# Higher-Order Numeric Solutions for Nonlinear Systems Based on the Modified Decomposition Method 

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#### Abstract

Higher-order numeric solutions for nonlinear differential equations based on the Rach-Adomian-Meyers modified decomposition method are designed in this work. The presented one-step numeric algorithm has a high efficiency due to the new, efficient algorithms of the Adomian polynomials, and it enables us to easily generate a higher-order numeric scheme such as a 10th-order scheme, while for the Runge-Kutta method, there is no general procedure to generate higher-order numeric solutions. Finally, the method is demonstrated by using the Duffing equation and the pendulum equation.


## KEYWORDS

Adomian Polynomials; Modified Decomposition Method; Adomian-Rach Theorem; Nonlinear Differential Equations; Numeric Solution

## 1. Introduction

The Adomian decomposition method (ADM) [1-4] is a practical technology for solving linear or nonlinear ordinary differential equations, partial differential equations, integral equations, etc. The ADM provides an efficient analytic approximate solution of nonlinear equations, which model real-world applications in engineering and the applied sciences. The Adomian decomposition series has been shown to be equivalent to a Banach-space analog of the Taylor series expansion about the initial solution component function, instead of the classic Taylor series expansion about a constant that is the initial point [5].

The ADM decomposes the pre-existent, unique, analytic solution into a series

$$
\begin{equation*}
u=\sum_{n=0}^{\infty} u_{n}, \tag{1}
\end{equation*}
$$

and decomposes the nonlinearity $N u=f(u)$ into the series of the Adomian polynomials

$$
\begin{equation*}
f(u)=\sum_{n=0}^{\infty} A_{n}, \tag{2}
\end{equation*}
$$

where the Adomian polynomials $A_{n}$, depend on the solution component functions $u_{0}, u_{1}, \cdots, u_{n}$, and are defined by the formula [3]

$$
\begin{equation*}
A_{n}=\left.\frac{1}{n!} \frac{d^{n}}{d \lambda^{n}} f\left(\sum_{k=0}^{\infty} u_{k} \lambda^{k}\right)\right|_{\lambda=0}, n \geq 0 \tag{3}
\end{equation*}
$$

For convenient reference, we list the first five Adomian polynomials

$$
\begin{gathered}
A_{0}=f\left(u_{0}\right) . \\
A_{1}=f^{\prime}\left(u_{0}\right) u_{1} . \\
A_{2}=f^{\prime}\left(u_{0}\right) u_{2}+f^{\prime \prime}\left(u_{0}\right) \frac{u_{1}^{2}}{2!} .
\end{gathered}
$$

$$
\begin{gathered}
A_{3}=f^{\prime}\left(u_{0}\right) u_{3}+f^{\prime \prime}\left(u_{0}\right) u_{1} u_{2}+f^{\prime \prime \prime}\left(u_{0}\right) \frac{u_{1}^{3}}{3!} . \\
A_{4}=f^{\prime}\left(u_{0}\right) u_{4}+f^{\prime \prime}\left(u_{0}\right)\left(\frac{u_{2}^{2}}{2!}+u_{1} u_{3}\right)+f^{\prime \prime \prime}\left(u_{0}\right) \frac{u_{1}^{2} u_{2}}{2!}+f^{(4)}\left(u_{0}\right) \frac{u_{1}^{4}}{4!} .
\end{gathered}
$$

Some algorithms for symbolic programming have since been devised to efficiently generate the Adomian polynomials quickly to high orders, such as in [5-10]. Rach's Rule for the Adomian polynomials reads

$$
\begin{equation*}
A_{n}=\sum_{k=1}^{n} f^{(k)}\left(u_{0}\right) C_{n}^{k}, n \geq 1 \tag{4}
\end{equation*}
$$

where the coefficients $C_{n}^{k}$ are the sums of all possible products of $k$ components from $u_{1}, u_{2}, \cdots, u_{n-k+1}$, whose subscripts sum to $n$, divided by the factorial of the number of repeated subscripts [6].

New, more efficient algorithms and subroutines in MATHEMATICA for fast generation of the one-variable and multi-variable Adomian polynomials to high orders have been provided in [8-10]. Here we list Corollary 3 algorithm [10] for the one-variable Adomian polynomials.

