

Weierstrass' Elliptic Function Solution to the Autonomous Limit of the String Equation of Type (2,5)*

Yoshikatsu Sasaki

Department of Mathematics, Hiroshima University, Higashi-Hiroshima, Japan
 Email: sasakiyo@hiroshima-u.ac.jp

Received 1 June 2014; revised 3 July 2014; accepted 15 July 2014

Copyright © 2014 by author and Scientific Research Publishing Inc.
 This work is licensed under the Creative Commons Attribution International License (CC BY).
<http://creativecommons.org/licenses/by/4.0/>



Open Access

Abstract

In this article, we study the string equation of type (2,5), which is derived from 2D gravity theory or the string theory. We consider the equation as a 4th order analogue of the first Painlevé equation, take the autonomous limit, and solve it concretely by use of the Weierstrass' elliptic function.

Keywords

Painlevé Hierarchy, String Equation, Elliptic Function

1. Introduction

1.1. The String Equation of Type (2,5)

Put $D = d/dz$. Consider the commutator equation of ordinary differential operators

$$[Q, P] = 1, \quad Q := \sum_{k=2}^q w_k D^{q-k}, \quad P := \sum_{k=2}^p v_k D^{p-k}.$$

We call it the string equation (or Douglas equation) of type (q, p) , which appears in the string theory or the theory of quantum gravity in 2D [1]-[9]. In the followings, we set $q = 2$, $p = 2g + 1$.

In the case where $q = 2$, $p = 3$, the string equation is written as an ODE satisfied by the potential w of Sturm-Liouville operator $Q = D^2 + w$, and then, by a fractional linear transformation, it is reduced to the first Painlevé equation [10] [11]

$$w'' = 6w^2 + z, \quad (\text{PI})$$

*Dedicated to Professor Masafumi Yoshino on the occasion of his 60th birthday.

which is equivalent to the Hamiltonian system:

$$dw/dz = \partial H / \partial v, \quad dv/dz = -\partial H / \partial w, \quad H = \frac{1}{2}v^2 - 2w^3 - zw.$$

In the case where $q = 2$, $p = 5$, $[Q, P] = 1$ yields

$$C_0 = w^{(4)} + 5w'^2 + 10(w' + C_1)(w'' + 3w^2) - 20w^3 + 16C_2w + 16z,$$

where C_0, C_1, C_2 are integral constants. By the fractional linear transformation $z \mapsto \alpha z + \beta$, $w \mapsto \gamma w + \delta$,

$$\alpha^7 = -\frac{1}{3}, \quad \gamma = 6\alpha^5, \quad \delta = -C_1, \quad 16\beta = C_0 - 20C_1^3 + 16C_1C_2$$

and putting $a = \alpha^4(8C_2 - 15C_1^2)/4$, the string equation is reduced to

$$w^{(4)} = 20w''w + 10w'^2 - 40w^3 - 8aw - \frac{8}{3}z, \quad (\text{S})$$

We also call it the string equation of type (2,5). Note that (S) coincides the 4th order equation of *the first Painlevé hierarchy* [12]-[15]

$$d_n[w] + 4z = 0, \quad (2_n\text{PI})$$

for $n \in \mathbb{N}$, where $d_n[w]$ is an expression of a given meromorphic function w defined by $d_0[w] = -4w$ and $Dd_{n+1}[w] = (D^3 - 8wD - 4w')d_n[w]$.

