

# **Minimum Mass of a Composite Boson**

## **Bo Lehnert**

Alfvén Laboratory, Royal Institute of Technology, Stockholm, Sweden Email: bo.lehnert@ee.kth.se

Received 19 October 2015; accepted 20 November 2015; published 23 November 2015

Copyright © 2015 by author and Scientific Research Publishing Inc. This work is licensed under the Creative Commons Attribution International License (CC BY). http://creativecommons.org/licenses/by/4.0/

## Abstract

A model of the Z boson is elaborated from a revised quantum electrodynamic theory (RQED) by the author. The electromagnetic steady field is derived from a separable generating function with a convergent radial part, resulting in a vanishing net electric charge and a nonzero spin and rest mass. From the superposition of the solutions of two Z bosons with antiparallel spin directions, a model is further formed of a composite boson, the computed mass  $m_c$  of which becomes connected with the mass of 91 GeV for each Z boson. This results in a composite boson which is likely to become identical with the heavy particle recently detected at CERN. Both these particles are thus lacking of net electric charge, magnetic field and spin, are purely electrostatic and highly unstable, and have masses close to the value of 125 GeV.

# **Keywords**

Quantum Electrodynamics, Zero Point Energy, Standard Model and Beyond

# **1. Introduction**

The heavy and unstable particle being recently detected experimentally at CERN [1] [2] has no electric charge, no spin, and a rest mass of 125 GeV. Even if the CERN result is generally being considered as a confirmation of the particle earlier proposed by Higgs [3], the value of its mass cannot be determined by the theory of Higgs. Quigg [4] has further pointed out that such a particle is perhaps not a truly fundamental one, but is built out of as yet unobserved constituents to form a composite particle.

Recently the author has proposed [5] [6] a composite boson to be formed from the superposition of two Z boson solutions with opposite spin directions. The resulting particle then has basic properties in common with the CERN particle, by having a vanishing spin and magnetic field and becoming purely electrostatic and unstable.

As based on a revised quantum electrodynamic theory (RQED) by the author [7]-[9], a model of the Z boson has further been developed which results in a relation between its mass and characteristic radial dimension. With

a mass of 91 GeV, this results in a characteristic radius of about  $10^{-18}$  m, in agreement with that estimated by Quigg [4]. This model will be used in the present investigation to determine the distribution of electrostatic and magnetostatic energy of the Z boson model. In its turn, this also results in a relation between the mass of the Z boson and that of the composite particle, as demonstrated by the following analysis. Such a mass relation becomes a function of the distribution of energy density within the Z boson. As seen from the analysis, a variation of the included parameters becomes associated with a minimum of the composite particle mass.

#### 2. A Model of the Z Boson

A characteristic feature of RQED theory, not being available from conventional theory on the vacuum state, is the existence of steady electromagnetic states, leading to models for massive particles at rest. The corresponding potentials can then be derived from a generating function [7]. Such a radially convergent function results in particle models having vanishing net electric charge but nonzero local and intrinsic electric charges of both polarities, a nonzero spin, as well as a nonzero rest mass.

#### 2.1. Basic Relations

The present analysis starts from a separable and axisymmetric generating function

$$F = G_0 G \qquad G = R(\rho) \cdot T(\theta) \tag{1}$$

in spherical coordinates  $(r, \theta, \varphi)$ . Here *G* is a dimensionless function,  $G_0$  is a characteristic amplitude, and  $\rho = r/r_0$  a normalized radial coordinate with  $r_0$  standing for a characteristic radial dimension. There is an electrostatic potential  $\phi$  and a magnetic vector potential  $\mathbf{A} = (0, 0, A)$ , being sources of the electric and magnetic field strengths  $\mathbf{E} = -\nabla \phi$  and  $\mathbf{B} = \operatorname{curl} \mathbf{A}$ . The vacuum state [7]-[9] further includes a nonzero electric charge density  $\overline{\rho} = \varepsilon_0 \operatorname{div} \mathbf{E}$ , and an electric current density  $\mathbf{j} = (0, 0, C\overline{\rho})$  where  $C = \pm c$  and *c* is the velocity constant of light. The general self-consistent solutions for the electromagnetic components in RQED theory [7] are then obtained from a generating function which takes the form

