

Algebras of Hamieh and Abbas Used in the Dirac Equation*

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ABSTRACT

Hamieh and Abbas [1] propose using a 3-dimensional real algebra in a solution of the Dirac equation. We show that this algebra, denoted by $G\mathbb{C}$, belongs to a large class of quadratic Jordan algebras with subalgebras isomorphic to the complex numbers and that the spinor matrices associated with the solution of the Dirac equation generate a six-dimensional real noncommutative Jordan algebra.

Keywords: Dirac Equation; Jordan Algebra; Quadratic Algebra

Non-associative algebras have long been used in the mathematical description of physical phenomena; first appearing as the “ r number algebra” in the seminal paper by Jordan, Wigner and von Neumann [2] of 1934. The r number algebra became known as a Jordan algebra from a 1946 paper by Albert [3]. The interested reader is referred to the books on non-associative algebras in physics Löhmas, Paal and Sorgsepp [4], Okubo [5]. A concise history of non-associative algebra is to be found in Tomber [6]; the standard introduction to non-associative algebra is the book by Schafer [7].

Hamieh and Abbas [1] present a “description of an algebra which can be used in a possible extension of local complex quantum field theories”. We further expand their description and show that these algebras are Type D Jordan algebras (see Jacobson [8]).

We construct a large family of quadratic Jordan algebras that contains the three-dimensional real algebra, the so called $G\mathbb{C}$ algebra, the generalized complex numbers, of Hamieh and Abbas [1], and show that the spinor matrices that arises from using the $G\mathbb{C}$ in a formulation of the Dirac equation generate a six-dimensional non-commutative quadratic Jordan algebra.

1. Introduction

Let \mathfrak{A} be an algebra over a field F not of characteristic two. The associator is a trilinear mapping

$$(x, y, z) = (xy)z - x(yz)$$

of $\mathfrak{A} \times \mathfrak{A} \times \mathfrak{A}$ into \mathfrak{A} that measures the lack of asso-

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ciativity in \mathfrak{A} .

One scheme of classifying nonassociative algebras involves placing conditions on the associator of certain sets of elements. Some of the better known algebras are:

1) Alternative algebras. In this variety of algebras, all elements x and y satisfy

$$(x, x, y) = (x, y, y) = 0$$

for all elements x and y . The octonion division ring is an alternative algebra. An interesting variation is pseudo-octonion algebra (Okubo [5,9]).

2) Jordan algebras. These are commutative algebras in which all x and y satisfy

$$(x, y, x^2) = 0.$$

A Type D Jordan algebra is the Jordan algebra of the symmetric bilinear form q on a vector space \mathfrak{B} . Albert [3] has shown that any algebra of Type D has a basis $\{e, b_1, b_2, \dots, b_n\}$ with multiplication given by

$$eb_i = b_i e = b_i, \text{ for all } 1 \leq i \leq n,$$

$$b_i b_j = \delta_{ij} \alpha_i e \text{ for } 1 \leq i, j \leq n.$$

The algebra will be semisimple if $\alpha_i \neq 0$ for all $1 \leq i \leq n$.

3) Noncommutative Jordan algebras. A generalization of the alternative and Jordan algebras that requires all x and y satisfy a generalization of the commutative law

$$(x, y, x) = 0,$$

that is, the algebras are flexible, and

$$(x, y, x^2) = 0.$$

The book by Zhevlakov, Slin’ko, Shestakov and Shirshov [10] provides a detailed analysis of the alternative and Jordan rings.

The above algebras are all power associative since each element a generates an associative subalgebra; equivalently, $a^m a^n = a^{n+m}$ for positive integers m, n . In any power associative algebra \mathfrak{A} with unit element we can introduce the series

$$e^x = \sum_{i=0}^{\infty} \frac{x^i}{i!}$$

for $x \in \mathfrak{A}$ ignoring the question of convergence.

An algebra \mathfrak{A} over a field F is called quadratic if, for every x in \mathfrak{A}

$$x^2 - 2t(x) + q(x)e = 0$$

where $t(x), q(x)$ are in F and e is the identity of \mathfrak{A} . The quantities $t(x)$ and $q(x)$ are called the trace and norm of the element x , respectively. The trace is a linear functional on \mathfrak{A} see Schafer [7]. The norm $q(x)$ defines a symmetric bilinear form $q(x, y)$ on \mathfrak{A} via

$$q(x, y) = q(x + y) - q(x) - q(y).$$

Say $q(x)$ is nondegenerate if $q(x, y)$ is. Any quadratic algebra is power associative and any flexible, quadratic algebra is a noncommutative Jordan algebra.

A quadratic algebra \mathfrak{A} is flexible if and only if the trace is associative; that is, $t((xy)z) = t(x(yz))$ for all x, y, z in \mathfrak{A} . If \mathfrak{A} is flexible then the mapping $x \rightarrow \bar{x} = 2t(x)e - x$ is an involution in \mathfrak{A} (see Braun and Koecher [11], p. 216).

