

Higher-Order Numeric Solutions for Nonlinear Systems Based on the Modified Decomposition Method

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ABSTRACT

Higher-order numeric solutions for nonlinear differential equations based on the Rach-Adomian-Meyers modified decomposition method are designed in this work. The presented one-step numeric algorithm has a high efficiency due to the new, efficient algorithms of the Adomian polynomials, and it enables us to easily generate a higher-order numeric scheme such as a 10th-order scheme, while for the Runge-Kutta method, there is no general procedure to generate higher-order numeric solutions. Finally, the method is demonstrated by using the Duffing equation and the pendulum equation.

KEYWORDS

Adomian Polynomials; Modified Decomposition Method; Adomian-Rach Theorem; Nonlinear Differential Equations; Numeric Solution

1. Introduction

The Adomian decomposition method (ADM) [1-4] is a practical technology for solving linear or nonlinear ordinary differential equations, partial differential equations, integral equations, etc. The ADM provides an efficient analytic approximate solution of nonlinear equations, which model real-world applications in engineering and the applied sciences. The Adomian decomposition series has been shown to be equivalent to a Banach-space analog of the Taylor series expansion about the initial solution component function, instead of the classic Taylor series expansion about a constant that is the initial point [5].

The ADM decomposes the pre-existent, unique, analytic solution into a series

$$u = \sum_{n=0}^{\infty} u_n,\tag{1}$$

and decomposes the nonlinearity Nu = f(u) into the series of the Adomian polynomials

$$f(u) = \sum_{n=0}^{\infty} A_n, \tag{2}$$

where the Adomian polynomials A_n , depend on the solution component functions u_0, u_1, \dots, u_n , and are defined by the formula [3]

$$A_n = \frac{1}{n!} \frac{d^n}{d\lambda^n} f\left(\sum_{k=0}^{\infty} u_k \lambda^k\right) \bigg|_{\lambda=0}, n \ge 0.$$
(3)

For convenient reference, we list the first five Adomian polynomials

$$A_0 = f(u_0).$$

$$A_1 = f'(u_0)u_1.$$

$$A_2 = f'(u_0)u_2 + f''(u_0)\frac{u_1^2}{2!}.$$

$$\begin{split} A_3 &= f'(u_0)u_3 + f''(u_0)u_1u_2 + f'''(u_0)\frac{u_1^3}{3!}.\\ A_4 &= f'(u_0)u_4 + f''(u_0)(\frac{u_2^2}{2!} + u_1u_3) + f'''(u_0)\frac{u_1^2u_2}{2!} + f^{(4)}(u_0)\frac{u_1^4}{4!}. \end{split}$$

Some algorithms for symbolic programming have since been devised to efficiently generate the Adomian polynomials quickly to high orders, such as in [5-10]. Rach's Rule for the Adomian polynomials reads

$$A_n = \sum_{k=1}^n f^{(k)}(u_0) C_n^k, \, n \ge 1.$$
(4)

where the coefficients C_n^k are the sums of all possible products of k components from $u_1, u_2, \dots, u_{n-k+1}$, whose subscripts sum to n, divided by the factorial of the number of repeated subscripts [6].

New, more efficient algorithms and subroutines in MATHEMATICA for fast generation of the one-variable and multi-variable Adomian polynomials to high orders have been provided in [8-10]. Here we list Corollary 3 algorithm [10] for the one-variable Adomian polynomials.

Corollary 3 algorithm [10]:

For $n \ge 1$,

$$C_n^1 = u_n. (5)$$

For $2 \le k \le n$,

$$C_n^k = \frac{1}{n} \sum_{j=0}^{n-k} (j+1)u_{j+1} C_{n-1-j}^{k-1}.$$
(6)

Then the Adomian polynomials are given by the formula

$$A_n = \sum_{k=1}^n f^{(k)}(u_0) C_n^k.$$

The recurrence procedure for C_n^k does not involve the differentiation operator, only requires the operations of addition and multiplication, which is eminently convenient for computer algebra systems.

