mg, Co, Ni, Cu, Zn) is studied by the method of isothermal decrease of supersaturation. It has been established that isostructural double compounds, $Rb_2M(SO_4)_2 \cdot 6H_2O$ (M = Mg, Co, Ni, Cu, Zn), $(SG P2_1/c (C_{2h}^5))$, crystallize from the ternary solutions within wide concentration ranges. The infrared spectra are discussed with respect to the normal vibrations of the sulfate ions and water molecules. The unit-cell group theoretical treatment of the double salts is presented. The extent of energetic distortions of SO_4^{2-} guest ions (about 2 mol%) matrix-isolated in the respective selenates, $M'_2M''(SeO_4)_2 \cdot 6H_2O$ (M' = K, Rb, NH_4^+ ; M'' = Mg, Co, Ni, Cu, Zn), is commented.

Keywords: Rb₂Me(SO₄)₂·6H₂O (Me = Mg, Co, Ni, Cu, Zn); Solubility Diagrams; X-Ray Powder Diffraction; Infrared Spectra; Matrix Infrared Spectroscopy

1. Introduction

The rubidium double sulfates belong to a large number of isomorphous compounds with a general formula

 $M'_2M''(XO_4)_2 \cdot 6H_2O$ (M' = K, Rb, NH_4^+ , Cs; M'' = Mg, Mn, Co, Ni, Cu, Zn; X = S, Se) known as Tutton salts. They crystallize in the monoclinic space group $P2_1/c(C_{2h}^5)$ with two formula units in the unit-cell. The crystal structures are built up from isolated octahedra, [$M''(H_2O)_6$], (three crystallographically different water molecules are coordinated to the M'' ions) and tetrahedra XO_4. The polyhedra are linked by hydrogen bonds. All atoms, except the divalent metal ions, which lie at centre of inversion C_i , are located at general positions C_1 . Recently, the crystal structures of some rubidium Tutton sulfates have been reported in Refs. [1,2]. As an example the crystal structures of Rb_2Co(SO_4)_2 \cdot 6H_2O is presented in **Figure 1**.

In this paper we present the results on the study of the solubility in the three-component systems

 Rb_2SO_4 -MSO_4-H₂O (M = Mg, Co, Ni, Cu, Zn) at 298 K. The double Tutton compounds $Rb_2M(SO_4)_2$ ·6H₂O are characterized by means of both the infrared spectroscopy and the X-ray powder diffraction methods. The vibrational behavior of SO_4^{2-} guest ions matrix-isolated in Tutton selenates, $M'_2M''(SeO_4)_2 \cdot 6H_2O$ (M' = K, Rb, NH_4^+ ; M'' = Mg, Co, Ni, Cu, Zn), is analyzed. A practical point of studying is that the Tutton compounds could be considered as proton conductors due to the existence of comparatively strong hydrogen bonds determined by the strong proton acceptor capabilities of the sulfate ions.

2. Experimental

 Rb_2SO_4 was prepared by neutralization of Rb_2CO_3 with dilute sulfuric acid solutions at 333 - 343 K. Then the solutions were filtered, concentrated at 323 - 333 K, and cooled to room temperature. The crystals were filtered, washed with alcohol and dried in air. The sulfates of divalent metals were commercial products. All reagents used were of reagent grade quality (Merck). The solubility in the three-component systems Rb_2SO_4 -MSO₄-H₂O (M = Mg, Co, Ni, Cu, Zn) at 298 K was studied using the method of isothermal decrease of supersaturation. Solutions containing different amount of the salt compounds corresponding to each point of the solubility diagrams

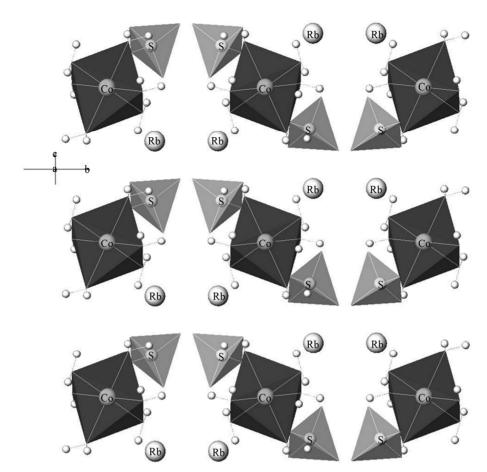


Figure 1. Crystal structure of Rb₂Co(SO₄)₂·6H₂O (structural data from [1]).

were heated at about 333 - 343 K and cooled to room temperature. Then the saturated solutions were vigorously stirred [3,4]. The equilibrium between the liquid and solid phases was reached in about 20 hours. The analysis of the liquid and the wet solid phases was performed as follows: the M" ion contents were determined complexonometrically at pH 9.5 - 10 using eriochrome black as indicator (magnesium ions) and at pH 5.5 - 6 using xylenol orange as indicator (cobalt, nickel, copper and zinc ions); the sulfate ions were precipitated as Ba- SO_4 with $Ba(NO_3)_2$ solutions and the concentrations of the excess of Ba²⁺ ions were determined complexonometrically using eriochrome black as indicator; the concentrations of the rubidium sulfate were calculated by difference [5]. The compositions of the solid phases were identified by means of both the X-ray diffraction and the infrared spectroscopy methods. Tutton compounds

