

ISSN Online: 2151-4844 ISSN Print: 2151-481X

Effect of Incidence Angle of Magnetic Field on the Performance of a Polycrystalline Silicon Solar Cell under Multispectral Illumination

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How to cite this paper: Sourabié, I., Zerbo, I., Zoungrana, M., Combari, D.U. and Bathiebo, D.J. (2017) Effect of Incidence Angle of Magnetic Field on the Performance of a Polycrystalline Silicon Solar Cell under Multispectral Illumination. *Smart Grid and Renewable Energy*, **8**, 325-335. https://doi.org/10.4236/sgre.2017.810021

Received: September 18, 2017 Accepted: October 28, 2017 Published: October 31, 2017

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Abstract

The aim of this work is to investigate, with a three-dimensional steady-state approach, the effect of the incidence angle of a magnetic field on the performance of a polycrystalline silicon solar cell under multispectral illumination. The magneto-transport and continuity equations of excess minority carriers are solved to find the expression of the density of excess minority carriers and the related electrical parameters, such as the photocurrent density, the photovoltage and the electric power, of a grain of the polycrystalline silicon solar cell. The influence of the incidence angle of the magnetic field on the diffusion coefficient, the short-circuit photocurrent density, the open-circuit photovoltage and the electric power-photovoltage is studied. Then, the curves of the electric power-photovoltage is used to find the maximum electric power allowing to calculate, according to the incidence angle of the magnetic field, the fill factor and the conversion efficiency. The study has shown that the increase of the incidence angle of the magnetic field from 0 rad to $\pi/2$ rad, can reduce the degradation of the performance of solar cells.

Keywords

Bifacial, Conversion Efficiency, Incidence Angle, Magnetic Field, Polycrystalline Silicon Solar Cell

1. Introduction

The performance of photovoltaic modules and cells depends on climatic and seasonal parameters but also on external factors as electric field, magnetic field,

or electromagnetic field.

Concerning the effect of climatic and seasonal parameters on the performance of photovoltaic modules, it has been proved that for PV cells/modules, the open circuit voltage decreases with an increase in the cells temperature but increases with an increase in solar radiation [1] [2]. As for the short-circuit, it increases with an increase in solar irradiation as well as an increase in temperature [1] [2]. Finally, the electrical efficiency and the power output of a PV module decrease with an increasing operating temperature of the PV cells/modules [3]. On the other hand, the maximum power produced by PV modules decreases under partially shaded conditions [4] [5] and when dust particles are deposited on them [6].

Various researchers used both experimental and theoretical methods to study the effect of magnetic field on the properties of solar cells. Betser et al. [7] used an experimental method to measure the mobility of minority carriers in the base of InP/GaInAs heterojunction bipolar transistor under constant magnetic field. The magnetic field was applied perpendicular to the direction of the current flow or parallel to the surface of junction emitter-base of the HBT. Vardanyan et al. [8] proposed a method for measuring all recombination parameters (diffusion length, diffusion coefficient, carriers mobility, back surface recombination velocity) in the base region of bifacial solar cell illuminated by its rear side and under constant magnetic field applied parallel to the surface of the p-n junction. Others researchers Madougou et al. [9], Zouma et al. [10], Zerbo et al. [11] [12], Sané et al. [13], Zoungrana et al. [14] [15] and Combari et al. [16], proposed studying the modelling of the negative effect of magnetic field on the performance of solar cells and photovoltaic modules. In their studies, the magnetic field was applied parallel to the surface of the p-n junction of the solar cells. In all these experimental and theoretical methods used to study the effect of magnetic field on the response of solar cells, the magnetic field was applied parallel to the surface of the p-n junction of the solar cells and it caused a degradation of the performance of solar cells.

As a magnetic field causes a degradation of the performance of solar cells, in this work, we showed how the increase of the incidence angle of a magnetic field can limited the degradation of the performance of a polycrystalline silicon solar cell in a three-dimensional steady-state approach.