In 1992, Rach, Adomian and Meyers [11] proposed a modified decomposition method based on the nonlinear transformation of series by the Adomian-Rach theorem [12,13]:

If
$$u(x) = \sum_{n=0}^{\infty} a_n (x - x_0)^n$$
,

then

$$f(u(x)) = \sum_{n=0}^{\infty} A_n (x - x_0)^n,$$
(7)

where $A_n = A_n(a_0, a_1, \dots, a_n)$ are the Adomian polynomials in terms of the solution coefficients.

The Rach-Adomian-Meyers modified decomposition method [11] combines the power series solution and the Adomian-Rach theorem [12,13], and has been efficiently applied to solve various nonlinear models [2,14-16]. In this work, higher-order numeric one-step methods are designed for solving nonlinear differential equations

based on the Rach-Adomian-Meyers modified decomposition method [11] and the previous research [17].

In the next section, we develop the numeric solution based on the modified decomposition method for nonlinear second-order differential equations, and demonstrate its application.

2. Higher-Order Numeric Solutions Based on the Modified Decomposition Method

We consider the IVP for the second-order ODE

$$\frac{d^2u}{dt^2} + \alpha(t)\frac{du}{dt} + \beta(t)u(t) + \gamma(t)f(u) = g(t),$$
(8)

$$u(t_0) = C_0, u'(t_0) = C_1,$$
(9)

where $t_0 \le t \le T$, $\alpha(t), \beta(t), \gamma(t), g(t)$ are specified bounded, analytic functions, and f is an analytic nonli-

near operator.

The modified decomposition method supposes an analytic solution

$$u(t) = \sum_{m=0}^{\infty} a_m (t - t_0)^m.$$
 (10)

Then the functions $\alpha(t), \beta(t), \gamma(t), g(t)$ are decomposed into the Taylor expansions

$$\alpha(t) = \sum_{m=0}^{\infty} \alpha_m (t - t_0)^m, \tag{11}$$

$$\beta(t) = \sum_{m=0}^{\infty} \beta_m (t - t_0)^m,$$
(12)

$$\gamma(t) = \sum_{m=0}^{\infty} \gamma_m (t - t_0)^m, \tag{13}$$

$$g(t) = \sum_{m=0}^{\infty} g_m (t - t_0)^m,$$
(14)

and the analytic nonlinearity f(u) is decomposed into the Taylor series

$$f(u) = \sum_{m=0}^{\infty} A_m (t - t_0)^m,$$
(15)

where the coefficients $A_n = A_n(a_0, a_1, \dots, a_n)$ are the Adomian polynomials in terms of the solution coefficients a_k due to the Adomian-Rach theorem [12,13].

Substituting Equations (10)-(15) in Equations (8), regrouping terms, equating the coefficients of like powers of $(t - t_0)$, and using the initial condition we obtain the recurrence scheme for the solution coefficients $a_0 = C_0$, $a_1 = C_1$,

$$a_{m+2} = \frac{1}{(m+1)(m+2)} [g_m - \sum_{l=0}^m ((m+1-l)\alpha_l a_{m+l-l} + \beta_l a_{m-l} + \gamma_l A_{m-l})],$$
(16)

where $m \ge 0$ and the $A_n = A_n(a_0, a_1, \dots, a_n)$ are the Adomian polynomials in terms of the coefficients a_k for the nonlinear function f(u).

In particular, if $\alpha(t) = \alpha$, $\beta(t) = \beta$ and $\gamma(t) = \gamma$ are constants, then the recurrence formula becomes $a_0 = C_0$, $a_1 = C_1$,

$$a_{m+2} = \frac{1}{(m+1)(m+2)} [g_m - (m+1)\alpha a_{m+1} - \beta a_m - \gamma A_m], \ m \ge 0.$$
(17)

Further if g(t) = g is also a constant, then the recurrence formula becomes

$$a_{0} = C_{0}, a_{1} = C_{1},$$

$$a_{2} = \frac{1}{2} [g - \alpha a_{1} - \beta a_{0} - \gamma A_{0}],$$

$$a_{m+2} = \frac{-1}{(m+1)(m+2)} [(m+1)\alpha a_{m+1} + \beta a_{m} + \gamma A_{m}], m \ge 1.$$
(18)

We denote the (n + 1)-term approximation of the solution as

$$\phi_{n+1}(t,t_0,C_0,C_1) = \sum_{m=0}^n a_m (t-t_0)^m.$$
(19)

We regard t_0, C_0, C_1 as three parameters, and generate the numeric solutions by using the (n + 1)-term approximation ϕ_{n+1} .