Corollary 3 algorithm [10]:
For $n \geq 1$,

$$
\begin{equation*}
C_{n}^{1}=u_{n} . \tag{5}
\end{equation*}
$$

For $2 \leq k \leq n$,

$$
\begin{equation*}
C_{n}^{k}=\frac{1}{n} \sum_{j=0}^{n-k}(j+1) u_{j+1} C_{n-1-j}^{k-1} \tag{6}
\end{equation*}
$$

Then the Adomian polynomials are given by the formula

$$
A_{n}=\sum_{k=1}^{n} f^{(k)}\left(u_{0}\right) C_{n}^{k}
$$

The recurrence procedure for $C_{n}^{k}$ does not involve the differentiation operator, only requires the operations of addition and multiplication, which is eminently convenient for computer algebra systems.

In 1992, Rach, Adomian and Meyers [11] proposed a modified decomposition method based on the nonlinear transformation of series by the Adomian-Rach theorem [12,13]:

$$
\text { If } u(x)=\sum_{n=0}^{\infty} a_{n}\left(x-x_{0}\right)^{n},
$$

then

$$
\begin{equation*}
f(u(x))=\sum_{n=0}^{\infty} A_{n}\left(x-x_{0}\right)^{n}, \tag{7}
\end{equation*}
$$

where $A_{n}=A_{n}\left(a_{0}, a_{1}, \cdots, a_{n}\right)$ are the Adomian polynomials in terms of the solution coefficients.
The Rach-Adomian-Meyers modified decomposition method [11] combines the power series solution and the Adomian-Rach theorem [12,13], and has been efficiently applied to solve various nonlinear models [2,14-16].

In this work, higher-order numeric one-step methods are designed for solving nonlinear differential equations based on the Rach-Adomian-Meyers modified decomposition method [11] and the previous research [17].

In the next section, we develop the numeric solution based on the modified decomposition method for nonlinear second-order differential equations, and demonstrate its application.

## 2. Higher-Order Numeric Solutions Based on the Modified Decomposition Method

We consider the IVP for the second-order ODE

$$
\begin{gather*}
\frac{d^{2} u}{d t^{2}}+\alpha(t) \frac{d u}{d t}+\beta(t) u(t)+\gamma(t) f(u)=g(t)  \tag{8}\\
u\left(t_{0}\right)=C_{0}, u^{\prime}\left(t_{0}\right)=C_{1} \tag{9}
\end{gather*}
$$

where $t_{0} \leq t \leq T, \quad \alpha(t), \beta(t), \gamma(t), g(t)$ are specified bounded, analytic functions, and $f$ is an analytic nonli-
near operator.
The modified decomposition method supposes an analytic solution

$$
\begin{equation*}
u(t)=\sum_{m=0}^{\infty} a_{m}\left(t-t_{0}\right)^{m} \tag{10}
\end{equation*}
$$

Then the functions $\alpha(t), \beta(t), \gamma(t), g(t)$ are decomposed into the Taylor expansions

$$
\begin{align*}
& \alpha(t)=\sum_{m=0}^{\infty} \alpha_{m}\left(t-t_{0}\right)^{m},  \tag{11}\\
& \beta(t)=\sum_{m=0}^{\infty} \beta_{m}\left(t-t_{0}\right)^{m},  \tag{12}\\
& \gamma(t)=\sum_{m=0}^{\infty} \gamma_{m}\left(t-t_{0}\right)^{m},  \tag{13}\\
& g(t)=\sum_{m=0}^{\infty} g_{m}\left(t-t_{0}\right)^{m}, \tag{14}
\end{align*}
$$

and the analytic nonlinearity $f(u)$ is decomposed into the Taylor series

$$
\begin{equation*}
f(u)=\sum_{m=0}^{\infty} A_{m}\left(t-t_{0}\right)^{m} \tag{15}
\end{equation*}
$$

where the coefficients $A_{n}=A_{n}\left(a_{0}, a_{1}, \cdots, a_{n}\right)$ are the Adomian polynomials in terms of the solution coefficients $a_{k}$ due to the Adomian-Rach theorem [12,13].

Substituting Equations (10)-(15) in Equations (8), regrouping terms, equating the coefficients of like powers of $\left(t-t_{0}\right)$, and using the initial condition we obtain the recurrence scheme for the solution coefficients $a_{0}=C_{0}, a_{1}=C_{1}$,

$$
\begin{equation*}
a_{m+2}=\frac{1}{(m+1)(m+2)}\left[g_{m}-\sum_{l=0}^{m}\left((m+1-l) \alpha_{l} a_{m+1-l}+\beta_{l} a_{m-l}+\gamma_{l} A_{m-l}\right)\right] \tag{16}
\end{equation*}
$$

where $m \geq 0$ and the $A_{n}=A_{n}\left(a_{0}, a_{1}, \cdots, a_{n}\right)$ are the Adomian polynomials in terms of the coefficients $a_{k}$ for the nonlinear function $f(u)$.