Lemma 1. *The Hamiltonian division ring is a quadratic algebra.*

Proof. Let $x = a + bI + cJ + dK$ be an element of the Hamiltonian division ring. Direct computation shows that

$$x^2 - 2ax + a^2 + b^2 + c^2 + d^2 = 0. \quad \square$$

Example 1. *The octonion division ring is a quadratic algebras.*

Example 2. *Domokos and Kövesi-Domokos [12] propose a quadratic algebra, the “algebra of color” as a candidate for the algebra obeyed by a quantized field describing quarks and leptons (see also Wene [13,14], and Schafer [15]).*

2. Construction of the Algebras

The elements of the algebra $G\mathbb{C}$ are the elements of the real vector space with basis $\{e, I, J\}$. The addition is the vector space addition and multiplication is defined by $IJ = JI = 0, I^2 = J^2 = -e, e$ is the identity and the distributive laws. We note that the algebra is commutative and has divisors of zero.

An immediate generalization of this algebra has a basis $\{e, b^i, i = 1, \dots, n\}, n \geq 2$ over the field \mathbb{R} of real

numbers and multiplication defined by $b^i b^j = -\delta_{ij}e$ where e is the identity. For want of a better name called these the Abbas algebras. As noted above, these algebras are Type D Jordan algebras. Note that the $G\mathbb{C}$ algebra is the construction for $n = 2$; the results for the Abbas algebras apply to the $G\mathbb{C}$. Each Abbas algebra contains a copy of the complex numbers.

Lemma 2. *The Abbas algebras are quadratic algebras.*

Proof. Let H denote a Abbas algebra. Then if $x \in H, x = \alpha_0 e + \alpha_i b^i$, Einstein summation convention where $i = 1, 2, \dots, n$. Then

$$\begin{aligned} x^2 &= \alpha_0 x + \alpha_0 \alpha_i b^i - (b^i b^i) \\ -2\alpha_0 x &= -2\alpha_0^2 - 2\alpha_0 \alpha_i b^i \end{aligned}$$

Adding both sides gives

$$x^2 - 2\alpha_0 x = -\alpha_0^2 - (\alpha_i \alpha_i)$$

and we see that $t(x) = \alpha_0$ and $q(x) = \alpha_0^2 + (\alpha_i \alpha_i)$.

A commutative quadratic algebra will be a Jordan algebra. \square

Since the algebra is commutative the trace is associative; the norm is symmetric.

Lemma 3. *The norm of a Abbas algebra is nondegenerate.*

Proof. Let H denote a Abbas algebra. Then if $x \in H, x = \alpha_0 e + \alpha_i b^i$ is arbitrary and $d = \delta_0 e + \delta_i b^i$ is fixed, then

$$\begin{aligned} q(d, x) &= q(d + x) - q(d) - q(x) \\ &= (\delta_0 + \alpha_0)^2 + \sum_{i=1}^n (\delta_i + \alpha_i)^2 \\ &\quad - (\delta_0^2 + \delta_i \delta_i) - (\alpha_0^2 + \alpha_i \alpha_i) \\ q(d, x) &= 2\delta_0 \alpha_0 + 2\delta_i \alpha_i \end{aligned}$$

\square

3. Special $G\mathbb{C}$ Algebras

Hamieh and Abbas [1] pass to a representation of the point $q = ae + bI + cJ$ of the algebra $G\mathbb{C}$ in spherical coordinates, $a = r \cos(\theta), b = r \sin(\theta) \cos(\varphi)$ and $c = r \sin(\theta) \sin(\varphi)$. The subalgebras, called special $G\mathbb{C}$ algebras and denoted by $SG\mathbb{C}$ are the subalgebras spanned by all elements in which the “azimutal phase angle φ is constant”. Each of these subalgebras is (isomorphic to) the complex numbers.

Lemma 4. *The algebra $G\mathbb{C}$ is isomorphic to an algebra of two by two matrices*

$$\begin{bmatrix} a & bI + cJ \\ bI + J & a \end{bmatrix}$$

under the usual matrix operation of addition and multiplication.

Proof. The straight forward verification that the mapping $\theta(ae + bI + cJ) = \begin{bmatrix} a & bI + cJ \\ bI + J & a \end{bmatrix}$ is an isomorphism is left to the reader. \square

Lemma 5. *Each of the algebras SGC is isomorphic to the complex numbers.*

Proof. We note that if $b = r \sin(\theta) \cos(\varphi)$ then $c = b \left(\frac{\sin(\varphi)}{\cos(\varphi)} \right)$ if $\cos(\varphi) \neq 0$. If $\varphi = \frac{\pi}{2}$ or $\frac{3\pi}{2}$, then $q = a + cJ$ and the subalgebra SGC is (isomorphic to) the complex numbers. Otherwise, $q = ae + bI + bsJU$ or $q = ae + b(I + sJ)$ for some $s \in \mathbb{R}$. Let $X = \frac{I + sJ}{\sqrt{1 + s^2}}$, then

$$X^2 = \left(\frac{I + sJ}{\sqrt{1 + s^2}} \right)^2 = \frac{-(1 + s^2)}{1 + s^2} e = -e.$$

The multiplication, using the basis $\{e, X\}$ will be given by

$$(ae + bX)(ce + dX) = (ac - bd)e + (ad + bc)X.$$

\square

4. The Spinor Matrices

The classical reference on spinors and wave equations is the book by Corson [16].

