# Reactions on Rigid Legs of Rectangular Tables 

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#### Abstract

A weight is placed on the top of a rectangular rigid ideal table with four legs, each leg is placed at each vertex of the rectangular table. It is assumed that the legs do not bend when the weight is added. The reactions are computed by assuming the table is supported on a beam, introducing two new beam parameters and minimizing a deflection function of the new parameters. A physical experiment is performed in the lab and the reactions on each leg are provided. The experimental results match the theoretical ones obtained by the proposed model. Geometrical interpretations of the results are given.


## Keywords

Reaction Forces, Rectangular Table, Deflection, Rigid Table, Indeterminate Case, Torque

## 1. Introduction

Euler [1] set a problem in St. Petersburg Academy of Sciences in 1773 in his dissertation De pressione ponderis in planum cui incumbit. The problem consisted of placing a weight on the surface of a rectangular table supported on four legs and determining the four reaction forces of the legs. Euler claimed that the solution to this problem "is not only much harder but also doubtful and misleading". He said that the imperfections of the feet might force the body to rest on only three legs and the pressure on the fourth one may be null. In Section 3, we observe that the problem arises because the system of equations given by the mechanics of static equilibrium is a linear system of equations that consists of three equations and four unknowns, being these four unknowns precisely the reaction forces at the legs of the table. This linear system of equations is called an indeterminate system.

Euler then changed the problem: the support is no longer on a rigid plane but on a soil. Euler argued that to understand the problem on a rigid body, one must pose the problem on a soft body, and that once this problem has been understood there, the solution from one body to another can be carried out by a mathematical calculus argument (a limit argument), "in the case of a soft ground where four feet penetrate, one must assume that the pressures brought about by the feet are proportional to those movements". This is the principle where Euler based his arguments. He stressed that even though here a soft ground is considered, this phenomena is independent of the ground because the movements, that were hypothetically introduced, "happen with the support of our imagination".

After Euler posed and discussed the problem, many people have studied it in depth, each contributing at some extent towards a richer interpretation of the problem. These people [2] include Jean Baptiste Lerond d'Alambert in 1780 who dedicated to the support problems part of a memoir entitled Sur quelques questions de Méchanique, Abbot Charles Bossut in 1788 who gives an interpretation of a triangular table in terms of two levers, Paolo Delanges of Brescia in 1790 who three times in 1790,1798 and 1811 made a proposal for an experimental investigation, Pietro Paoli in 1792 who checked the validity of the computations done by others regarding the problem, Anton Maria Lorogna in 1794 who attempted to formulate a "reasonable" principle that can not be obtained from the principle of mechanics, Mariano Fontana [3] in 1798 who provides a formula to obtain the reactions on a parallelogramic table by imagining that the weight is transmitted to the sides of the parallelogram, later Fontana writes "I started to meditate more carefully about this problem and I have found that, the solution to the problem is not determined only by the principles of mechanics". We will refer to Fontana later in this paper.

Other people that also worked on this problem (see [2]) were Guiseppe Venturoli in 1806, Gabrio Piola in 1824, Barsotti in 1842 and Ottaviano Fabrizio Mossotti in 1858; Malfatti in 1805 published three memoirs, in one of them he stresses that if the object is placed at the center of a regular polygon, all the reactions at each of the vertices of the polygon will be equal and hence they can not be "assigned arbitrarily", hence in this case there is no statics indetermination. It was finally Navier [4] in his Résumé des leçons données à l’Ecole royale des ponts et chaussées in 1825 who recognized that the indeterminate system will no longer be indeterminate if one takes into account the "elasticity of the reacting bodies". Even though the solution to the problem of the rectangular table appears for the first time in Castigliano's thesis in 1873, Augustin Cournot in 1828 is the one who describes an extrema point in the pressure involved in the problem, he talks about a minimum for the pressure exerted in the body. A derivation of the solution can be appreciated in Du Bois 1902 [5], he considers that each foot of the table is hence a spring obeying Hooks’ Law and the work expressed in terms of the potential energy is computed. Using The Principle of Least Work, a fourth equation is obtained by obtaining the derivative of the total potential energy and setting it to zero. Therefore we now have a system of four independent equations and four unknowns which can be solved in this case providing a unique solution.

There is a list of people who contributed to the solution of the problem during the 1830-1873 period, the list includes Moseley, Pagani and Menabrea, Massotti and Dorna among others. For a more detailed history of the problem you can read [2].

The paper by Benvenuto [2] contains not only more details about the history of the problem itself but also gives a geometrical interpretation of the triangular table solution (as the one we have in this paper), it provides a short discussion and a geometrical interpretation of a solution given by Fontana to the general parallelogramic case (hence the rectangular case is included) which we will address in Section 5.

Recently, three NASA engineers Surya N. Patnaik, Dale A. Hopkins and Gary R. Halford published a paper in 2004 [6] where they discussed the problem and introduce a compatibility condition (a fourth equation) that they obtained based on the deformation displacements of the legs, these deformations are written in terms of the angles at which the table tilts. When this compatibility condition is expressed in terms of the forces, it is then concluded that the sum of the reactions of two opposite corners of the table matches the sum of the reactions of the other two opposite corners. After all, the solution to the new system of four equations with four unknowns that Patnaik, Hopkins and Halford obtained agrees with the one described in Du Bois [5].