 $Rb_2M(SeO_4)_2 \cdot 6H_2O$ were prepared according to the solubility diagrams of the three-component systems

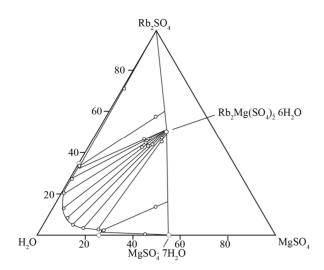
Rb₂SeO₄-MSeO₄-H₂O (M = Mg, Co, Ni, Cu, Zn) [6-8]. The samples, $M'_2M''(SeO_4)_{1.98}(SO_4)_{0.02} \cdot 6H_2O$ (M' = K, Rb, NH_4^+ ; M'' = Mg, Co, Ni, Cu, Zn), were prepared by crystallization from ternary selenate solutions in the presence of SO_4^{2-} ions (the data for the potassium and ammonium selenates are taken from [9]).

The infrared spectra were recorded on a Bruker model IFS 25 Fourier transform interferometer (resolution $< 2 \text{ cm}^{-1}$) at ambient temperature using KBr discs as matrices. Ion exchange or other reactions with KBr have not been observed. In some cases Lorentz band profile for multi peak data was used to determine the correct band positions corresponding to asymmetric stretches of the included SO₄²⁻ ions (ORIGIN PRO 6.1). The X-ray powder diffraction spectra were collected within the range from 5° to 50° 2 θ with a step 0.02° 2 θ and counting time 35 s/step on Bruker D8 Advance diffractometer with CuK α radiation and LynxEye detector. The lattice parameters of the double salts were calculated using the program ITO and refined with the program LSUCR.

3. Results and Discussions

3.1. Solubility Diagrams of the Three-Component Systems Rb₂SO₄-MSO₄-H₂O (M = Mg, Co, Ni, Cu, Zn) at 298 K

The solubility diagrams of the above systems are presented in Figures 2-6 (the respective experimental data



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Figure 2. Solubility diagram of the three-component system Rb₂SO₄-MgSO₄-H₂O at 298 K.

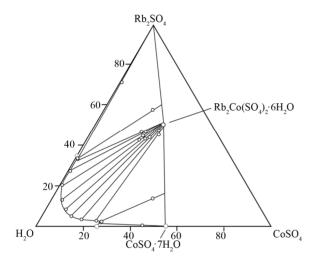


Figure 3. Solubility diagram of the three-component system Rb_2SO_4 - $CoSO_4$ - H_2O at 298 K.

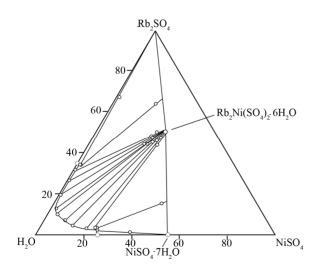


Figure 4. Solubility diagram of the three-component system Rb_2SO_4 -NiSO_4-H_2O at 298 K.

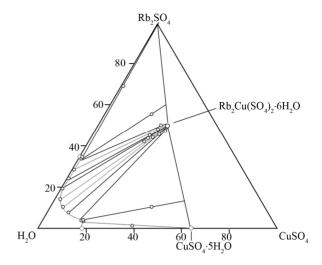


Figure 5. Solubility diagram of the three-component system Rb₂SO₄-CuSO₄-H₂O at 298 K.

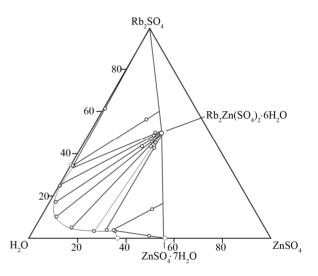


Figure 6. Solubility diagram of the three-component system Rb₂SO₄-ZnSO₄-H₂O at 298 K.

are summarized in **Tables 1-5**). It is seen that the simple salts Rb_2SO_4 , $MgSO_4$ · $7H_2O$, $CoSO_4$ · $7H_2O$, $NiSO_4$ · $7H_2O$, $CuSO_4$ · $5H_2O$ and $ZnSO_4$ · $7H_2O$ crystallize within very narrow concentration ranges, whereas the rubidium double sulfates crystallize within wide concentration ranges, thus indicating that strong complex formation processes occur in the ternary solutions.

3.2. X-Ray Powder Diffraction Data of Rubidium Tutton Compounds

The X-ray powder diffraction patterns of the rubidium Tutton compounds are shown in **Figure 7**. The double salts form monoclinic crystals $\left(\text{SG }P2_1/c\left(C_{2h}^5\right)\right)$.

The calculated unit-cell parameters are presented in **Table 6**. Our results coincide well with those determined from single crystal X-ray diffraction data [1,2].