ł

$$F = CA - \phi \tag{2}$$

and the electrostatic and magnetostatic potentials are determined by

$$CA = -(\sin\theta)^2 DF = F + \phi.$$
(3)

Here the operator  $D = D_{\rho} + D_{\theta}$  has the parts

$$D_{\rho} = -\frac{\partial}{\partial \rho} \left( \rho^2 \frac{\partial}{\partial \rho} \right) \qquad D_{\theta} = -\frac{\partial^2}{\partial \theta^2} - \frac{\cos \theta}{\sin \theta} \frac{\partial}{\partial \theta}$$
(4)

and the potentials are given by

$$\phi/G_0 = -RT - \left(\sin\theta\right)^2 T \left(D_\rho R\right) - \left(\sin\theta\right)^2 \left(D_\theta T\right) R \tag{5}$$

$$CA/G_0 = -(\sin\theta)^2 T(D_\rho R) - (\sin\theta)^2 (D_\theta T)R$$
(6)

in terms of the parts R and T of the generating function.

#### 2.2. The Field Strengths

From expressions (3)-(6) the components of the field strengths are now determined by

$$e_r = r_0 E_r / G_0 = \frac{\partial R}{\partial \rho} T + \left(\sin\theta\right)^2 T \frac{\partial}{\partial \rho} \left(D_\rho R\right) + \left(\sin\theta\right)^2 \left(D_\theta T\right) \frac{\partial R}{\partial \rho}$$
(7)

$$e_{\theta} = r_0 E_{\theta} / G_0 = \frac{R}{\rho} \frac{\partial T}{\partial \theta} + \frac{\partial}{\partial \theta} \Big[ (\sin \theta)^2 T \Big] \frac{1}{\rho} \Big( D_{\rho} R \Big) + \frac{\partial}{\partial \theta} \Big[ (\sin \theta)^2 \big( D_{\theta} T \big) \Big] \frac{R}{\rho}$$
(8)

$$cb_{r} = r_{0}cB_{r}/G_{0} = -\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\Big[\left(\sin\theta\right)^{3}T\Big]\frac{1}{\rho}\Big(D_{\rho}R\Big) - \frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\Big[\left(\sin\theta\right)^{3}\Big(D_{\theta}T\Big)\Big]\frac{R}{\rho}$$
(9)

$$cb_{\theta} = r_0 cB_{\theta} / G_0 = \frac{1}{\rho} \frac{\partial}{\partial \rho} \Big[ \rho \Big( D_{\rho} R \Big) \Big] (\sin \theta)^2 T + \frac{1}{\rho} \frac{\partial}{\partial \rho} \Big( \rho R \Big) (\sin \theta)^2 \Big( D_{\theta} T \Big).$$
(10)

With  $c^2 = 1/\mu_0 \varepsilon_0$  the energy density becomes

$$v_f = \frac{1}{2} \varepsilon_0 \left( \boldsymbol{E}^2 + c^2 \boldsymbol{B}^2 \right) \tag{11}$$

corresponding to a mass density  $w_f/c^2$ . The total rest mass of the Z boson is then

$$m_Z = \left(2/c^2\right) \int_0^{\infty \pi/2} w_f dV \qquad dV = 2\pi r^2 \left(\sin\theta\right) d\theta dr.$$
 (12)