In particular, if $\alpha(t)=\alpha, \beta(t)=\beta$ and $\gamma(t)=\gamma$ are constants, then the recurrence formula becomes $a_{0}=C_{0}, a_{1}=C_{1}$,

$$
\begin{equation*}
a_{m+2}=\frac{1}{(m+1)(m+2)}\left[g_{m}-(m+1) \alpha a_{m+1}-\beta a_{m}-\gamma A_{m}\right], m \geq 0 \tag{17}
\end{equation*}
$$

Further if $g(t)=g$ is also a constant, then the recurrence formula becomes

$$
\begin{gather*}
a_{0}=C_{0}, a_{1}=C_{1}, \\
a_{2}=\frac{1}{2}\left[g-\alpha a_{1}-\beta a_{0}-\gamma A_{0}\right] \\
a_{m+2}=\frac{-1}{(m+1)(m+2)}\left[(m+1) \alpha a_{m+1}+\beta a_{m}+\gamma A_{m}\right], m \geq 1 \tag{18}
\end{gather*}
$$

We denote the $(n+1)$-term approximation of the solution as

$$
\begin{equation*}
\phi_{n+1}\left(t, t_{0}, C_{0}, C_{1}\right)=\sum_{m=0}^{n} a_{m}\left(t-t_{0}\right)^{m} \tag{19}
\end{equation*}
$$

We regard $t_{0}, C_{0}, C_{1}$ as three parameters, and generate the numeric solutions by using the ( $n+1$ )-term approximation $\phi_{n+1}$.

Partition the interval $\left[t_{0}, t_{N}\right.$ ] into $t_{0}<t_{1}<\cdots<t_{N}$. Here we consider an equal step-size partition with $h=t_{k}-t_{k-1}$. The numeric solution generated by $\phi_{n+1}$ is of order $n$. We denote the $n$ th-order numeric solution by $u_{k}^{<n>}, k=0,1, \ldots, N$. The one-step recurrence scheme is as follows:

$$
u_{0}^{<n>}=C_{0}, \dot{u}_{0}=C_{1}
$$

$$
\begin{gathered}
u_{k}^{<n>}=\phi_{n+1}\left(t_{k}, t_{k-1}, u_{k-1}^{<n>}, \dot{u}_{k-1}\right) \\
=u_{k-1}^{<n \gg}+\dot{u}_{k-1} h+\sum_{m=2}^{n} a_{m}^{(k-1)} h^{m}, \\
\dot{u}_{k}=\left.\frac{d}{d t} \phi_{n+1}\left(t, t_{k-1}, u_{k-1}^{<n>}, \dot{u}_{k-1}\right)\right|_{t=t_{k}} \\
=\dot{u}_{k-1}+\sum_{m=2}^{n} m a_{m}^{(k-1)} h^{m-1}, \\
\mathrm{k}=1,2, \cdots, N,
\end{gathered}
$$

where $a_{m}^{(0)}$ are the $a_{m}$ in (16), and for $k=2, \cdots, N, a_{m}^{(k-1)}, m=2,3, \cdots, n$, are determined by a recursion similar to (16) with $a_{0}^{(k-1)}=u_{k-1}^{<n>}$ and $a_{1}^{(k-1)}=\dot{u}_{k-1}$.

Example 1. Consider the IVP for the Duffing equation

$$
\begin{gathered}
\frac{\mathrm{d}^{2} u}{\mathrm{~d} t^{2}}-u+2 u^{3}=-\sin t \sin 2 t, \\
u(0)=1, u^{\prime}(0)=0
\end{gathered}
$$

The IVP has the exact solution $u^{*}(t)=\cos t$. The Adomian polynomials for the nonlinearity $f(t)=u^{3}$ are

$$
\begin{gathered}
A_{0}=a_{0}^{3}, \\
A_{1}=3 a_{0}^{2} a_{1}, \\
A_{2}=3 a_{0} a_{1}^{2}+3 a_{0}^{2} a_{2}, \\
A_{3}=a_{1}^{3}+6 a_{0} a_{1} a_{2}+3 a_{0}^{2} a_{3}, \\
A_{4}=3 a_{1}^{2} a_{2}+3 a_{0} a_{2}^{2}+6 a_{0} a_{1} a_{3}+3 a_{0}^{2} a_{4},
\end{gathered}
$$

The 8th-order numeric solutions on the interval $[0,45]$ are plotted in Figure 1 with the step-size $h=0.5$. The numeric solution is suitable for a larger domain as the order increases.