		-		
Liquid phase mass%		Wet solid p	hase mass%	Composition of the
Rb_2SO_4	$MgSO_4$	Rb_2SO_4	$MgSO_4$	solid phases
35.10				Rb ₂ SO ₄
34.58	1.12	74.41	0.79	-
34.42	2.08	59.17	16.58	$\begin{array}{c} Rb_2SO_4 + \\ Rb_2Mg(SO_4)_2{\cdot}6H_2O \end{array}$
34.07	2.03	53.26	21.09	$Rb_2Mg(SO_4)_2$ ·6H ₂ O
29.39	2.48	50.02	19.96	-
25.85	4.13	48.62	20.14	-
21.19	6.11	49.57	21.75	-
17.17	8.66	46.68	21.24	-
13.37	12.96	47.87	23.36	-
9.45	19.25	46.71	23.96	-
6.97	24.01	48.37	24.83	-
7.38	27.54	49.58	25.68	-
7.51	28.51	20.54	37.42	$\begin{array}{c} MgSO_4{\cdot}7H_2O + \\ Rb_2Mg(SO_4)_2{\cdot}6H_2O \end{array}$
6.87	28.73	3.19	42.21	$MgSO_4 \cdot 7H_2O$
3.49	27.61	0.64	45.27	-
	27.50			-

Table 1. Solubility in the Rb₂SO₄-MgSO₄-H₂O system.

Table 3. Solubility in the Rb₂SO₄-NiSO₄-H₂O system.

Liquid phase mass%		Wet solid pl	nase mass%	Composition of the	
Rb ₂ SO ₄	$NiSO_4$	Rb_2SO_4	$NiSO_4$	solid phases	
35.10				Rb ₂ SO ₄	
34.72	1.53	67.42	1.17	-	
34.68	1.57	63.65	18.67	$\begin{array}{c} Rb_2SO_4 + \\ Rb_2Ni(SO_4)_2{\cdot}6H_2O \end{array}$	
31.92	0.50	49.96	25.73	$Rb_2Ni(SO_4)_2{\cdot}6H_2O$	
26.44	0.68	48.18	24.58	-	
19.85	0.81	47.26	26.22	-	
13.00	2.40	43.72	23.19	-	
10.52	3.45	44.36	24.65	-	
7.59	9.04	45.61	26.29	-	
5.08	13.37	46.44	28.07	-	
3.97	19.82	47.56	29.45	-	
4.02	23.37	43.74	29.67	-	
4.06	24.08	15.72	45.67	$\begin{array}{l} NiSO_4{\cdot}7H_2O + \\ Rb_2Ni(SO_4)_2{\cdot}6H_2O \end{array}$	
3.85	24.57	1.43	38.79	NiSO ₄ ·7H ₂ O	
	26.47			-	

Table 2. Solubility in the Rb₂SO₄-CoSO₄-H₂O system.

Liquid phase mass%		Wet solid pl	hase mass%	Composition of the	
Rb_2SO_4	$CoSO_4$	Rb_2SO_4	$CoSO_4$	solid phases	
35.10				Rb ₂ SO ₄	
33.54	1.02	71.49	0.76	-	
33.28	1.09	57.36	21.37	$\begin{array}{c} Rb_2SO_4 + \\ Rb_2Co(SO_4)_2{\cdot}6H_2O \end{array}$	
33.21	1.04	47.24	21.69	$Rb_2Co(SO_4)_2 \cdot 6H_2O$	
27.43	0.98	47.21	23.56	-	
20.68	0.78	45.47	23.38	-	
13.32	4.93	42.78	22.84	-	
8.48	8.59	43.17	25.31	-	
5.44	12.97	44.25	26.83	-	
3.79	17.07	47.36	29.14	-	
2.83	24.21	45.28	30.07	-	
2.40	26.68	14.55	42.73	$\begin{array}{c} CoSO_4{\cdot}7H_2O\\ Rb_2Co(SO_4)_2{\cdot}6H_2O \end{array}$	
2.19	26.49	0.84	44.93	$CoSO_4$ ·7H ₂ O	
	26.14			-	

Table 4. Solubility in the Rb₂SO₄-CuSO₄-H₂O system.

Liquid phase mass%		Wet solid pl	nase mass%	Composition of the	
Rb ₂ SO ₄	CuSO ₄	Rb_2SO_4	CuSO ₄	solid phases	
35.10				Rb ₂ SO ₄	
34.81	0.98	69.11	1.02	-	
33.72	1.04	55.74	20.14	$\begin{array}{c} Rb_2SO_4 + \\ Rb_2Cu(SO_4)_2{}\cdot 6H_2O \end{array}$	
33.64	1.12	50.47	26.73	$Rb_2Cu(SO_4)_2 \cdot 6H_2O$	
28.68	1.06	48.26	26.44	-	
24.59	0.96	45.49	24.36	-	
19.53	0.35	43.72	24.21	-	
14.46	2.48	42.64	23.55	-	
10.72	4.92	44.26	25.87	-	
8.75	9.28	45.82	28.33	-	
4.17	16.16	47.48	29.94	-	
4.43	16.57	10.45	42.48	$\begin{array}{c} CuSO_4{\cdot}5H_2O + \\ Rb_2Cu(SO_4)_2{\cdot}6H_2O \end{array}$	
3.97	17.01	1.74	38.76	$CuSO_4{}^{\cdot}5H_2O$	
	18.30			-	