#### 2.3. The Energy Distributions

A radially convergent generating function of the form

$$R = \rho^{\gamma} e^{-\rho} \qquad T = \left(\sin\theta\right)^{\alpha} \tag{13}$$

is now introduced [7] which includes the two parameters  $\gamma \ge 0$  and  $\alpha \ge 0$ . The radial part *R* is finite at r = 0, approaches zero at large *r*, and has its maximum

$$R_m = \gamma^{\gamma} \mathrm{e}^{-\gamma} \tag{14}$$

at the normalized radius  $\rho_m = \gamma$ , *i.e.* at  $r = r_m = \gamma r_0$ . To obtain a measure of the relative spatial extension of *R* near the maximum  $R_m$ , we define the value of *R* at the fraction  $f_m < 1$  of the radius  $\rho_m$ , *i.e.* 

$$R(f_m\gamma) = (f_m\gamma)^{\gamma} e^{-f_m\gamma}.$$
(15)

The relative extension then becomes

$$g_m = R(f_m \gamma) / R_m = f_m^{\gamma} e^{\gamma (1 - f_m)}.$$
(16)

For a fixed value of  $f_m$  the value of  $g_m$  then decreases at an increasing  $\gamma$ , *i.e.* for an *R* representing an ever decreasing thickness of a shell localized around the maximum radius  $\rho_m$ .

In its turn the polar part T becomes more concentrated to the equatorial plane at  $\theta = \pi/2$  for increasing values of  $\alpha$ . At large values of both  $\gamma$  and  $\alpha$ , the generating function RT thus tends to a spatial distribution being mainly concentrated to a thin ring at the equatorial plane.

With the generating function (13) the normalized field strengths (7)-(10) can be written as

$$e_{r}\rho^{-\gamma}e^{\rho} = (\gamma - \rho)(\sin\theta)^{\alpha} + \left[-\gamma^{2}(\gamma + 1) + (3\gamma^{2} + 5\gamma + 2)\rho - (3\gamma + 4)\rho^{2} + \rho^{3}\right](\sin\theta)^{\alpha+2} + (\gamma - \rho)\left[-\alpha^{2}(\sin\theta)^{\alpha} + \alpha(\alpha + 1)(\sin\theta)^{\alpha+2}\right]$$
(17)

$$e_{\theta}\rho^{-\gamma}e^{\rho} = \alpha \left(\sin\theta\right)^{\alpha-1} \left(\cos\theta\right) + \left[-\gamma(\gamma+1) + 2(\gamma+1)\rho - \rho^{2}\right] (\alpha+2) \left(\sin\theta\right)^{\alpha+1} \left(\cos\theta\right) + \left[-\alpha^{3} \left(\sin\theta\right)^{\alpha-1} + \alpha(\alpha+1)(\alpha+2) \left(\sin\theta\right)^{\alpha+1}\right] (\cos\theta)$$
(18)

$$cb_{r}\rho^{-\gamma}e^{\rho} = \left[-\gamma(\gamma+1)+2(\gamma+1)\rho-\rho^{2}\right](\alpha+3)(\sin\theta)^{\alpha+1}(\cos\theta) + \left[-\alpha^{2}(\alpha+1)(\sin\theta)^{\alpha-1}+\alpha(\alpha+1)(\alpha+3)(\sin\theta)^{\alpha+1}\right](\cos\theta)$$
(19)

$$cb_{\theta}\rho^{-\gamma}e^{\rho} = \left[\gamma(\gamma+1)^{2} - (\gamma+1)(3\gamma+4)\rho + (3\gamma+5)\rho^{2} - \rho^{3}\right](\sin\theta)^{\alpha+2} + \left[\rho - (\gamma+1)\right]\left[-\alpha^{2}(\sin\theta)^{\alpha} + \alpha(\alpha+1)(\sin\theta)^{\alpha+2}\right].$$
(20)

The mass of Equation (12) is now distributed among the four field components of expressions (7)-(10) and (17)-(20) as given by

B. Lehnert

$$2\pi\varepsilon_0 r_0 \left(G_0/c\right)^2 \int_0^{\infty\pi/2} \left(e_r^2, e_\theta^2, c^2 b_r^2, c^2 b_\theta^2\right) \rho^{2\gamma} e^{-2\rho} \left(\sin\theta\right) d\theta d\rho$$
  
=  $2\pi\varepsilon_0 r_0 \left(G_0/c\right)^2 \left(M_{er}, M_{e\theta}, M_{br}, M_{b\theta}\right)$  (21)

where  $(M_{er}, M_{e\theta}, M_{br}, M_{b\theta})$  are the corresponding normalized partial masses.