1.2. Degenerated Garnier System

Equation (S) is also obtained as follows. Consider a 2D degenerated Garnier system [16] [17]:

$$\partial q_i / \partial t_j = \partial H_j / \partial p_i, \quad \partial p_i / \partial t_j = -\partial H_j / \partial q_i, \quad (i, j \in \{1, 2\}), \quad (\text{dG}_{9/2})$$

$$H_1 = \frac{1}{3} \left(q_2^2 - q_1 - \frac{1}{3}t_1 \right) p_1^2 + \frac{2}{3} q_2 p_1 p_2 + \frac{1}{3} p_2^2 + 3 \left(q_1 + \frac{1}{3}t_1 \right) q_2 \left(q_2^2 - 2q_1 + \frac{1}{3}t_1 \right) - t_2 q_1,$$

$$H_2 = \frac{1}{3} q_2 p_1^2 + \frac{2}{3} p_1 p_2 + 3q_2^4 - 9q_1 q_2^2 + 3q_1^2 - t_1 q_1 - t_2 q_2.$$

which is a 2D analogue of (PI) in the theory of isomonodromic deformations. If we fix one of the independent variables $t_1 \equiv a (= \text{const.})$, we get a Hamiltonian system with only one independent variable $t_2 \equiv z$ as follows:

$$\partial q_i / \partial z = \partial H / \partial p_i, \quad \partial p_i / \partial z = -\partial H / \partial q_i, \quad (i \in \{1, 2\}),$$

$$H (= H_2) = \frac{1}{3} q_2 p_1^2 + \frac{2}{3} p_1 p_2 + 3q_2^4 - 9q_1 q_2^2 + 3q_1^2 - a q_1 - z q_2.$$

From the above system, eliminating q_1, p_1, p_2 and putting $w = q_2$, we obtain (S). So, Equation (S) is 4th order analogue of (PI) in the double sences.

It is already known by Shimomura [18] that every solution to (S) is meromorphic on \mathbb{C} , and that every pole of every solution is double one with its residue 0.

1.3. Autonomous Limit of the First Painlevé Equation

The first Painlevé equation (PI) has the autonomous limit [11]. Replacing (w, v, z, H) by $(\varepsilon^{-2}w, \varepsilon^{-3}v, \varepsilon z + \varepsilon^{-4}b, \varepsilon^{-6}H)$ with a constant $b \in \mathbb{C}$, and taking limit $\varepsilon \rightarrow 0$, we obtain $w'' = 6w^2 + b$ which is solved by the Weierstrass' elliptic function [10] [11]. The relation between the fundamental 2-form before and after the replacement is

$$dw \wedge dv - dH \wedge dz \mapsto \varepsilon^{-5} (dw \wedge dv - dH \wedge dz).$$

1.4. Results

It is quite natural to think that:

Conjecture. Each equation of the first Painlevé hierarchy has the autonomous limit, and which is satisfied by the Weierstrass' elliptic function.

For $n = 2$, the statement is valid, *i.e.*

Theorem A. Replacing (w, z, a) by $(\varepsilon^{-2}w, \varepsilon z + \varepsilon^{-6}b, \varepsilon^{-4}a)$, or replacing $(q_1, q_2, p_1, p_2, z, H)$ by $(\varepsilon^{-4}q_1, \varepsilon^{-2}q_2, \varepsilon^{-3}p_1, \varepsilon^{-5}p_2, \varepsilon z + \varepsilon^{-6}b, \varepsilon^{-8}H)$ with a constant $b \in \mathbb{C}$, and taking limit $\varepsilon \rightarrow 0$, we obtain the autonomous limit of the 4th order equation of the first Painlevé hierarchy (S). Moreover, the relation between the fundamental 2-form before and after the replacement is

$$dp_1 \wedge dq_1 + dp_2 \wedge dq_2 - dH \wedge dz \mapsto \varepsilon^{-7} (dp_1 \wedge dq_1 + dp_2 \wedge dq_2 - dH \wedge dz).$$

It is easy to show the above. The autonomous limit is given by

$$w^{(4)} = 20w''w + 10w'^2 - 40w^3 - 8aw - \frac{8}{3}b, \quad (\text{A})$$

Theorem B. The autonomous limit Equation (A) has a solution concretely described by the Weierstrass' elliptic function as

$$w(z) = 4\wp(z)/a_1,$$

where $a_1 = (8 \pm 4\sqrt{-2})/3$.