Liquid phase mass%		Wet solid phase mass%		Composition of the solid		
Rb_2SO_4	$ZnSO_4$	Rb_2SO_4	ZnSO ₄	phases		
35.10				Rb_2SO_4		
34.61	1.98	61.48	1.28	-		
34.15	1.72	56.57	20.46	$Rb_2SO_4 + Rb_2Zn(SO_4)_2 \cdot 6H_2O_4$		
33.85	1.34	50.26	27.19	$Rb_2Zn(SO_4)_2 \cdot 6H_2O$		
25.02	0.50	48.43	27.54	-		
17.72	2.63	43.67	25.71	-		
10.44	6.36	45.72	28.44	-		
4.87	15.06	43.85	29.16	-		
3.77	25.45	45.05	30.11	-		
4.01	30.86	42.67	30.85	-		
3.96	33.75	13.74	44.26	$\frac{ZnSO_4 \cdot 7H2O +}{Rb_2Zn(SO_4)_2 \cdot 6H_2O}$		
3.79	34.05	1.28	29.04	$ZnSO_4 \cdot 7H_2O$		
	36.43			-		

Table 5. Solubility in the Rb₂SO₄-ZnSO₄-H₂O system.

Table 6. Lattice parameters of the Tutton compounds.

Tutton salts	a (Å)	<i>b</i> (Å)	c (Å)	β (°)	$V(Å^3)$
Mg	9.235(2)	12.493(2)	6.227(1)	105.95(2)	690.8(2)
Co	9.198(3)	12.483(5)	6.245(1)	106.12(2)	688.9(2)
Ni	9.147(4)	12.418(4)	6.223(3)	106.06(2)	679.3(3)
Cu	9.269(4)	12.365(4)	6.229(2)	105.38(2)	688.4(3)
Zn	9.198(2)	12.461(4)	6.248(2)	105.9(2)	688.5(1)

3.3. Infrared Spectra of Neat Rubidium Tutton Sulfates

The free tetrahedral ions (XO_4^{n-}) under perfect T_d symmetry exhibit four internal vibrations: $v_1(A_1)$, the symmetric X-O stretching modes, $v_2(E)$, the symmetric XO₄ bending modes, $v_3(F_2)$ and $v_4(F_2)$, the asymmetric stretching and bending modes, respectively. The normal vibrations of the free sulfate ions in aqueous solutions are reported to appear as follows: $v_1 = 983$ cm⁻¹, $v_2 = 450$ cm⁻¹, $v_3 = 1105$ cm⁻¹, $v_4 = 611$ cm⁻¹ [10]. On going into solid state, the normal modes of the (XO_4^{n-}) ions are expected to shift to higher or lower frequencies.

The static field (related to the symmetry of the site on which the XO_4^{n-} ions are situated) will cause a removal of the degeneracy of both the doubly degenerate v_2 modes and the triply degenerate v_3 and v_4 modes. Since the tetrahedral ions in the structures of Tutton compounds occupy site symmetry C_1 , two bands for v_2 (2A)

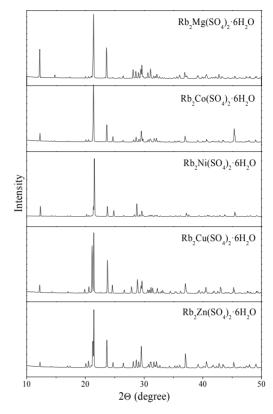


Figure 7. X-ray powder diffraction patterns of Rb₂M(SO₄)₂·6H₂O (M = Mg, Co, Ni, Cu, Zn).

and three bands for v_3 and v_4 (3A), respectively, are expected to appear in the vibrational spectra as predicted from the site group analysis (the nondegenerate v_1 mode is activated). Additionally, the factor group analysis (C_{2h} factor group symmetry) predicts a splitting of each species of A symmetry into $A_g + A_u + B_g + B_u$ (related to interactions of identical oscillators, correlation field effects). Consequently, 18 infrared bands $(9A_u + 9B_u)$ and 18 Raman bands $(9A_g + 9B_g)$ correspond to the normal vibrations of the tetrahedral ions. The correlation diagram between the T_d point symmetry, C_1 site symmetry of the sulfate ions and C_{2h} factor group symmetry is presented in Figure 8. Unit-cell theoretical treatment for the translational lattice modes (Rb, M", SO₄²⁻ and H₂O) and librational lattice modes (SO_4^{2-} and H_2O) yields: 69 modes of Ag, Au, Bg and Bu symmetry and 48 modes of A_g, A_u, B_g and B_u symmetry for the translational and librational modes respectively [11].

Infrared spectroscopic investigations of the potassium Tutton sulfates and selenates are widely discussed in the literature [9,12-17] and those of the rubidium selenates are commented in [8,13,18].