One of our main hypotheses (and here we differ from the results obtained by all the previous people) is that the legs of the table are completely rigid, i.e. they do not deflect or deform. The way we approach the problem is by considering the fact that the reactions on the legs are known and given in terms of two parameters. These parameters are obtained by first assuming that the four legs are replaced by a beam with end points at opposite edges of the table, then changing roles with the other two sides of the table. One important and natural assump-
tion is a consistency condition on the forces in terms of these two parameters. We finally give a condition (and here is where our main contribution comes in) on the deflection of the table (not of the legs). This process allows us to minimize such deflections obtaining the values of the parameters we started with initially, providing therefore one more equation which, together with the three that one has by considering mechanical assumptions, one is able to find a solution for the table of rigid legs. In a forthcoming paper, we will generalize these results to the quadrilateral table and we are also considering increasing the number of legs on a future paper.

Regarding the structure of the paper, in Section 2, we discuss the Beam and the Triangular case, in each case we solve the system and we provide a geometrical interpretation that matches the ones given in Benvenuto [2]. In Section 3, we carefully analyze and discuss the indeterminate case of a rectangular table. Right after, in Section 4, we establish the main hypotheses under which we will solve the rectangular table case. It is precisely in Section 5 where we give conditions that will help us to solve the problem.

Finally, in Section 6 we provide some lab results that are consistent with the theoretical values obtained by the proposed theoretical model. A discussion on how these lab values are obtained is contained in this section too.

## 2. Beam and Triangular Tables

In this section we will discuss the beam (one dimensional table) and the triangular table case.

### 2.1. Beam Case

A load $W$ of magnitude $W$ is placed at a point (at a distance $x$ from one of the end points) on the beam of length $a$. The load is placed perpendicular to the beam. Two reaction forces $\boldsymbol{R}_{\boldsymbol{A}}$ and $\boldsymbol{R}_{\boldsymbol{B}}$ at the end points of the beam occur in the opposite directions of the load $\boldsymbol{W}$. Figure 1 illustrates the beam and the forces.

Here we are interested in finding the magnitudes $R_{A}$ and $R_{B}$ of the reaction forces $\boldsymbol{R}_{A}$ and $\boldsymbol{R}_{B}$ on each leg of the beam. We assume that the beam and the legs are rigid. This problem can be easily solved using torques. The equilibrium equation system

$$
\begin{aligned}
& R_{A}+R_{B}-W=0 \\
& a \cdot R_{B}-x \cdot W=0
\end{aligned}
$$

has the following solution

$$
\begin{gathered}
R_{A}=W \cdot \frac{a-x}{a} \\
R_{B}=W \cdot \frac{x}{a}
\end{gathered}
$$

### 2.2. Geometric Interpretation

Given a leg of the beam, the point where the load is placed divides the beam into two segments, one adjacent to the given leg and one opposite to it, in each case, the solution takes the form

$$
R_{X}=W \cdot \frac{\text { Opposite Segment Length }}{\text { Total Beam Length }}
$$



Figure 1. The beam case consists of a two legs table and a load $W$ located at distance $x$ from one of the legs.

### 2.3. Triangular Table

Consider a triangular table of sides $a_{1}, a_{2}, a_{3}$ and legs at vertices $V_{1}, V_{2}, V_{3}$ of the table. A load $\boldsymbol{W}$ of magnitude $W$ is placed at a point on the table's surface and perpendicular to it. The point determines three heights $y_{1}, y_{2}$ and $y_{3}$ with respect the sides $a_{1}, a_{2}$ and $a_{3}$ respectively as in Figure 2 (we only draw $y_{1}$ and $h_{1}$ for simplicity).

Here we have a rigid surface and we are assuming that the legs do not bend. For $i=1,2,3$ denote by $R_{i}$ the magnitude of the force $\boldsymbol{R}_{\boldsymbol{i}}$. If we calculate the resultant torque on the body about an axis through $V_{2}$ and $V_{3}$ and we also calculate the resultant torque on the body about an axis through $V_{1}$ and $V_{3}$ we obtain the equilibrium equation system

$$
\begin{align*}
& R_{1}+R_{2}+R_{3}-W=0 \\
& h_{1} \cdot R_{1}-y_{1} \cdot W=0  \tag{2.1}\\
& h_{2} \cdot R_{2}-y_{2} \cdot W=0 .
\end{align*}
$$

When we solve the system we find that

$$
\begin{align*}
& R_{1}=W \cdot \frac{y_{1}}{h_{1}} \\
& R_{2}=W \cdot \frac{y_{2}}{h_{2}}  \tag{2.2}\\
& R_{3}=W-R_{1}-R_{2}
\end{align*}
$$

By adding the three small areas of the triangles $V_{1} V_{2} P, V_{1} V_{3} P$, and $V_{2} V_{3} P$ in Figure 3, we obtain that $\frac{a_{1} h_{1}}{2}=\frac{a_{1} y_{1}}{2}+\frac{a_{2} y_{2}}{2}+\frac{a_{3} y_{3}}{2}$. We note that the area of the triangle $V_{1} V_{2} V_{3}$ is $\frac{a_{1} h_{1}}{2}=\frac{a_{2} h_{2}}{2}=\frac{a_{3} h_{3}}{2}$. Using these identities and the previous equation, we obtain that

$$
\begin{aligned}
\frac{a_{1} h_{1}}{2} & =\frac{a_{1} y_{1}}{2}+\frac{a_{2} y_{2}}{2}+\frac{a_{3} y_{3}}{2}=\frac{a_{1} h_{1}}{2} \cdot \frac{y_{1}}{h_{1}}+\frac{a_{2} h_{2}}{2} \cdot \frac{y_{2}}{h_{2}}+\frac{a_{3} h_{3}}{2} \cdot \frac{y_{3}}{h_{3}} \\
& =\frac{a_{1} h_{1}}{2} \cdot \frac{y_{1}}{h_{1}}+\frac{a_{1} h_{1}}{2} \cdot \frac{y_{2}}{h_{2}}+\frac{a_{1} h_{1}}{2} \cdot \frac{y_{3}}{h_{3}} .
\end{aligned}
$$



Figure 2. In this triangular table, the point where the load $W$ is placed, determines height $y_{1}$ with respect to the vertex $V_{1}$.