The distributions of each of these masses can be demonstrated in a two-dimensional space defined by the parameters  $\gamma$  and  $\alpha$ , where  $\gamma = \text{const.}$  and  $\alpha = \text{const.}$  represent nested surfaces. The latter surfaces have then different geometries for each partial mass. This also applies to a comparison between the electrostatic mass  $M_E = M_{er} + M_{e\theta}$  and the magnetostatic mass  $M_B = M_{br} + M_{b\theta}$ .

#### 3. The Composite Boson

#### **3.1. Basic Relations**

A superposition is now made of two Z bosons having the same electrostatic potentials given by Equation (5) and opposite cancelling magnetostatic potentials due to Equation (6) with  $C = \pm c$ . This results in a purely electrostatic composite boson, having an electric field strength of double the value given by Equations (7) and (8), and with no total magnetic field and spin. The resulting particle is expected to become highly unstable, due to its lack of a counter-balancing force due to the magnetic field [7].

With the definitions (21) the normalized mass of the Z boson becomes  $M_Z = M_{er} + M_{e\theta} + M_{br} + M_{b\theta}$ . For a doubled electric field strength due to the superposition, the electrostatic mass of the composite particle becomes  $M_C = 4(M_{er} + M_{e\theta}) = 4M_E$ . Using the experimental value of  $m_Z = 91 \text{ GeV}$  for the Z boson, the corresponding value of the composite particle becomes

$$m_{C} = 4m_{Z}M_{E}/(M_{E} + M_{B}) = 364/|1 + (M_{B}/M_{E})| \text{GeV}.$$
(22)

Here the ratio  $M_B/B_E$  as well as the mass  $m_C$  does not become represented by a system of nested surfaces. Instead the surfaces of constant  $\alpha$  will at some points intersect in a representation where  $m_C$  is plotted as a function of  $\gamma$ .

#### 3.2. Computed Mass Distributions

The partial masses (21) have been computed from Equations (17)-(21) for various values of  $\gamma$  and  $\alpha$ . In most cases  $M_{er}$  and  $M_{b\theta}$  are found to be dominating as compared to  $M_{e\theta}$  and  $M_{br}$ . The electrostatic mass  $M_E$  further differs generally from the magnetostatic mass  $M_B$ , and there is no equipartition between electrostatic and magnetostatic energy. This affects the computed value of the composite particle mass.

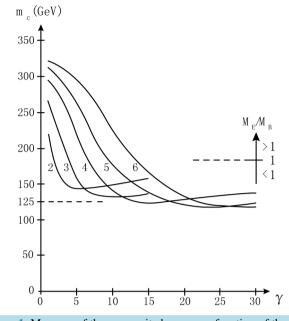
In the experimental investigations on elementary particles of heavy mass, such as those at CERN [1] [2], an increasing available energy corresponds to a situation in which heavy particle masses are approached from below. Therefore a minimum value of the computed mass  $m_c$  will be of main interest. Obtained values of  $m_c$  as functions of the radial parameter  $\gamma$  for a set of values of the polar parameter  $\alpha$  are demonstrated in Figure 1. Here three domains should be considered [1]:

1) For  $0 < \gamma < 7$  and  $0 < \alpha < 2$  the computed values of  $m_c$  are higher than the experimentally detected level of 125 GeV at CERN.