Infrared spectra of $Rb_2M(SO_4)_2 \cdot 6H_2O$ (M = Mg, Co, Ni, Cu, Zn) in the region of 4000 - 400 cm⁻¹ are shown in **Figures 9** and **10**. Some structural and spectroscopic data are summarized in **Table 7** (for comparison the

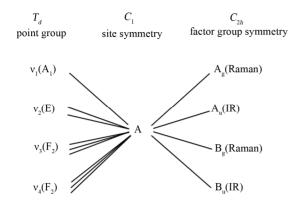


Figure 8. Correlation diagram between T_d point symmetry, C_1 site symmetry and C_{2h} factor group symmetry (SO₄²⁻ ions).

structural and spectroscopic data for the respective potassium and ammonium sulfates are presented; the data are taken from [9]). It is readily seen that the shape of the spectra and the band positions are similar owing to the isostructureness of the double salts. The six infrared bands expected according to the factor group analysis for the asymmetric stretching modes of the sulfate ions coalesce into three bands: $Rb_2Mg(SO_4)_2 \cdot 6H_2O$ (1141, 1111, 1097 cm⁻¹); $Rb_2Co(SO_4)_2 \cdot 6H_2O$ (1142, 1112, 1097 cm⁻¹); $Rb_2Ni(SO_4)_2 \cdot 6H_2O$ (1139, 1109, 1097 cm⁻¹); $Rb_2Zn(SO_4)_2 \cdot 6H_2O$ (1139, 1111, 1099 cm⁻¹), and into

two bands for Rb₂Cu(SO₄)₂· $6H_2O$ (1144, 1097 cm⁻¹). The v₁ modes appear at 984 cm⁻¹. The bending modes v₄ are detected in the spectral intervals of 632 - 612 cm⁻¹. v₂ appear at about 450 cm⁻¹ (see **Figure 9**).

Our infrared spectroscopic findings differ slightly from those reported by Brown and Ross with respect to the number of the bands corresponding to v_3 of the sulfate ions. The spectra of the rubidium compounds commented in [13] show more bands for v_3 (some of them assigned as shoulders) as compared to our results. We believe that the difference between our spectra and those discussed by Brown and Ross is due probably to the temperatures at which the spectra are recorded. It is mentioned in [13] that some spectra are run at liquid nitrogen temperature. However, there is no indication which spectra are obtained at LNT.

The normal vibrations of the water molecules appear in the high frequency region of $3000 - 4500 \text{ cm}^{-1}$. The three crystallographically different water molecules (C_1 site symmetry) in the structures of the Tutton sulfates are expected to display six infrared bands corresponding to the asymmetric and symmetric modes, v_3 and v_1 respectively. However, due to the strong interactions of identical oscillators O-H the different normal modes overlap and as a result one broad band centered at about 3230 cm⁻¹ is observed in the spectra of the double Tutton salts

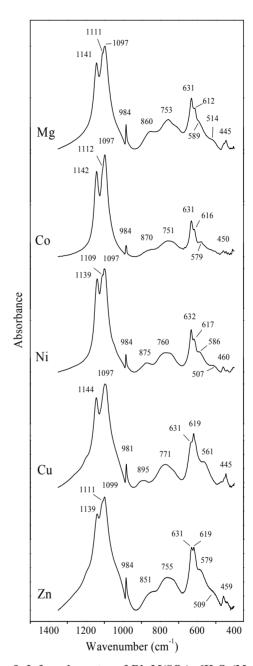


Figure 9. Infrared spectra of $Rb_2M(SO_4)_2 \cdot 6H_2O$ (M = Mg, Co, Ni, Cu, Zn) in the region of 1400 - 400 cm⁻¹ (normal vibrations of the SO_4^{2-} ions and water librations).

(with exception of the copper compound) (**Figure 10**). Three bands corresponding to v_2 of the three crystallographically different water molecules are observed in the spectra of the compounds under study: 1709, 1630 and 1553 cm⁻¹ (magnesium); 1734, 1641 and 1560 cm⁻¹ (cobalt); 1734, 1631 and 1574 cm⁻¹ (nickel); 1734, 1626 and 1590 cm⁻¹ (copper), and 1734, 1621 and 1548 cm⁻¹ (zinc). The band positions of the stretching modes indicate that comparatively strong hydrogen bonds are formed in the sulfates and the hydrogen bond strengths 142

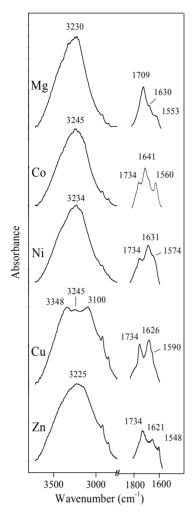


Figure 10. Infrared spectra of $Rb_2M(SO_4)_2 \cdot 6H_2O$ (M = Mg, Co, Ni, Cu, Zn) in the region of the stretching and bending modes of the water molecules.

do not depend on the M" chemical nature. The appearance of a band at a lower frequency (3100 cm⁻¹) in the spectrum of the copper compound shows that stronger hydrogen bonds are formed in this salt as compared to other rubidium compounds. This spectroscopic finding is owing to the stronger *synergetic* effect of the Cu²⁺ ions, *i.e.* to the strong Cu-OH₂ interactions (increasing of the acidity of the water molecules) [19,20]. The formation of comparatively strong hydrogen bonds in the rubidium compounds is due to the strong proton acceptor capacity of the sulfate ions [19,20].