Remark. Modulus of the elliptic function is determined by the constants a and b . g_2 and g_3 in the elliptic function theory are as follows:

$$g_2 = -4a_1a/(40 - 9a_1), \quad g_3 = -a_1^2b/6(10 - 3a_1).$$

The next section is devoted to give the proof of Theorem B.

2. Proof of Theorem B

Put $\varphi = -[\text{l.h.s. of (A)}] + [\text{r.h.s. of (A)}]$, *i.e.*

$$\varphi := w^{(4)} + 20w''w + 10w'^2 - 40w^3 - 8aw - \frac{8}{3}b = 0. \quad (1)$$

Multiplying both sides of $\varphi = 0$ by w' , and integrating it, we obtain a first integral of (A)

$$\int \varphi w' dz := -w'w''' + \frac{1}{2}w''^2 + 10w'^2w - 10w^4 - 4aw^2 - \frac{8}{3}bw = c : \text{const.} \quad (2)$$

In order to find the elliptic function solution, let w satisfy the relation:

$$w'^2 = a_0w^4 + a_1w^3 + a_2w^2 + a_3w + a_4 =: A(w). \quad (3)$$

Substituting (3), $w'' = \frac{1}{2}A_w(w)$ and $w''' = \frac{1}{2}A_{ww}(w)w'$ into (2), we have

$$\begin{aligned} \int \varphi w' dz &= -\frac{1}{2}A_{ww}(w)A(w) + \frac{1}{8}A_w(w)^2 + 10wA(w) - 10w^4 - 4aw^2 - \frac{8}{3}bw = c \\ &= [-4a_0^2]w^6 + [-6a_0a_1 + 10a_0]w^5 + \left[-5a_0a_2 - \frac{15}{8}a_1^2 + 10a_1 - 10\right]w^4 \\ &\quad + \left[-5a_0a_3 - \frac{5}{2}a_1a_2 + 10a_2\right]w^3 + \left[-6a_0a_4 - \frac{9}{4}a_1a_3 - \frac{1}{2}a_2^2 + 10a_3 - 4a\right]w^2 \\ &\quad + \left[-3a_1a_4 - \frac{1}{2}a_2a_3 + 10a_4 - \frac{8}{3}b\right]w + \left[-a_2a_4 + \frac{1}{8}a_3^2\right] \cdot 1. \end{aligned}$$

So, if we take

$$a_0 = a_2 = 0, \quad a_1 = \frac{1}{3}(8 \pm 4\sqrt{-2}), \quad a_3 = 16a/(40 - 9a_1), \quad a_4 = 8b/3(10 - 3a_1), \quad \text{and } c = \frac{1}{8}a_3^2,$$

then solutions of (3) satisfy (2). Now, in order to reduce $w'^2 = a_1w^3 + a_3w + a_4$ to $\wp'^2 = 4\wp^3 - g_2\wp - g_3$, we

use the scale transformation $w = \chi \wp$, $\chi \in \mathbb{C} \setminus \{0\}$. Immediately we obtain $\chi = 4/a_1$, and also $g_2 = -a_3/\chi$, $g_3 = -a_4/\chi$. \square