Figure 3. The point $P$ divides the triangle in three triangles by drawing lines from $P$ to each vertex.

Hence

$$
\begin{equation*}
1=\frac{y_{1}}{h_{1}}+\frac{y_{2}}{h_{2}}+\frac{y_{3}}{h_{3}} . \tag{2.3}
\end{equation*}
$$

From Equation (2.3) and Equation (2.2) we can also conclude that $R_{3}=W \cdot \frac{y_{3}}{h_{3}}$.

### 2.4. Geometric Interpretation

On the table's surface, let's draw three lines from the point where the load is located to the three corners of the table. We obtain three smaller triangles (shaded differently on the figure) inside the big triangle as in Figure 3.

Among these three triangles, the triangle $P V_{2} V_{3}$ is not adjacent to the vertex $V_{1}$, we call this triangle the opposite triangle to $V_{1}$.

We observe that

$$
\frac{y_{1}}{h_{1}}=\frac{\left(y_{1} \cdot a_{1}\right) / 2}{\left(h_{1} \cdot a_{1}\right) / 2}=\frac{\text { Area of Triangle } P V_{2} V_{3}}{\text { Area of Triangle } V_{1} V_{2} V_{3}} .
$$

From the solution to the triangular case given by Equation (2.2) and the previous observation, we can establish a formula to obtain the magnitude of the reaction forces at any single leg in the triangular table according to

$$
R_{X}=W \cdot \frac{\text { Opposite Area }}{\text { Total Area }} .
$$

## 3. Rectangular Table, Indeterminate Case

Consider a rectangular table of sides $B$ and $H$ and legs at each of its vertices $V_{1}, V_{2}, V_{3}$ and $V_{4}$. A load $\boldsymbol{W}$ of magnitude $W$ is placed at a point on the table's surface and perpendicular to it. The point has distances to the sides $x$ and $y$ respectively as in Figure 4.

Here once more we have a rigid surface and we are assuming that the legs do not bend. For $i=1, \cdots, 4$ denote by $R_{i}$ the magnitude of the force $\boldsymbol{R}_{\boldsymbol{i}}$. If we calculate the resultant torque on the body about an axis through $V_{1}$ and $V_{2}$ and we also calculate the resultant torque on the body about an axis through $V_{1}$ and $V_{4}$ we obtain the equilibrium equation system

$$
\begin{align*}
& R_{1}+R_{2}+R_{3}+R_{4}-W=0 \\
& H \cdot R_{3}+H \cdot R_{4}-y \cdot W=0  \tag{3.4}\\
& B \cdot R_{2}+B \cdot R_{3}-x \cdot W=0 .
\end{align*}
$$



Figure 4. Rectangular table and load $W$ on the table placed at a point with distances $x$ and y to the sides of the table.

This is all the mechanics says about the system. Systems of three equations and four unknowns might have no solutions or have an infinite number of solutions according to the consistency of the system and the existence of non-trivial solutions to the homogeneous system. It is the case that this system is consistent and the homogeneous system has non-trivial solutions, hence it has an infinite amount of solutions. We call these type of systems indeterminate systems.

Patnaik and Hopkins [6] discussed the problem under different hypotheses. They assume that the legs deform and take in consideration the modulus of elasticity and the cross-sectional areal of the legs. By using the principle of least work, they arrive at a fourth equation which they refer to it as compatibility condition:

$$
\begin{equation*}
R_{1}+R_{3}=R_{2}+R_{4} \tag{3.5}
\end{equation*}
$$

Adding this equation to the system 3.4 we obtain solutions where sometimes the reactions magnitudes are negative according to where the load is placed (a negative magnitude of a reaction means that the reaction force is pointing in the same direction as $\boldsymbol{W}$ ).

We should note that the problem we are approaching here is different. We will define the problem in the next section.

## 4. Main Hypotheses

We are considering a rectangular table whose legs bent, shrink or deflect a sufficiently small amount that we are neglecting it, therefore in our case, the legs of this rectangular table deflect zero, i.e., they do not deflect at all. We are also assuming that the table's surface is rigid, uniform and homogeneous enough to disregard any conditions about the material, hence we are also assuming that our table is an ideal rigid table. Let's assume for the moment that the table weights 0 (at the end of the problem, to obtain the final reactions magnitudes at each leg, we could add $1 / 4$ of the table weight to each reaction by symmetry).

The fact that in our case the legs do not deflect is the main difference with the work of Patnaik and Hopkins [6]. Whereas in the results of Patnaik and Hopkins, the reactions on some legs might point downwards because their flexible-legs table tilts, in our case, that does not happen, all the reactions point upwards.

Main Hypothesis 4.1. The rectangular table consists of a rigid flat surface and rigid legs, hence neither the legs nor the surface deform or deflect. An ideal object of weight $W$ is placed at a point $P$ on the surface of the table causing a force to point downwards and perpendicular to the table's surface.