The water librations (rocking, twisting and wagging) appear in the region below 1000 cm^{-1} and a strong overlapping with vibrations of other entities in the structure is expected. Two types of water librations for the Tutton sulfates are discussed briefly in the literature—rocking and wagging, the former observed at higher frequencies [14]. Each type is characterized with two broad bands. The water molecules bonded to the M" ions via shorter

M"-OH2 bonds display water librations at higher frequencies as compared to those forming longer M"-OH2 bonds (equatorial water molecules). The former M"-OH₂ bonds are much more polarized due to the stronger synergetic effect of the M" ions (stronger metal-water interactions). The mean wavenumbers for the rocking librations are reported to have values of 855 and 740 $\rm cm^{-1}$, and 770 and 680 cm^{-1} for the potassium and ammonium sulfates, respectively. The respective wagging modes have mean values of 570 and 441 cm^{-1} for the potassium compounds, and 544 and 425 cm⁻¹ for the ammonium ones [14]. Thus, the bands in the interval of 895 - 751 cm⁻¹ are attributed to the rocking modes of the water molecules and the bands in the region of 589 - 507 cm^{-1} to the wagging modes. The close wavenumbers of the water librations confirm the claim that the hydrogen bonds formed in the rubidium Tutton sulfates are of close strength (see Figure 9).

3.4. Infrared Spectra of SO₄²⁻ Ions Matrix Isolated in Tutton Selenates

The method of crystal matrix-spectroscopy provides important information about the local potential at the lattice sites where the guest ions are located as deduced from their extent of distortion and the chemical nature of the ligand environments in the lattice. When the polyatomic ions are doped in host lattices at low concentration (up to 2 - 7 mol%) the correlation field splitting, the dispersion of phonon curves (due to the interactions between identical oscillators) and LO/TO splitting effects (due to the long-range forces of electrostatic origin) are neglected. Thus, the vibrational spectra of the guest ions are determined by the site symmetry, which is assumed to be the same as that of the respective host ions (substitutional mixed crystals). The spectra of matrix-isolated polyatomic ions are an excellent probe of the local potential at called an *energetic* distortion in order to distinguish it from the geometrical distortion revealed by structural data [21-23]. Both the site group splitting of the asymmetric modes (Δv_{as}) and the value of Δv_{max} (the difference between the highest and the lowest wavenumbered components of the stretching and bending modes, respectively) are an adequate measure for the degree of energetic distortion of the polyatomic ions [22,24,25]. Recently, the value of the ratio $\Delta v_{as}/v_c$ (where v_c is the centro-frequency value of the asymmetric modes) has been proposed to calculate the relative splitting of the dopant ions [26].

Numerous papers in the literature are devoted to the vibrational spectroscopic studies of polyatomic ions (for example, XO_3^- , XO_4^{2-} , XO_4^{3-}) doped in various host lattices [21,23-40]. Lutz *et al.* performed solid solution spectroscopic experiments on monoclinic

Ba(ClO₃)₂·H₂O-type and orthorhombic Sr(ClO₃)₂-type

Compounds	V/n (Å ³)	S-O (Å)	$\Delta r(SO_4)$ (Å)	$v_3 ({\rm cm}^{-1})$	\overline{V}_{3} (cm ⁻¹)	$v_1 (\mathrm{cm}^{-1})$	$\Delta v_3 (\mathrm{cm}^{-1})$	$\Delta v_{\rm max} ({\rm cm}^{-1})$
K ₂ Mg(SO ₄) ₂ ·6H ₂ O	164	1.474	0.065	1147, 1108, 1098	1115	984	49	163
Rb ₂ Mg(SO ₄) ₂ ·6H ₂ O	172	1.470	0.014	1141, 1111, 1097	1116	984	45	160
$(NH_4)_2Mg(SO_4)_2{\cdot}6H_2O$	174	1.473	0.021	1147, 1108, 1098	1118	984	49	163
$K_2Co(SO_4)_2 \cdot 6H_2O$	165	1.474	0.020	1144, 1100	1122	984	44	160
Rb ₂ Co(SO ₄) ₂ ·6H ₂ O	172	1.472	0.015	1142, 1112, 1097	1117	984	45	158
$(NH_4)_2Co(SO_4)_2 \cdot 6H_2O$	175	1.510	0.065	1146, 1102	1124	982	44	164
$K_2Ni(SO_4)_2{\cdot}6H_2O$	162	1.473	0.016	1144, 1111, 1101	1119	982	43	164
Rb ₂ Ni(SO ₄) ₂ ·6H ₂ O	169	1.470	0.012	1139, 1109, 1097	1115	984	42	155
$(NH_4)_2Ni(SO_4)_2\cdot 6H_2O$	171	1.476	0.020	1144, 1108, 1100	1117	982	44	164
$K_2Cu(SO_4)_2 \cdot 6H_2O$	164	1.471	0.024	1144, 1102, 1097	1114	984	47	160
$Rb_2Cu(SO_4)_2{\cdot}6H_2O$	171	1.473	0.024	1144, 1097	1121	981	47	163
$(NH_4)_2Cu(SO_4)_2 \cdot 6H_2O$	173	1.473	0.022	1144, 1102, 1095	1114	981	49	163
$K_2Zn(SO_4)_2 \cdot 6H_2O$	164	1.470	0.018	1141, 1108, 1102	1117	982	39	159
Rb ₂ Zn(SO ₄) ₂ ·6H ₂ O	169	1.472	0.013	1139, 1111, 1099	1116	984	40	155
$(NH_4)_2Zn(SO_4)_2\cdot 6H_2O$	173	1.474	0.021	1144, 1105	1124	984	39	160