2) In the integral of expression (21) the square of the generating function *RT* of Equation (13) is broadly speaking included. This implies that the domain defined by  $\gamma > 30$  and  $\alpha > 7$  will be represented by thin ring-shaped configurations localized close to the equatorial plane. This is confirmed by an example applied to Equation (16) with  $f_m = 0.8$ , resulting in a relative extension of  $g_m = 0.25$ . Such a geometry is far from that of bulky particles and may be put into doubt.

3) Between the two domains 1) and 2) there is a broad window in parameter space given by the range  $7 < \gamma < 30$ . It represents a high probability of occurrence, also for various forms of geometry as determined by  $\gamma$  and  $\alpha$ . This window is in **Figure 1** seen to include minimum values of  $m_c$ , all being close to the mass of 125 GeV of the particle detected at CERN. The average deviation from this value is only  $\pm 5$  per cent.

The result of **Figure 1** can be analyzed as follows. Each curve in the figure represents the values 2 to 6 given to the polar parameter  $\alpha$  of the part *T* in relations (13), this as a function of the radial parameter  $\gamma$  of the part *R*. The set of curves thus shows that there is generally a minimum of the composite boson mass being close to



**Figure 1.** Mass  $m_C$  of the composite boson as a function of the radial parameter  $\gamma$  for some values (2, 3, 4, 5, 6) of the polar parameter  $\alpha$ . The dashed line at 125 GeV in the left-hand part represents the particle detected at CERN. The dashed line at 182 GeV in the right-hand part represents the situation  $M_E = M_B$  of equipartition for the Z boson.

the limit 125 GeV for the manifold of particle geometries determined by varying values of  $\alpha$  and  $\gamma$ . Further, when performing experiments at increasing available energies, this minimum level will first be approached from below. This implies that a composite boson can and will be created when reaching the same level, before passing to higher levels of available energy.

#### **4.** Conclusions

From superposition of the solutions for two Z bosons with antiparallel spin directions, a model of a composite boson has been formed in terms of the present RQED theory. It connects the computed mass  $m_c$  of the composite boson with the mass of 91 GeV given for the Z boson. Within a large window of the prevailing parameter space, the composite boson mass is found to have a minimum close to the mass of 125 GeV found for the heavy particle detected at CERN.

Due to these results, the composite boson of the present theory is thus likely to become identical with the heavy particle found at CERN. Both particles are namely lacking of net electric charge, magnetic field and spin, are purely electrostatic and highly unstable, and have rest masses close to 125 GeV. The present result has no relation to the theory by Higgs.

### Acknowledgements

The author is indebted to MSc Yushan Zhou for a valuable work with the computations of the present analysis.

#### References

- [1] Aad, G., et al. (2012) Physics Letters, B716, 1-29. http://dx.doi.org/10.1016/j.physletb.2012.08.020
- [2] Chatrchyan, S., et al. (2012) Physics Letters, B716, 30-61. <u>http://dx.doi.org/10.1016/j.physletb.2012.08.021</u>
- [3] Higgs, P.W. (1966) *Physical Review*, **154**, 1156-1168. <u>http://dx.doi.org/10.1103/PhysRev.145.1156</u>
- [4] Quigg, C. (2008) Scientific American, 298, 46-53. http://dx.doi.org/10.1038/scientificamerican0208-46
- [5] Lehnert, B. (2013) Progress in Physics, 3, 31-32.
- [6] Lehnert, B. (2014) Progress in Physics, 10, 5-7.

- [7] Lehnert, B. (2013) Revised Quantum Electrodynamics. Nova Science Publications, Inc., New York.
- [8] Lehnert, B. (2014) Journal of Electromagnetic Analysis and Applications, **6**, 319-327. http://dx.doi.org/10.4236/jemaa.2014.610032
- [9] Lehnert, B. (2015) Journal of Modern Physics, 6, 448-452. <u>http://dx.doi.org/10.4236/jmp.2015.64048</u>