## 5. Solving the Rectangular Table Case

Consider a view from the top of the rectangular table described on Section 3. We are labeling the vertices with
$V_{1}, V_{2}, V_{3}, V_{4}$, starting with the bottom left corner and going around counter clockwise. These vertices are precisely where the reactions, perpendicular to the table's surface, occur. Assume then that the reaction $\boldsymbol{R}_{\boldsymbol{i}}$ of magnitude $R_{i}$ occur at the vertex $V_{i}$. For simplicity, we can assume that the vertices are in the cartesian plane and they have the following coordinates, $V_{1}=(0,0), \quad V_{2}=(B, 0), \quad V_{3}=(B, H)$, and $V_{4}=(0, H)$. Let $l_{i}$ be the segment from $V_{i}$ to $V_{i+1}$ for $i=1,2,3$ and let $l_{4}$ be the segment from $V_{4}$ to $V_{1}$. Let $P=(b, h)$ be the point where the load $\boldsymbol{W}$ is placed as in Figure 5.

Since the legs are rigid, the table does not tilt, hence each magnitude $R_{i}$ of the reaction $\boldsymbol{R}_{\boldsymbol{i}}$ is non-negative (a negative magnitude would mean that the reaction points downwards).

We can imagine for a second that legs 3 and 4 disappear and that we want to equilibrate the table by placing an extra leg, leg 34 at a point $P_{3}$ on the segment $l_{3}$ in such a way that our new triangular table supported on leg 1 , leg 2 and $\operatorname{leg} 34$ is still in equilibrium and leg 34 has a reaction of magnitude equal to $R_{3}+R_{4}$. Where do we place such new leg 34 ? There is certainly a geometrical interval on the segment $l_{3}$ determined by leg 1 , leg 2 and the point $P$ where this new leg can be placed with those conditions. This geometrical interval is in fact found by drawing a straight line passing through $V_{1}$ and $P$ and another one passing through $V_{2}$ and $P$. We extend these lines until they intersect also the prolonged segment $l_{3}$. Such interval must be within the segment $l_{3}$ and between these two intersection points. The reasoning behind this is that the point $P$ must be inside the triangle determined by $V_{1}, V_{2}$ and $P_{3}$. This interval can be appreciated in Figure 5.

Let's place such a leg at the point $P_{3}$ on the segment $l_{3}$ within the above interval, and let's call $t$ the number between 0 and 1 such that $P_{3}=V_{3}+t \cdot\left(V_{4}-V_{3}\right)=(B(1-t), H)$. In other words, $P_{3}$ is the vector $V_{3}$ plus $t$ times the vector that goes from $V_{3}$ to $V_{4}$. Let's call $\boldsymbol{R}_{34}(t)$ the reaction corresponding to the leg 34, and $R_{34}(t)$ its magnitude, of course these are functions of the parameter $t$. Figure 6 exemplifies this situation.

From Figure 6, we can assume that the table we are dealing with is a triangular table, therefore the solved case applies (Figure 7).

If we calculate the resultant torque on the body about the $x$-axis as well as calculating the torque about the $y$-axis, we can also obtain an equilibrium equation system (equivalent to System 2.1) for this triangular table


Figure 5. In this view from the top of the rectangular table described before, a load $W$ is placed on the table. On segment $l_{3}$ is drawn an interval where leg 34 will be placed.


Figure 6. In this view from the top of the rectangular table described before, reaction force $R_{34}(t)$ on segment $l_{3}$ replaces forces $R_{3}$ and $R_{4}$ for some $t$ in the interval $[0 ; 1]$.


Figure 7. Triangular table supported at $V 1 ; V 2$ and $P_{3}=P_{3}(t)$ has reaction forces $R_{34}(t) ; R_{1}(t)$ and $R_{3}(t)$.

$$
\begin{align*}
& R_{1}(t)+R_{2}(t)+R_{34}(t)-W=0 \\
& B \cdot R_{2}(t)+(B-B t) \cdot R_{34}(t)-b \cdot W=0  \tag{5.6}\\
& H \cdot R_{34}(t)-h \cdot W=0
\end{align*}
$$

We can easily solve for $R_{1}(t), R_{2}(t)$ and $R_{34}(t)$. Notice that all the magnitudes will be functions of the
coordinates of $P_{3}$, namely:

$$
\begin{aligned}
& R_{1}(t)=W \cdot \frac{(H-h t) B-H b}{H B} \\
& R_{2}(t)=W \cdot \frac{(h t-h) B+H b}{H B} \\
& R_{34}(t)=W \cdot \frac{h}{H}
\end{aligned}
$$

If we now remove leg 34 and place back the leg 3 and leg $4, R_{3}(t)$ and $R_{4}(t)$ can be computed using the case of the one dimensional beam as

$$
\begin{aligned}
& R_{3}(t)=W \cdot \frac{h-h t}{H} \\
& R_{4}(t)=W \cdot \frac{h t}{H}
\end{aligned}
$$

To summarize, the solution to the rectangular case that depends on the parameter $t$, is given by

$$
\begin{align*}
& R_{1}(t)=W \cdot \frac{(H-h t) B-H b}{H B} \\
& R_{2}(t)=W \cdot \frac{(h t-h) B+H b}{H B}  \tag{5.7}\\
& R_{3}(t)=W \cdot \frac{h-h t}{H} \\
& R_{4}(t)=W \cdot \frac{h t}{H}
\end{align*}
$$

The solution $\left\{R_{1}, R_{2}, R_{3}, R_{4}\right\}$ to the rectangular case that we are looking for (and that does not depend on any parameter) satisfies the following conditions:

## Condition 5.1.

$$
R_{1}, R_{2}, R_{3}, R_{4} \geq 0
$$

The reactions $\boldsymbol{R}_{3}$ and $\boldsymbol{R}_{4}$ can be replaced by a reaction $\boldsymbol{R}_{34}$ of magnitude $R_{3}+R_{4}$ at some point $P_{3}$ on the segment $l_{3}$.