2. Theoretical Background

The bifacial polycrystalline silicon solar cell n^+ -p- p^+ , studied, is modelled as a regular array of parallelepipedic grains connected in parallel [17] [18]. Each grain has the same electronic and electrical characteristics as the solar cell and that allowed us to conduct the study on a grain and then extrapolate the results to a solar cell [19]. The isolated grain of polycrystalline silicon solar cell illustrated in **Figure 1** is used in the study with the following assumptions:

• the grain has a square section $(g_x = g_y)$ and a thickness H [17] [18];

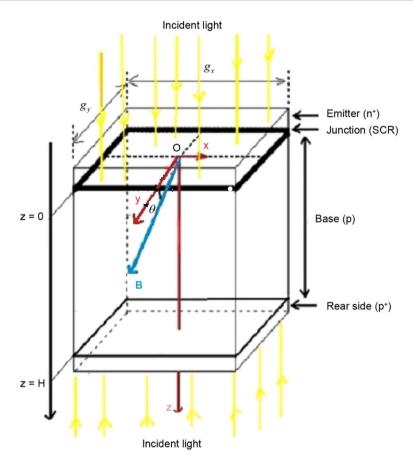


Figure 1. Grain of bifacial polycrystalline silicon solar cell illuminated by multispectral light and under the influence of incidence angle of magnetic field.

- the grain boundaries are perpendicular to the junction and their recombination velocity S_{gb} is constant along the grain boundaries and independent of illumination up to AM 1 [20] so under AM 1.5 which is the standard spectrum for measuring the efficiency of solar cells used on the surface of the earth [21];
- the grain and by extrapolation the polycrystalline silicon solar cell have a high n-type doped emitter so that the contribution on the performance of the cell comes from the p-type base region [18];
- the p-type base is quasi-neutral (quasi-neutral base assumption) so that only the junction electric field will be taken into account [22];
- the effect of temperature is not taken into account on the performance of the solar cell;
- the magnetic field is constant and it is applied into the base region in the plan (O, y, z) with an incidence angle θ with the axis Oy: $\theta = (\widehat{B}, \widehat{e_y})$.

When the grain of the bifacial polycrystalline silicon solar cell represented in **Figure 1** is illuminated simultaneously on both sides with a multispectral light, the continuity equation for the density of excess minority carriers photogenerated in the base is given by Equation (1) [20]:

$$\frac{\partial^{2} \delta(x, y, z)}{\partial x^{2}} + C_{y} \frac{\partial^{2} \delta(x, y, z)}{\partial y^{2}} + \left\{ 1 + \left[\mu_{n} B_{0} \sin(\theta) \right]^{2} \right\} \frac{\partial^{2} \delta(x, y, z)}{\partial z^{2}} + \frac{G(z)}{D_{n}^{*}} - \frac{\delta(x, y, z)}{L_{n}^{*}^{*}} = 0$$
(1)

with
$$C_y = 1 + \left[\mu_n B_0 \cos(\theta)\right]^2$$
, $D_n^* = \frac{D_n}{1 + \left(\mu_n B_0\right)^2}$ and $L_n^* = \frac{L_n}{1 + \left(\mu_n B_0\right)^2}$ are

the electron diffusion coefficient and diffusion length depending on the applied magnetic field.

 $\delta(x, y, z)$ and G(x) are respectively the density of carriers and the optical generation rate of carriers for a multispectral incident light. B_0 and θ are respectively the intensity and the incidence angle of the applied magnetic field while μ_n is the electrons mobility.