Example 2. Consider the IVP for the pendulum equation

$$
\frac{d^{2} u}{d t^{2}}+25 \sin u=0, u(0)=0, u^{\prime}(0)=9
$$

The exact solution can be expressed in terms of a Jacobi elliptic function as

$$
u^{*}(t)=2 \arcsin \left(\frac{9}{10} \operatorname{sn}\left(5 t, \frac{81}{100}\right)\right)
$$

The Adomian polynomials in terms of the decomposition coefficients $a_{k}$ for the sinusoidal nonlinearity $\sin u$ are


Figure 1. The exact solution (solid line) and the 8th-order numeric solution on [0,45] with $\mathbf{h}=\mathbf{0 . 5}$ (dots).

$$
\begin{gathered}
A_{0}=\sin a_{0} \\
A_{1}=a_{1} \cos a_{0} \\
A_{2}=a_{2} \cos a_{0}-\frac{a_{1}^{2}}{2} \sin a_{0} \\
A_{3}=-\frac{a_{1}^{3}}{6} \cos a_{0}-a_{1} a_{2} \sin a_{0}-a_{3} \cos a_{0}
\end{gathered}
$$

Using the initial conditions, the coefficients of solution series are calculated to be

$$
\begin{aligned}
& a_{0}=0, a_{1}=9, a_{2}=0, a_{3}=-75 / 2 \\
& a_{4}=0, a_{5}=795 / 4, \ldots
\end{aligned}
$$

The 5-term, 10-term and 20-term approximations $\phi_{5}(t, 0,0,9), \quad \phi_{10}(t, 0,0,9), \quad \phi_{20}(t, 0,0,9)$ are plotted in Figure 2. It is shown that the decomposition solution has a radius of convergence of more than 0.2.

The 5-term approximation $\phi_{5}\left(t, t_{0}, C_{0}, C_{1}\right)$ under the general initial conditions are calculated to be

$$
\begin{aligned}
\phi_{5}\left(t, t_{0}, C_{0}, C_{1}\right)= & \left.C_{0}+C_{1}\left(t-t_{0}\right)-\frac{1}{2} \gamma\left(t-t_{0}\right)^{2} \sin C_{0}\right)\left(-\frac{1}{6} \gamma C_{1}\left(t-t_{0}\right)^{3} \text { со }\left(C_{0}\right)\right. \\
& +\frac{1}{24} \gamma\left(t-t_{0}\right)^{4} \sin \left(C_{0}\right)\left(\gamma \cos \left(C_{0}\right)+C_{1}^{2}\right)
\end{aligned}
$$

The 4th-order and 5th-order numeric solutions on the interval [ 0,6 ] with $\mathrm{h}=0.1$ are plotted in Figures 3 and 4, respectively. The 9th-order numeric solutions on the interval [0,10] with $h=0.2$ are plotted in Figure 5. We observe that the higher-order numeric solutions permit a larger step-size, and enlarge the effective region.


Figure 2. The exact solution $u^{*}(t)$ (solid line), the 5-term approximation $\phi_{s}(t, 0,0,9)$ (dot line), the 10-term approximation $\phi_{10}(\mathbf{t}, \mathbf{0 , 0 , 9})$ (dash line) and the $\mathbf{2 0 - t e r m}$ approximation $\boldsymbol{\phi}_{20}(\mathbf{t}, \mathbf{0 , 0 , 9}$ ) (dot-dash line).


Figure 3. The exact solution (solid line) and the 4 th-order numeric solution on $[0,6]$ with $h=0.1$ (dots).


Figure 4. The exact solution (solid line) and the 5 th-order numeric solution on $[0,6]$ with $\mathbf{h}=\mathbf{0 . 1}$ (dots).


Figure 5. The exact solution (solid line) and the 9th-order numeric solution on [0,10] with $\mathbf{h}=0.2$ (dots).

## 3. Conclusions

We have developed higher-order numeric solutions for nonlinear differential equations based on the Rach-Adomian-Meyers modified decomposition method. Due to the new, efficient algorithms of the Adomian polynomials, the one-step numeric algorithm has a high efficiency, and permits us to easily generate a higher-order numeric scheme such as a 10th-order scheme, while for the Runge-Kutta method, there is no general procedure to generate higher-order numeric solutions. We demonstrated the presented numeric method by two nonlinear physical models.

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