The procedure that follows consists of finding such point $P_{3}$ by establishing some physical conditions on the net torque.

Similar to the description above, the reactions $\boldsymbol{R}_{1}$ and $\boldsymbol{R}_{2}$ can be replaced by a reaction $\boldsymbol{R}_{12}$ of magnitude $R_{1}+R_{2}$ at some point $P_{1}$ on the segment $l_{1}$ within a geometrical interval. We might also say that the point $P_{1}$ depends on a parameter $r$, where $P_{1}=V_{1}+r \cdot\left(V_{2}-V_{1}\right)$; however, if we think that the reaction $\boldsymbol{R}_{12}$ substitutes the reactions $\boldsymbol{R}_{1}, \boldsymbol{R}_{2}$ and the reaction $\boldsymbol{R}_{34}$ substitutes the reactions $\boldsymbol{R}_{3}$ and $\boldsymbol{R}_{4}$, then the reactions $\boldsymbol{R}_{12}, \boldsymbol{R}_{34}$ are substituting $\boldsymbol{R}_{1}, \boldsymbol{R}_{2}, \boldsymbol{R}_{3}$ and $\boldsymbol{R}_{4}$, and moreover, the table is still in equilibrium (we are assuming that the table weights zero), hence we do have a new condition.

Condition 5.2. The points $P_{3}, P_{1}$ and $P$ are collinear! (as if the whole table were supported on a bar only).
Under this condition, the point $P_{1}$ does not depend on $r$ anymore, the point $P_{3}$ determines the point $P_{1}$, and hence $P_{1}$ depends on $t$, in fact, we have

$$
P_{1}=\left(\frac{H b+B(h t-h)}{H-h}, 0\right) .
$$

When we do the same analysis as before, but this time using leg 3 , leg 4 and $\operatorname{leg} 12$ (at the point $P_{1}$ ), we obtain the same values $R_{1}(t), R_{2}(t), R_{3}(t)$ and $R_{4}(t)$, as in the system 5.7.

Similarly, we can now think that the reactions $\boldsymbol{R}_{4}$ and $\boldsymbol{R}_{1}$ can be replaced by a reaction $\boldsymbol{R}_{41}$ of magnitude $R_{4}+R_{1}$ at some point $P_{4}(u)$ on the segment $l_{4}$, (within some interval) where $P_{4}=V_{4}+u \cdot\left(V_{1}-V_{4}\right)=(0, H(1-u))$ as in Figure 8.
Performing the same analysis as before, we calculate the resultant torque on the body about the axis determined by $V_{1}$ and $V_{2}$ as well as calculating the torque about the axis determined by $V_{2}$ and $V_{3}$. We obtain the equilibrium system

$$
\begin{align*}
& R_{2}(u)+R_{3}(u)+R_{41}(u)-W=0 \\
& H \cdot R_{3}(u)+(H-H u) \cdot R_{41}(u)-h \cdot W=0  \tag{5.8}\\
& B \cdot R_{41}(u)-(B-b) \cdot W=0,
\end{align*}
$$

Proceeding as before, we arrive at the solution to the system 5.8

$$
\begin{align*}
& R_{1}(u)=W \cdot \frac{u(B-b)}{B} \\
& R_{2}(u)=W \cdot \frac{B(H-h)-H u(B-b)}{H B}  \tag{5.9}\\
& R_{3}(u)=W \cdot \frac{h B-H(1-u)(B-b)}{H B} \\
& R_{4}(u)=W \cdot \frac{(1-u)(B-b)}{B} .
\end{align*}
$$

Analogously to Condition 5.2, the point $P_{4}$ determines a point $P_{2}$ on the segment $l_{2}$ collinear with $P$ and $P_{4}$. In fact (Figure 9)

$$
P_{2}=\left(B, \frac{B h-H(B-b)(1-u)}{b}\right) .
$$



Figure 8. In this view from the top of the rectangular table described before, reaction force $R_{41}(u)$ on segment $l_{4}$ replaces forces $R_{4}$ and $R_{1}$ for some $u$ in the interval [0;1].


Figure 9. Points $P_{3} ; P$ and $P_{1}$ are collinear as well as points $P_{4} ; P$ and $P_{2}$.
When doing the same analysis, such point also determines the same values $R_{1}(u), R_{2}(u), R_{3}(u)$ and $R_{4}(u)$.
These conditions of collinearity, do not provide any extra information, it only allows us reduce to two parameters $(t, u)$ instead of four.

A natural assumption is that the solution given by considering the point $P_{3}(t)$ will also match the solutions given by considering the point $P_{4}(u)$, in other words.

Condition 5.3. The solutions to the system 5.9 agree with the solutions to system 5.7, in other words, the system is self-consistent.