The density of the excess minority carriers photogenerated in the base of the polycrystalline silicon grain, general solution of the continuity equation, is expressed by Equation (2) [20]:

$$\delta(x, y, z) = \sum_{j} \sum_{k} \left[A_{j,k} \cdot \cosh\left(\frac{z}{L_{j,k}^{*}}\right) + B_{j,k} \cdot \sinh\left(\frac{z}{L_{j,k}^{*}}\right) + \sum_{i=1}^{3} K_{i} \cdot \left(e^{-b_{i} \cdot z} + e^{-b_{i} \cdot (H-z)}\right) \right]$$

$$\cdot \cos\left(C_{xj}x\right) \cdot \cos\left(C_{yk}y\right)$$

$$(2)$$

with
$$L_{j,k}^{*2}(\theta) = \left[\frac{L_{n}^{*-2} + C_{k}^{2} + C_{j}^{2}}{1 + \left[\mu_{n} B_{0} \sin(\theta)\right]^{2}}\right]^{-\frac{1}{2}}, \quad K_{i} = -\frac{a_{i} \cdot L_{j,k}^{*2}(\theta)}{D_{j,k}^{*}(\theta) \left[b_{i}^{2} \cdot L_{j,k}^{*2}(\theta) - 1\right]},$$

$$\frac{1}{D_{j,k}^{*}(\theta)} = \frac{16 \cdot \sin\left(C_{j} \frac{g_{x}}{2}\right) \cdot \sin\left(C_{k} \cdot \frac{g_{y}}{2}\right)}{D_{n}^{*}(\theta) \cdot \left[\sin\left(C_{j} g_{x}\right) + C_{j} g_{x}\right] \cdot \left[\sin\left(C_{k} \cdot g_{y}\right) + C_{k} \cdot g_{y}\right]} \text{ and}$$

$$D_{n}^{*}(\theta) = D_{n} \frac{1 + \left[\mu_{n} B_{0} \sin(\theta)\right]^{2}}{1 + \left(\mu_{n} B_{k}\right)^{2}}$$

 $D_n^*(\theta)$ is the electron diffusion coefficient depending on the magnetic field and its incidence angle. Coefficients a_i and b_i are tabulated values obtained from modelling of the generation rate considered for over all the solar radiation spectrum under Air Mass 1.5 standard conditions [11] [23]. Coefficients $A_{j,k}$ and $B_{j,k}$ are solved for using the boundary conditions of the solar cell [14] [20] while coefficients C_{xj} and C_{yk} are found using the grain boundary conditions [14] [20].

The expression of the photocurrent density at the junction of the polycrystalline silicon grain is given by Equation (3) [20]:

$$Jph(Sf,B,\theta) = \frac{q \cdot D_n^*(\theta)}{g_x \cdot g_y} \int_{-\frac{g_y}{2} - \frac{g_x}{2}}^{\frac{g_y}{2} - \frac{g_x}{2}} \left[\frac{\partial \delta(x,y,z)}{\partial z} \right]_{z=0} dxdy$$
 (3)

The expression of the photovoltage across the junction of the polycrystalline silicon grain is expressed using Boltzmann's relation [14]:

$$Vph(Sf, B, \theta) = V_T \cdot \ln \left| 1 + \frac{N_B}{n_i^2} \int_{-\frac{g_y}{2}}^{\frac{g_x}{2}} \int_{-\frac{g_x}{2}}^{\frac{g_x}{2}} \delta(x, y, 0) dxdy \right|$$
(4)

 V_T is the thermal voltage, N_B the base doping density and n_i the intrinsic concentration of electrons at thermodynamic equilibrium.

The expression of the electric power provided by the polycrystalline silicon grain to an external circuit is expressed by Equation (5):

$$P(Sf, B, \theta) = Vph(Sf, B, \theta) \cdot Jph(Sf, B, \theta)$$
(5)

The conversion efficiency of the polycrystalline silicon grain is calculated, for various incidence angle of the magnetic field, using Equation (6):

$$\eta = \frac{P_{\text{max}}}{P_{\text{inc}}} \tag{6}$$

The power of the incident light under AM 1.5 standard conditions, used, is $P_{inc} = 100 \text{ mW/cm}^2$.

Knowing the short-circuit photocurrent density, the open-circuit photovoltage and using Equation (7), the fill factor of the polycrystalline silicon grain is calculated according to the incidence angle of the magnetic field:

$$FF = \frac{P_{\text{max}}}{J_{sc} \cdot V_{oc}} \tag{7}$$

3. Results and Discussion

3.1. Effect of the Incidence Angle of the Magnetic Field on the Diffusion Coefficient

The curve of diffusion coefficient versus incidence angle of the magnetic field is plotted in Figure 2.