Partition the interval $[t_0, t_N]$ into $t_0 < t_1 < \cdots < t_N$. Here we consider an equal step-size partition with $h = t_k - t_{k-1}$. The numeric solution generated by ϕ_{n+1} is of order *n*. We denote the *n*th-order numeric solution by $u_k^{<n>}$, $k = 0, 1, \dots, N$. The one-step recurrence scheme is as follows:

$$u_0^{} = C_0, \dot{u}_0 = C_1,$$

$$u_{k}^{\langle n \rangle} = \phi_{n+1}(t_{k}, t_{k-1}, u_{k-1}^{\langle n \rangle}, \dot{u}_{k-1})$$
$$= u_{k-1}^{\langle n \rangle} + \dot{u}_{k-1}h + \sum_{m=2}^{n} a_{m}^{(k-1)}h^{m},$$
$$\dot{u}_{k} = \frac{d}{dt}\phi_{n+1}(t, t_{k-1}, u_{k-1}^{\langle n \rangle}, \dot{u}_{k-1})\Big|_{t=t_{k}}$$
$$= \dot{u}_{k-1} + \sum_{m=2}^{n} ma_{m}^{(k-1)}h^{m-1},$$

 $k = 1, 2, \dots, N,$ where $a_m^{(0)}$ are the a_m in (16), and for $k = 2, \dots, N, a_m^{(k-1)}, m = 2, 3, \dots, n$, are determined by a recursion similar to (16) with $a_0^{(k-1)} = u_{k-1}^{(n)}$ and $a_1^{(k-1)} = \dot{u}_{k-1}$. **Example 1.** Consider the IVP for the Duffing equation

$$\frac{d^2 u}{dt^2} - u + 2u^3 = -\sin t \sin 2t,$$

$$u(0) = 1, u'(0) = 0.$$

The IVP has the exact solution $u^*(t) = \cos t$. The Adomian polynomials for the nonlinearity $f(t) = u^3$ are

$$A_{0} = a_{0}^{3},$$

$$A_{1} = 3a_{0}^{2}a_{1},$$

$$A_{2} = 3a_{0}a_{1}^{2} + 3a_{0}^{2}a_{2},$$

$$A_{3} = a_{1}^{3} + 6a_{0}a_{1}a_{2} + 3a_{0}^{2}a_{3},$$

$$A_{4} = 3a_{1}^{2}a_{2} + 3a_{0}a_{2}^{2} + 6a_{0}a_{1}a_{3} + 3a_{0}^{2}a_{4},$$
...

The 8th-order numeric solutions on the interval [0,45] are plotted in Figure 1 with the step-size h = 0.5. The numeric solution is suitable for a larger domain as the order increases.

Example 2. Consider the IVP for the pendulum equation

$$\frac{d^2u}{dt^2} + 25\sin u = 0, u(0) = 0, u'(0) = 9.$$

The exact solution can be expressed in terms of a Jacobi elliptic function as

$$u^{*}(t) = 2 \arcsin(\frac{9}{10} \operatorname{sn}(5t, \frac{81}{100})).$$

The Adomian polynomials in terms of the decomposition coefficients a_k for the sinusoidal nonlinearity sin*u* are

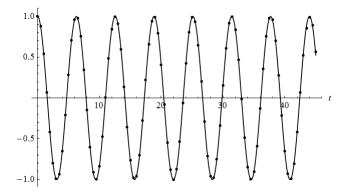


Figure 1. The exact solution (solid line) and the 8th-order numeric solution on [0,45] with h = 0.5 (dots).

$$A_{0} = \sin a_{0},$$

$$A_{1} = a_{1} \cos a_{0},$$

$$A_{2} = a_{2} \cos a_{0} - \frac{a_{1}^{2}}{2} \sin a_{0},$$

$$A_{3} = -\frac{a_{1}^{3}}{6} \cos a_{0} - a_{1}a_{2} \sin a_{0} - a_{3} \cos a_{0},$$

...