This brings us to the equation

$$
\begin{equation*}
t=\frac{H}{h B} \cdot(B-b) \cdot(1-u) . \tag{5.10}
\end{equation*}
$$

This does not solve the problem; however, whichever the solution is, there are parameters $t$ and $u$ such that the solution can be expressed as in the system 5.7 or 5.9, and those parameters should be related by Equation (5.10).

Since we are considering a rigid surface, we are expecting that the surface strains only a little bit, we really want the surface to strain the least to be in equilibrium, it is then a natural requirement to ask that the table strains the least in two directions. Let's consider only the line that goes from $P_{3}$ to $P_{1}$. Recall that

$$
\begin{aligned}
R_{34}(t)=W \cdot \frac{h}{H}, \quad R_{12}(t) & =W \cdot \frac{H-h}{H} \text { and that } \\
P & =(b, h), \quad P_{3}=(B(1-t), H), \quad P_{1}=\left(\frac{H b+B(h t-h)}{H-h}, 0\right) .
\end{aligned}
$$

A way of measuring the deflection of the surface along this line is by considering the sum of the magnitudes of the torques at the points $P_{3}$ and $P_{1}$, with respect to the point $P$, this is

$$
\text { distance }\left(P, P_{3}\right) \cdot R_{34}(t)+\operatorname{distance}\left(P, P_{1}\right) \cdot R_{12}(t)
$$

In our case, since $P_{3}, P$ and $P_{1}$ are collinear, we can use similar triangles to obtain that this previous quantity is

$$
\frac{2 \sqrt{\left(t^{2}-2 t+1\right) B^{2}+2(b t-b) B+H^{2}-2 H h+h^{2}+b^{2}} \cdot h}{H}
$$

Since we are asking this quantity to be minimum, we need to minimize this as a function of $t$. The previous expression is minimum when the inside of the square root is minimum, hence to minimize the previous expression, we set to zero the derivative of

$$
\left(t^{2}-2 t+1\right) B^{2}+2(b t-b) B+H^{2}-2 H h+h^{2}+b^{2}
$$

with respect to $t$, which gives $2(t-1) B^{2}+2 B b=0$, hence

$$
t=\frac{B-b}{B}
$$

therefore the point $P_{3}$ is located at $P_{3}=(b, H)$. Plugging these values in Solution 5.7, we obtain

$$
\begin{align*}
& R_{1}=W \cdot \frac{(B-b)(H-h)}{H B} \\
& R_{2}=W \cdot \frac{b(H-h)}{H B}  \tag{5.11}\\
& R_{3}=W \cdot \frac{h b}{H B} \\
& R_{4}=W \cdot \frac{(B-b) h}{H B} .
\end{align*}
$$

If we now consider the other line that goes from $P_{4}$ to $P_{2}$, we can also measure the deflection of the surface along this line by considering the sum of the magnitudes of the torques at the points $P_{4}$ and $P_{2}$, with respect to the point $P$, this is

$$
\operatorname{distance}\left(P, P_{4}\right) \cdot R_{41}(u)+\operatorname{distance}\left(P, P_{2}\right) \cdot R_{23}(u)
$$

Once more, using collinearity and similar triangles, this quantity becomes

$$
\frac{2(B-b) \sqrt{\left(u^{2}-2 u+1\right) H^{2}+2(h u-h) H+h^{2}+b^{2}}}{B}
$$

By a similar argument, minimizing the previous expression is equivalent to minimize the expression inside the square root as a function of $u$. Therefore, we need to set to zero the derivative of

$$
\left(u^{2}-2 u+1\right) H^{2}+2(h u-h) H+h^{2}+b^{2}
$$

which gives $2(u-1) H^{2}+2 H h=0$, hence

$$
u=\frac{H-h}{H}
$$

Notice that this value of $u$ agrees with the one obtained by substituting $t=\frac{B-b}{B}$ in Equation (5.10).
Therefore the point $P_{4}$ is located at

$$
P_{4}=(0, H-h)
$$

which also provides, according to System 5.7, the solutions

$$
\begin{align*}
& R_{1}=W \cdot \frac{(B-b)(H-h)}{H B} \\
& R_{2}=W \cdot \frac{b(H-h)}{H B}  \tag{5.12}\\
& R_{3}=W \cdot \frac{h b}{H B} \\
& R_{4}=W \cdot \frac{(B-b) h}{H B} .
\end{align*}
$$

In a forthcoming paper, we will generalize these results.

## Geometric Interpretation

We notice that the solutions achieved here are self consistent, the points $P_{1}, P, P_{3}$ are collinear as well as the points $P_{2}, P, P_{4}$, the net-torque quantities are minimum. We also have that the lines joining the collinear points determined four rectangles.

From Figure 10, it can be easily seen, by drawing the corresponding lines, that the points $P_{3}$ and $P_{4}$ lie within the geometrical interval discussed at the beginning of the previous section.

To each vertex of the table, among the four areas of those rectangles, one area is where the vertex is located, two areas are adjacent to the vertex and one area is opposite to the vertex. The formula obtained in Solution 5.11 or in 5.12 to compute the magnitudes of the reactions at any single leg in the rectangular table is given by

$$
R_{X}=W \cdot \frac{\text { Opposite Area }}{\text { Total Area }}
$$

as proposed by Fontana [3] (when the parallelogram becomes a rectangle); however, Fontana imagined that this formula had to be true without any hypothesis and never justified such formula. We, in this paper, stated the conditions under which this formula is true.