The curve in **Figure 2** shows that the diffusion coefficient increases for incidence angle varying from $\theta = 0$ rad to $\theta = \pi/2$ rad and from $\theta = \pi$ rad to $\theta = 3\pi/2$ rad while it decreases for incidence angle varying from $\theta = \pi/2$ rad to $\theta = \pi$ rad and from $\theta = 3\pi/2$ rad to $\theta = 2\pi$ rad. In fact, for incidence angle varying from $\theta = 0$ rad to $\theta = \pi/2$ rad and from $\theta = \pi$ rad to $\theta = 3\pi/2$ rad the strength of Lorentz, that deviates the electrons photogenerated in the base of the solar cell, decreases until its cancellation resulting in an increase of the quantity of carriers that diffuse through a unit surface per second. For incidence angle varying from $\theta = \pi/2$ rad to $\theta = \pi$ rad and from $\theta = 3\pi/2$ rad to $\theta = 2\pi$ rad, the reverse behaviour is observed. Thus, the diffusion coefficient versus incidence angle of the magnetic field is a positive periodic sinusoidal function which period is π . Consequently, all electronic and electrical parameters proportional to the diffusion coefficient will have the same behaviour as the diffusion coefficient.

3.2. Effect of the Incidence Angle of the Magnetic Field on the Short-Circuit Photocurrent Density

Figure 3 plots the curve of short-circuit photocurrent density versus incidence angle of the magnetic field.

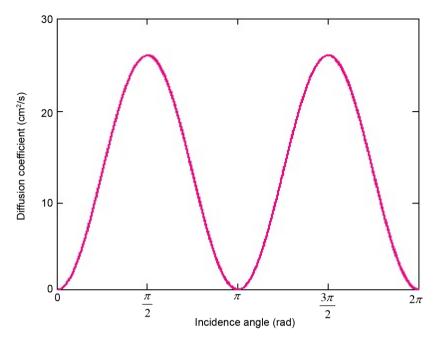


Figure 2. Diffusion coefficient versus incidence angle of magnetic field (D_n = 26 cm²/s; μ_n = 10³ cm²/V.s, B_0 = 7.5 mT).

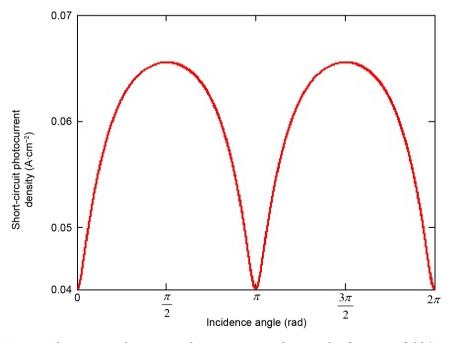


Figure 3. Short-circuit photocurrent density versus incidence angle of magnetic field (D_n = 26 cm²/s; μ_n = 10³ cm²/V.s, B_0 = 7.5 mT, L_n = 0.015 cm; H = 0.03 cm, $g_x = g_y = 3 \times 10^{-3}$ cm; $S_b = 10^3$ cm/s; $S_{gb} = 10^2$ cm/s, $S_b = 10^3$ cm/s, $S_b =$

The curve in Figure 3 has the same behaviour than the one of the diffusion coefficient in Figure 2 because the photocurrent density and therefore the short-circuit photocurrent density are proportional to the diffusion coefficient. So, the strength of Lorentz which influences the carriers' diffusion towards the junction will also influence the short circuit photocurrent. Indeed, for the incidence angles of the applied magnetic field: $\theta = 0$ rad, $\theta = \pi$ rad and $\theta = 2\pi$ rad, the direction of the magnetic field vector is orthogonal to the direction of the velocity vector of the electrons resulting in a maximum effect of the strength of Lorentz. In this situation, the strength of Lorentz deviates the electrons photogenerated resulting in a reduction of the quantity of carriers which can cross the junction and then the short-circuit photocurrent density is minimum. On the other hand, for the incidence angles of the applied magnetic field: $\theta = \pi/2$ rad and $\theta = 3\pi/2$ rad, the direction of the magnetic field vector is collinear to the direction of the velocity vector of the electrons. For this orientation, the strength of Lorentz is null and so it does not has any effect on the movement of the electrons towards the junction. The maximum of carriers can reach to the junction and cross it and then the short-circuit photocurrent density is maximum.