Table 7. Some structural and spectroscopic characteristics for the $SO_4^{2^-}$ ions in the neat Tutton salts (*V/n*, unit-cell volumes divided by the numbers of the $SO_4^{2^-}$ ions; S-O, mean values of the S-O bond lengths; $\Delta r(SO_4)$, the difference between the longest and the shortest S-O bond lengths in the respective tetrahedra; Δv_{max} , the difference between the highest and the lowest wave numbered components of the stretches of the $SO_4^{2^-}$ ions; the structural data are taken from [1,2,9]).

host crystals with incorporated ClO_3^- , BrO_3^- and $\text{IO}_3^$ ions [21,24,25]. The energetic distortions of tetrahedral ions matrix-isolated in different alkali metal lattices are widely discussed in [26-34]. Recently, infrared spectra of isomorphously included species— NH_4^+ ions isolated in KMPO₄·6H₂O (M = Mg, Ni) and PO_4^{3-} ions isolated in MgNH₄AsO₄·6H₂O have been reported [37].

Infrared spectra of mixed crystals

 $M'_2M''(\text{SeO}_4)_{2-x}(\text{SO}_4)_x \cdot 6\text{H}_2\text{O}$ (M' = K, Rb, NH₄⁺; M'' = Mg, Co, Ni, Cu, Zn; x is approximately 0.02) are presented in **Figure 11**. The infrared spectroscopic characteristics of the matrix-isolated SO₄²⁻ ions are summarized in **Table 8** (the data for potassium and ammonium Tutton salts are taken from [9]).

The matrix-isolated SO_4^{2-} ions exhibit three bands corresponding to v_3 in agreement with the low site symmetry C_1 of the SeO_4^{2-} host ions. Bands of small intensities around 980 cm⁻¹ appear in the spectra which are assigned to v_1 of the guest sulfate ions (the spectra are recorded at higher concentrations of the samples in KBr in order to distinguish the v_1 mode; the spectra are not shown; see **Table 8**). When the larger SeO_4^{2-} host ions are replaced by the smaller SO_4^{2-} guest ions the mean values of the asymmetric stretching modes v_3 of the SO_4^{2-} guest ions are slightly shifted to lower frequencies as compared to those of the same ions in the neat sulfate compounds due to the smaller repulsion potential of the selenate matrices, *i.e.* to the larger unit-cell volumes of the respective selenates (compare **Tables 7** and **8**).

Several factors are expected to influence on the values of Δv_3 and $\Delta v_3/v_c$: 1) the chemical nature of the metal ions; 2) the repulsion potential of the selenate matrices, and 3) the strength of the hydrogen bonds. The spectroscopic experiments show that the distortion of the SO_4^{2-} guest ions as deduced from the values of Δv_3 and $\Delta v_3/v_c$ increase on going from the potassium to the ammonium compounds (see Table 8). The formation of hydrogen bonds between the SO_4^{2-} guest ions and both the water molecules of the host compounds and the NH₄⁺ host ions obviously facilitate the extent of energetic distortion of the guest ions (for example, Δv_3 of the sulfate guest ions have values of 30 and 51 cm⁻¹ in the nickel potassium and ammonium matrices, and 33 and 49 cm^{-1} in the zinc potassium and ammonium matrices, respectively). The larger values of Δv_3 and $\Delta v_3/v_c$ of the sulfate guest ions matrix-isolated in rubidium selenate matrices as compared to those of the SO₄²⁻ ions included in the potassium matrices are owing probably to the smaller

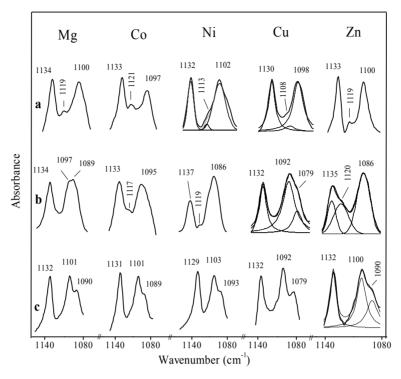


Figure 11. Infrared spectra of SO_4^{2-} ions matrix-isolated (about 2 mol%) in $M'_2M''(SeO_4)_{1.98}(SO_4)_{0.02} \cdot 6H_2O$ (M' = K, Rb, NH₄⁺; M'' = Mg, Co, Ni, Cu, Zn) in the region of v_3 : potassium selenate matrices (row a); ammonium selenate matrices (row b); rubidium selenate matrices (row c).