## 6. Experimental Results

### 6.1. Summary of Experimental Results

To test the predictions of the Equation (5.11) above, load weights of mass 0.200 kg up to 3.50 kg were suspended from a rectangular aluminum plate. Reaction forces were measured by digital scales affixed to the corners of the plate. A detailed procedure is described below. The resulting experimental reaction forces agree well with the theory, as seen in Figure 11. The experimental values versus the theoretical ones are fit by a straight line and the resulting fitting parameters are consistent with agreement. The slope is consistent with one (best fit slope of $1.02 \pm 0.03$ ) and $y$-intercept is consistent zero (best fit $-0.62 \pm 1.0$ ) to within the one sigma uncertainties of the slope and intercept parameters.

In Figure 11, the solid line is a best fit line to the data and is consistent with agreement between theory and experiment: The slope is consistent with one (best fit slope of $1.02 \pm 0.03$ ) and $y$-intercept is consistent zero (best fit $-0.62 \pm 1.0$ ) to within the one sigma uncertainties of the slope and intercept parameters. Each experi-


Figure 10. The point $P$ where the load $W$ is placed divides the rectangle into four rectangles. In this gure, rectangle $P_{1} P P_{2} V_{2}$ is opposite to vertex $V_{4}$.


Figure 11. Experimental versus theoretical fractions of the load weight found at each reaction force.
mental fraction reaction force is the slope of a linear fit such as those in Figure 12 (reaction forces versus load weights). The $y$-error bars on the data are uncertainties from the resulting linear fits, such as in Figure 12. The $x$-error bars are from calculations of the uncertainties of the positions of the load weights, limited by the apparatus.

Each reaction force in Figure 11 is derived from a series of data of the measurement of reaction force versus load weight. In Figure 12, the four reaction forces are plotted versus load weight for one of the three locations (called Location 2, details below). The slope of each linear fit is included in Table 1. To form the fraction supported by one reaction force, the slope corresponding to that force is divided by the sum of the four slopes.

In Figure 12, total plate size is $(34.35,22.9) \mathrm{cm}$. Data are summarized in Table 1.
On Table 1, the slopes and $y$-intercepts are from the linear regression fits of the experimental reaction forces versus load weight, plotted in Figure 12. The fractions of the load weight supported by each leg (the $R_{n}$ 's) are listed as $f(\exp )$ and $f$ (theory). The experimental values are the slopes divided by the sum of slopes. Sum of slopes is $34.96 \mathrm{oz} / \mathrm{kg}$ which is close to the true value of $35.274 \mathrm{oz} / \mathrm{kg}$, and the difference is just outside the 0.25 $\mathrm{oz} / \mathrm{kg}$ one sigma uncertainty of the sum of the uncertainties of the slopes. The sum of the $y$-intercepts is zero to two significant figures.

On Table 2, the three locations load weights were placed are labeled 2, 3, and 6. Experimental fractions of support of the load weight (labeled $f(\exp )$ in this table) are computed as in Table 1: by dividing the slope of the individual reaction force versus load weight graph by the sum of the slopes. Theory is computed by Equation (5.12). The uncertainty is solely from the uncertainty in the experimental slope. Although the agreement is quite good between experiment and theory, there is additional uncertainty beyond only that of the linear regression fits. Possible sources are discussed below.

On Table 3, all values have a measurement uncertainty of 0.1 cm , but the nuts affixing the machine screws to the plate have diameters of 0.8 cm , thus this latter uncertainty is employed to produce the $x$-error bars on Figure 12.

### 6.2. Estimation of Uncertainties

The slopes and $y$-intercept values in the data Table 1 and Table 2 have uncertainties which arise from the linear regression fitting (performed with Originlabs Origin software). Individual reaction force measurements have


Figure 12. Reaction forces for the load weight placed at $(20.7,12.6) \mathrm{cm}$.
Table 1. The table of reaction forces for weight at location two, given by $(20.7,12.6) \mathrm{cm}$ in the $(34.35,22.9) \mathrm{cm}$ plate.

| Force | $y$-int (oz) | slope (oz) | $f(\exp )$ | $f($ theory $)$ | unc |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Units: | oz | kg/oz | - | - | - |
| $R_{1}$ | 0.30 | 6.62 | 18.9 | 17.9 | 0.51 |
| $R_{3}$ | -0.29 | 9.15 | 26.2 | 27.1 | 0.29 |
| $R_{2}$ | 0.28 | 11.98 | 34.3 | 33.2 | 0.20 |
| $R_{4}$ | -0.29 | 7.213 | 20.6 | 21.9 | 0.23 |

Table 2. Summary of all three experiments.

| Loc | Support | $f(\exp )$ | $f$ (theory) | unc(exp) | unc(theory) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | $R_{1}$ | 18.9 | 17.9 | 0.51 | 0.87 |
| 2 | $R_{3}$ | 26.2 | 27.1 | 0.29 | 1.2 |
| 2 | $R_{2}$ | 34.3 | 33.2 | 0.20 | 1.2 |
| 2 | $R_{4}$ | 20.6 | 21.9 | 0.23 | 0.94 |
| 6 | $R_{2}$ | 6.8 | 9.1 | 0.69 | 0.60 |
| 6 | $R_{3}$ | 23.2 | 21.1 | 0.82 | 0.97 |
| 6 | $R_{1}$ | 46.8 | 48.7 | 0.80 | 1.5 |
| 6 | $R_{4}$ | 23.3 | 21.0 | 0.68 | 1.3 |
| 3 | $R_{2}$ | 10.9 | 10.2 | 0.48 | 1.2 |
| 3 | $R_{3}$ | 56.3 | 56.7 | 0.53 | 1.5 |
| 3 | $R_{1}$ | 28.5 | 28.0 | 0.62 | 1.1 |
| 3 | $R_{4}$ | 4.2 | 5.1 | 0.53 | 0.60 |