3.3. Effect of Incidence Angle of the Magnetic Field on the Open-Circuit Photovoltage

The curve of open-circuit photovoltage is plotted versus incidence angle of the magnetic field in **Figure 4**.

The curve in **Figure 4** shows that the open-circuit photovoltage decreases for values of the incidence angle θ of the magnetic field included between the intervals $[0, \pi/2]$ and $[\pi, 3\pi/2]$, while it increases for incidence angle θ included between

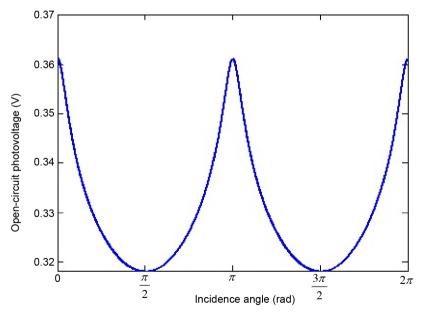


Figure 4. Open-circuit photovoltage versus incidence angle of magnetic field ($D_n = 26 \text{ cm}^2/\text{s}$; $\mu_n = 10^3 \text{ cm}^2/\text{V.s}$, $B_0 = 7.5 \text{ mT}$, $L_n = 0.015 \text{ cm}$; H = 0.03 cm, $g_x = g_y = 3 \times 10^{-3} \text{ cm}$; $S_b = 10^3 \text{ cm/s}$; $S_{gb} = 10^2 \text{ cm/s}$, $S_f = 0 \text{ cm/s}$).

the intervals $[\pi/2, \pi]$ and $[3\pi/2, 2\pi]$. The maximum values of the open-circuit photovoltage are observed at the values of the incidence angle θ equal to 0 rad, π rad and 2π rad and the minimum one's at the incidence angle θ equal to $\pi/2$ rad and $3\pi/2$ rad. It was proved in previous studies [11] [15] that, the increase of the magnetic field leads to an increase of carriers' storage near the junction. Thus, the increase of the incidence angle θ of the magnetic field from $\theta = 0$ rad to $\theta =$ $\pi/2$ rad and from $\theta = \pi$ rad to $\theta = 3\pi/2$ rad which means a reduction of the magnetic field effect on the carriers photogenerated in the base leads then to the decrease of the carriers concentration in the base. The open-circuit photovoltage being a function of carriers concentration near the junction, the increase of the incidence angle θ of the magnetic field leads to the decrease of the open circuit photovoltage in the intervals $[0, \pi/2]$ and $[\pi, 3\pi/2]$ while it increases in the intervals $[\pi/2, \pi]$ and $[3\pi/2, 2\pi]$. The maximum values of the open-circuit photovoltage are observed at the maximum influence of magnetic field which corresponds to the incidence angle θ equal to 0 rad, π rad and 2π rad and the minimum one's at the incidence angle θ equal to $\pi/2$ rad and $3\pi/2$ rad.

The curve in **Figure 4** has the reverse behaviour of the curve of diffusion coefficient in **Figure 2** and short-circuit photocurrent density in **Figure 3**. Despite this reverse behaviour, the open-circuit photovoltage versus incidence angle of the magnetic field is a positive periodic sinusoidal function which period is π .

3.4. Effect of the Incidence Angle of the Magnetic Field on the Electric Power-Photovoltage Characteristics

The curves of electric power versus photovoltage for various incidence angle of the magnetic field are plotted in **Figure 5**.