Using the initial conditions, the coefficients of solution series are calculated to be

$$a_0 = 0, a_1 = 9, a_2 = 0, a_3 = -75 / 2$$

 $a_4 = 0, a_5 = 795 / 4, \dots$

The 5-term, 10-term and 20-term approximations $\phi_5(t,0,0,9)$, $\phi_{10}(t,0,0,9)$, $\phi_{20}(t,0,0,9)$ are plotted in **Figure 2**. It is shown that the decomposition solution has a radius of convergence of more than 0.2. The 5-term approximation $\phi_5(t,t_0,C_0,C_1)$ under the general initial conditions are calculated to be

$$\phi_{5}(t,t_{0},C_{0},C_{1}) = C_{0} + C_{1}(t-t_{0}) - \frac{1}{2}\gamma(t-t_{0})^{2}\sin C_{0}(t-t_{0})^{3}\cos (C_{0}) + \frac{1}{24}\gamma(t-t_{0})^{4}\sin(C_{0})(\gamma\cos(C_{0}) + C_{1}^{2}).$$

The 4th-order and 5th-order numeric solutions on the interval [0,6] with h = 0.1 are plotted in Figures 3 and 4, respectively. The 9th-order numeric solutions on the interval [0,10] with h = 0.2 are plotted in Figure 5. We observe that the higher-order numeric solutions permit a larger step-size, and enlarge the effective region.

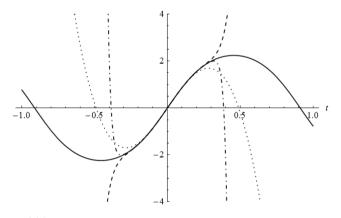


Figure 2. The exact solution $u^*(t)$ (solid line), the 5-term approximation $\phi_0(t,0,0,9)$ (dot line), the 10-term approximation $\phi_0(t,0,0,9)$ (dot line), the 10-term approximation $\phi_0(t,0,0,9)$ (dot-dash line).

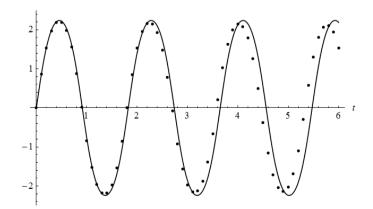


Figure 3. The exact solution (solid line) and the 4th-order numeric solution on [0,6] with h = 0.1 (dots).

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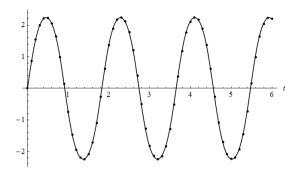


Figure 4. The exact solution (solid line) and the 5th-order numeric solution on [0,6] with h = 0.1 (dots).

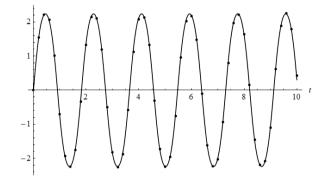


Figure 5. The exact solution (solid line) and the 9th-order numeric solution on [0,10] with h = 0.2 (dots).

3. Conclusions

We have developed higher-order numeric solutions for nonlinear differential equations based on the Rach-Adomian-Meyers modified decomposition method. Due to the new, efficient algorithms of the Adomian polynomials, the one-step numeric algorithm has a high efficiency, and permits us to easily generate a higher-order numeric scheme such as a 10th-order scheme, while for the Runge-Kutta method, there is no general procedure to generate higher-order numeric solutions. We demonstrated the presented numeric method by two nonlinear physical models.

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