Table 3. Coordinates of the three load weights (Loc 2, 6 and 3) and the four corner supports (Labeled AD) in rectangular form.

| Label | $x(\mathrm{~cm})$ | $y(\mathrm{~cm})$ |
| :---: | :---: | :---: |
| Loc 2 | 20.7 | 12.6 |
| Loc 6 | 10.4 | 6.9 |
| Loc 3 | 23.0 | 3.5 |
| $R_{1}$ | 0 | 0 |
| $R_{3}$ | 34.35 | 0 |
| $R_{2}$ | 34.3 | 22.9 |
| $R_{4}$ | 0 | 22.9 |

uncertainties of 10 grams due to the limitations of the accuracy of the scales (this corresponds, roughly, to the size of the symbols in Figure 12). The masses used as load weights are accurate to under 1\% which is negligible compared to other sources.

In calculating the theoretical rigid plate reaction forces, from Equation (5.11), above, the dominant source of uncertainty is the location of the load weights. These are constrained by the size of the nuts, which are $8.0 \pm 0.2$ mm across (from corner to corner, the longest distance across the hexagonal nut) thus the uncertainty of the load weight position is taken to be $\pm 4 \mathrm{~mm}$, which produces the horizontal error bars in Figure 12. The other measurements dealing with the dimensions of the plate are, somewhat arbitrarily, assigned zero uncertainty. If the uncertainty due to the nuts (same size as above) at the corners of the plate are included, significantly larger error bars are produced.

### 6.3. Procedure

The goal of the procedure is to measure the distribution of weight from a single load weight onto supports places at the four corners of the table. To do so, it is important that the table remain level during measurements. Several methods were investigated to reach optimal conditions the most successful of which are detailed below (Figure 13).

Four PASCO SE-9372 laboratory jacks are places close to, but outside, where the four corners of the rectangular plate will lie. Height of each jack is set to near the middle of the range of travel, allowing future up or down adjustment. At approximately equal height, a right angle clamp is affixed to each ring stand. Steel crossbars are connected to pairs of ring stands. A carpenter's level is used to check for levelness. With meter stick, ensure the crossbars are roughly parallel and separated by the width of the plate.

The longer side of the rectangle is here termed the length; the shorter side the width. The longer sides should be directly under the crossbars. PE scales are clamped to crossbars. Scales are connected to the plate via steel machine hooks with threads which screw into the plate. To avoid stripping the threads, washers are placed below the plate and tightened. The threaded screws give a second method (after the lab jacks) to level the plate. At this point, level the plate adjusting only the screws. Ensure the hooks below the PE scales are vertical. To detect any deviation from levelness, one or more carpenter's levels can be left on the plate. The four PE scales are tarred at this point.

Now the load weight can be suspended from the plate and the reaction forces can be measured by the four scales. However, if the load weight is not placed near the center of the plate, large load weights cause the plate to rotate out of level. For these larger load weights, re-level the plate using the lab jacks, then remove load weight. Plate will rotate somewhat out of level. Tare scales. Add load weight (which will rotate the plate back very close to level) and record reaction force readings.

### 6.4. Materials

Plates are fabricated from 6061 aluminum with a thickness of 3.18 mm with threaded holes for the load weights and measuring reaction forces of inner diameter 2.69 mm , outer diameter of 3.51 mm . Steel washers connecting


Figure 13. One of the PE scales supporting the rectangular plate and a dark hanging weight below. The black threaded eye bolt has black washers clamping it to the plate both above and below the plate. However, these can be loosened and the threaded bolt can be adjusted slightly to help level the plate, initially. Later re-leveling of the plate for large masses is done via lab jacks (not pictured) which support the ring stands which in turn support the horizontal crossbars (seen near the top of the gure).
screws to plate have a width from corner to corner of 8.0 mm and a width from flat side to parallel flat side of 7.0 mm and a height of about 2.5 mm .

Weights and scales are affixed to the plate via steel machine screws, clamped to plate via washers. Scales are West-Boao Science \& Technology Co., Ltd. brand ETWT001 portable electronic scales, with maximum load of 50 kg and accuracy to 10 grams.

## 7. Conclusion

Through this paper the well known beam case was provided, the triangular table case was also provided and the geometrical interpretations of the reaction forces at any given leg were given. The indeterminate case of the rectangular table, also well known, was discussed. The hypotheses to our problem in Section 4 clarified the situation of the problem we wanted to solve. By employing the method of introducing two parameters and supporting the table on ideal beams and minimizing a quantity that somehow measures the deflection of the table surface, we were able to solve the problem in Section 5. The main difference of our method and the one provided by Patnaik, Hopkins and Halford [6] is that we have different hypotheses. The results from Patnaik, Hopkins and Halford [6] and our result do not contradict each other, we just approach the same problem but with different conditions. The method that we provide here will be generalized in a forthcoming paper by the same authors. A physical experiment was performed and described in Section 6. In the same section, laboratory results were obtained and these results are consistent with the theoretical values obtained by our proposed theoretical model. Although the solution provided here has not been analyzed before, its novelty might contribute to some applications on the mechanics and engineering of solid structures.

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