References

- [1] Douglas, M.R. (1990) String in Less than One-Dimensions and K-dV Hierarchies. *Physics Letters B*, **238**, 176-180. [http://dx.doi.org/10.1016/0370-2693\(90\)91716-O](http://dx.doi.org/10.1016/0370-2693(90)91716-O)
- [2] Moore, G. (1990) Geometry of the String Equations. *Communications in Mathematical Physics*, **133**, 261-304. <http://dx.doi.org/10.1007/BF02097368>
- [3] Moore, G. (1991) Matrix Models of 2D Gravity and Isomonodromic Deformations. *Progress of Theoretical Physics Supplement*, **102**, 255-285. <http://dx.doi.org/10.1143/PTPS.102.255>
- [4] Fukuma, M., Kawai, H. and Nakayama, R. (1991) Infinite Dimensional Grassmannian Structure of Two Dimensional String Theory. *Communications in Mathematical Physics*, **143**, 371-403. <http://dx.doi.org/10.1007/BF02099014>
- [5] Kac, V. and Schwarz, A. (1991) Geometric Interpretation of Partition Functions of 2D Gravity. *Physics Letters B*, **257**, 329-334. [http://dx.doi.org/10.1016/0370-2693\(91\)91901-7](http://dx.doi.org/10.1016/0370-2693(91)91901-7)
- [6] Schwarz, A. (1991) On Solutions to the String Equations. *Modern Physics Letters A*, **29**, 2713-2725. <http://dx.doi.org/10.1142/S0217732391003171>
- [7] Adler, M. and van Moerbeke, P. (1992) A Matrix Integral Solution to Two-Dimensional W_p -Gravity. *Communications in Mathematical Physics*, **147**, 25-26. <http://dx.doi.org/10.1007/BF02099527>
- [8] van Moerbeke, P. (1994) Integrable Foudations of String Theory. In: Babelon, O., et al., Ed., *Lectures on Integrable Systems*, World Science Publisher, Singapore, 163-267.
- [9] Takasaki, K. (2007) Hamiltonian Structure of PI Hierarchy. *SIGMA*, **3**, 42-116.
- [10] Ince, E.L. (1956) Ordinary Differential Equations. Dover Publications, New York.
- [11] Conte, R. and Mussette, M. (2008) The Painlevé Handbook. Springer Science + Business Media B.V., Dordrecht.
- [12] Weiss, J. (1984) On Classes of Integrable Systems and the Painlevé Property. *Journal of Mathematical Physics*, **25**, 13-24. <http://dx.doi.org/10.1063/1.526009>
- [13] Kudryashov, N.A. (1997) The First and Second Painlevé Equations of Higher Order and Some Relations between Them. *Physics Letters A*, **224**, 353-360. [http://dx.doi.org/10.1016/S0375-9601\(96\)00795-5](http://dx.doi.org/10.1016/S0375-9601(96)00795-5)
- [14] Gromak, V.I., Laine, I. and Shimomura, S. (2002) Painlevé Differential Equations in the Complex Plane. Walter de Gruyter, Berlin. <http://dx.doi.org/10.1515/9783110198096>
- [15] Shimomura, S. (2004) Poles and α -Points of Meromorphic Solutions of the First Painlevé Hierarchy. *Publications of the Research Institute for Mathematical Sciences, Kyoto University*, **40**, 471-485. <http://dx.doi.org/10.2977/prims/1145475811>
- [16] Kimura, H. (1989) The Degeneration of the Two Dimensional Garnier System and the Polynomial Hamiltonian Structure. *Annali di Matematica Pura ed Applicata*, **155**, 25-74. <http://dx.doi.org/10.1007/BF01765933>
- [17] Suzuki, M. (2006) Spaces of Initial Conditions of Garnier System and Its Degenerate Systems in Two Variables. *Journal of the Mathematical Society of Japan*, **58**, 1079-1117. <http://dx.doi.org/10.2969/jmsj/1179759538>
- [18] Shimomura, S. (2000) Painlevé Property of a Degenerate Garnier System of (9/2)-Type and a Certain Fourth Order Non-Linear Ordinary Differential Equation. *Annali della Scuola Normale Superiore di Pisa, Classe di Scienze*, **29**, 1-17.

Scientific Research Publishing (SCIRP) is one of the largest Open Access journal publishers. It is currently publishing more than 200 open access, online, peer-reviewed journals covering a wide range of academic disciplines. SCIRP serves the worldwide academic communities and contributes to the progress and application of science with its publication.

Other selected journals from SCIRP are listed as below. Submit your manuscript to us via either submit@scirp.org or [Online Submission Portal](#).