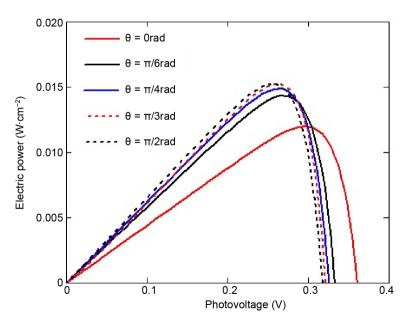


Figure 5. Electric power-photovoltage curves for various incidence angle of magnetic field ($D_n = 26 \text{ cm}^2/\text{s}$; $\mu_n = 10^3 \text{ cm}^2/\text{V.s}$, $B_0 = 7.5 \text{ mT}$, $L_n = 0.015 \text{ cm}$; H = 0.03 cm, $g_x = g_y = 3 \times 10^{-3} \text{ cm}$; $S_b = 10^3 \text{ cm/s}$; $S_{gb} = 10^2 \text{ cm/s}$).

Curves in **Figure 5** have the same shape than those of previous works [12] [16]. We observe on the curves in **Figure 5** that for the values of the incidence angle of the applied magnetic field varying from $\theta = 0$ rad to $\theta = \pi/2$ rad, the maximum electric power increases with the increase of incidence angle while the open-circuit photovoltage decreases. This behaviour is the consequence of the decrease of the strength of Lorentz until its cancellation.

3.5. Electrical Parameters of the Polycrystalline Silicon Grain

As the diffusion coefficient, the short-circuit photocurrent density and the open-circuit photovoltage are symmetrical compared with $\theta=\pi/2$ rad, consequently the modelling of the effect of the incidence angle of the magnetic field on the electrical parameters of the polycrystalline silicon grain will be conducted in the range of $\theta=0$ rad to $\theta=\pi/2$ rad.

By means of the curves of **Figures 3-5** we find successively, for various incidence angle of the magnetic field, the short-circuit photocurrent density, the open-circuit photovoltage and the maximum electric power. The conversion efficiency and the fill factor of the polycrystalline silicon grain are calculated using respectively Equations (6) and (7). The electrical parameters of the polycrystalline silicon grain for various incidence angle of the magnetic field are given in **Table 1**.

These results show that the maximum electric power and consequently the conversion efficiency increase with the increase of the incidence angle of the magnetic field while the fill factor decreases. The increase of the maximum electric power and the conversion efficiency of the polycrystalline silicon grain with the increase of the incidence angle of the magnetic field shows a way to decrease the negative effect of a magnetic field on solar cells. One must tilt solar cells in order that the lines of magnetic field make an incidence angle non-null with the junction of these solar cells.

4. Conclusion

Three-dimensional modelling of the effect of incidence angle of a magnetic field on the performance of a grain of polycrystalline silicon solar cell has been investigated. This study showed that the maximum electric power and the conversion

Table 1. Electrical parameters of the polycrystalline silicon grain according to the incidence angle of magnetic field.

θ (rad)	0	π/6	π/4	π/3	π/2
J_{sc} (mA)	44.074	58.238	62.082	64.196	65.536
V_{oc} (mV)	361.060	330.400	326.000	321.400	317.89
$P_{ m max}~({ m mW\cdot cm^{-2}})$	11.969	14.350	14.884	15.118	15.215
η (%)	11.969	14.350	14.884	15.118	15.215
FF (%)	75.213	74.577	73.542	73.272	73.032

efficiency of the polycrystalline silicon grain increase with the increase of the incidence angle of the magnetic field. It appears also through this study that the degradation of the solar cells performance induced by the application of a magnetic field can be reduced by increasing the incidence angle of the magnetic field. These results show also that the negative effect of a magnetic field on solar cells can be limited by tilting the cells so that the lines of magnetic field make an incidence angle non-null with the junction of the cells.

Acknowledgements

The authors are grateful to International Science Program (ISP) for supporting their research group (energy and environment) and allowing them to conduct this work.

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