Table 8. Some spectroscopic characteristics of SO ₄ ²⁻	guest ions matrix-isolated in selenate matrices (for the assignments see
Table 1; the data for the potassium and ammonium c	ompounds are taken from [9]).

		SO ₄ ²⁻ guest	t ions (approximation	ately 2 mol%)		
Host compounds —	$v_3 ({\rm cm}^{-1})$	\overline{v}_{3} (cm ⁻¹)	$v_1 ({\rm cm}^{-1})$	$\Delta v_3 (\mathrm{cm}^{-1})$	$\Delta v_{\rm max} ({\rm cm}^{-1})$	$\Delta v_3/c$ (%)
K ₂ Mg(SeO ₄) ₂ ·6H ₂ O	1134, 1119, 1100	1118	980	34	154	3.04
Rb ₂ Mg(SeO ₄) ₂ ·6H ₂ O	1132, 1101, 1090	1108	981	42	151	3.79
$(NH_4)_2Mg(SeO_4)_2{\cdot}6H_2O$	1134, 1097, 1089,	1107	980	45	154	4.06
K ₂ Co(SeO ₄) ₂ ·6H ₂ O	1133, 1121, 1097	1117	982	36	151	3.22
Rb ₂ Co(SeO ₄) ₂ ·6H ₂ O	1131, 1101,1089	1107	980	42	151	3.79
$(NH_4)_2Co(SeO_4)_2 \cdot 6H_2O$	1133, 1117, 1095	1115	981	38	152	3.40
K ₂ Ni(SeO ₄) ₂ ·6H ₂ O	1132, 1113, 1102	1116	982	30	150	2.68
Rb ₂ Ni(SeO ₄) ₂ ·6H ₂ O	1129, 1103, 1093	1108	982	36	147	3.25
(NH ₄) ₂ Ni(SeO ₄) ₂ ·6H ₂ O	1137, 1119, 1086	1114	979	51	158	4.58
$K_2Cu(SeO_4)_2 \cdot 6H_2O$	1130, 1198, 1098	1112	981	32	149	2.88
$Rb_2Cu(SeO_4)_2 \cdot 6H_2O$	1135, 1100, 1084	1106	980	51	155	4.61
$(NH_4)_2Cu(SeO_4)_2\cdot 6H_2O$	1132, 1092, 1079	1101	981	53	151	4.81
$K_2Zn(SeO_4)_2 \cdot 6H_2O$	1133, 1119, 1100	1117	982	33	151	2.95
Rb ₂ Zn(SeO ₄) ₂ ·6H ₂ O	1132, 1100, 1090	1107	983	42	149	3.79
$(NH_4)_2Zn(SeO_4)_2 \cdot 6H_2O$	1135, 1120, 1086	1114	982	49	153	4.40

local potential at the lattice sites where the sulfate ions are located (*i.e.* to the larger unit-cell volumes of the rubidium selenates) [39].

As far as the influence of the M" ion nature on the values of Δv_3 and $\Delta v_3/v_c$ of the sulfate ions is concerned the spectroscopic experiments show that the sulfate ions are stronger distorted in the case of

 $(NH_4)_2M(SeO_4)_{1.98}(SO_4)_{0.02} \cdot 6H_2O$ (M = Ni, Cu, Zn) and Rb₂Cu(SeO₄)_{1.98}(SO₄)_{0.02} $\cdot 6H_2O$. These findings are due probably to the formation of stronger hydrogen bonds in these compounds, *i.e.* to the stronger interaction between the water molecules of the host compounds and the sulfate guest tetrahedra (stronger proton donor capacity of the water molecules coordinated to the copper, zinc and nickel cations, *i.e.* to the stronger *synergetic* effect of these ions) [19,20].

4. Conclusions

1) The solubility in the three-component systems Rb_2SO_4 -MSO₄-H₂O (M = Mg, Co, Ni, Cu, Zn) was studied at 298 K.

2) Isostructural compounds, $Rb_2M(SO_4)_2 \cdot 6H_2O$ (M = Mg, Co, Ni, Cu, Zn), crystallize from the ternary solutions within wide concentration ranges due to the strong complex formation processes in the solutions.

3) Comparatively strong hydrogen bonds are formed in the rubidium Tutton sulfates as deduced from both the wavenumbers of the stretching modes of the water molecules and the water librations due to the strong proton acceptor strength of the sulfate ions.

4) The degree of energetic distortion of the SO_4^{2-} guest ions in $M'_2M''(SeO_4)_{1.98}(SO_4)_{0.02} \cdot 6H_2O$ ($M' = K_2$, Rb, NH₄; M'' = Mg, Co, Ni, Cu, Zn) is analyzed with respect to the values of Δv_{as} and $\Delta v_{as}/v_c$.

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