The associator spinor matrices are

$$C_t = \begin{pmatrix} 0 & J \\ J & 0 \end{pmatrix}, \quad C_x = \begin{pmatrix} 0 & -J \\ J & 0 \end{pmatrix},$$

$$C_y = \begin{pmatrix} 0 & -I \\ I & 0 \end{pmatrix}, \quad C_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$xy = \begin{pmatrix} -bs - ct & bvI + cvJ \\ dsI + dtJ & dv - ey - fz \end{pmatrix}$$

$$(xy)x = \begin{pmatrix} -ebv + fcv & (-bs - ct)(bI + cJ) + d[(bvI + cvJ)] \\ (dv - ey - fz)(eI + fJ) & bds + cdt + d(dv - ey - fz) \end{pmatrix}$$

$$yx = \begin{pmatrix} -ey - fz & (xb + dy)I + (cx + dz)J \\ evI + fvJ & dv - bs - ct \end{pmatrix}$$

$$x(yx) = \begin{pmatrix} -bev - cfv & (dv - bs - ct)(bI + cJ) \\ (-ey - fz)(eI + fJ) + (devI + dfvJ) & d(dv - bs - ct) - edy - fdz \end{pmatrix} \quad \square$$

Theorem 1. *The algebra $SP(6)$ is a quadratic non-commutative Jordan algebra.*

5. The Dirac Equation

We proceed as in Hamieh and Abbas [1]. The Dirac

where $1, I, J \in GC$. Denoting the 2 by 2 identity matrix by I_2 , these matrices satisfy

$$C_i^2 = -I_2 \text{ and } \{C_\mu, C_\nu\} = C_\mu C_\nu + C_\nu C_\mu = 2\delta_{\mu\nu} I_2$$

for $\mu, \nu = x, y, z$.

The spinor matrices generate a 6-dimensional real algebra with elements

$$\left\{ \begin{pmatrix} a & bI + cJ \\ eI + fJ & d \end{pmatrix} \mid a, b, c, d, e, f \in \mathbb{R}, I, J \in GC \right\}$$

that contains the matrix representation of the GC algebra. Denote this algebra by $SP(6)$.

Lemma 6. *The algebra $SP(6)$ is a quadratic algebra.*

Proof. If $x = \begin{pmatrix} a & bI + cJ \\ eI + fJ & d \end{pmatrix}$ is an element of $SP(6)$, then

$$x^2 = \begin{pmatrix} a^2 - be - cf & (a + d)(bI + cJ) \\ (a + d)(fI + cJ) & d^2 - be - cf \end{pmatrix}$$

$$-2\left(\frac{1}{2}\right)(a + d)x = -\begin{pmatrix} a^2 + ad & (a + d)(bI + cJ) \\ (a + d)(eI + fJ) & ad + d^2 \end{pmatrix}$$

Adding the left and right sides gives

$$x^2 - 2\left(\frac{1}{2}\right)(a + d)x = \begin{pmatrix} -ad - be - cf & 0 \\ 0 & -ad - cb - cf \end{pmatrix}$$

\square

Lemma 7. *The algebra $SP(6)$ is flexible.*

Proof. Because of the trilinearity of the associator, we can write the elements x and y of the associator (x, y, x)

$$\text{as } x = \begin{pmatrix} 0 & bI + cJ \\ eI + fJ & d \end{pmatrix} \text{ and } y = \begin{pmatrix} 0 & yI + zJ \\ sI + tJ & v \end{pmatrix}.$$

Then

equation over the complex numbers is often written as

$$(i\gamma^\mu \partial_\mu - m)\Psi = 0$$

utilizing the Einstein summation convention for $\mu = x, y, z, t$. A more general form is, setting

$$\gamma^\mu = C_\mu, H\Psi = (C_\mu \partial_\mu)\Psi = m\Psi$$

$$(C_\mu \partial_\mu - m)\Psi = 0,$$

where $\mu = x, y, z, t$.

Upon substituting the matrices for C_μ and simplifying we get

$$\begin{pmatrix} \partial_z - m & \partial_z \\ J\partial_x + I\partial_y + J\partial_t & -\partial_z - m \end{pmatrix} \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix} = 0.$$

In dimensions x and t , the solution is given by

$$\Psi(x, t) = N \begin{pmatrix} E + p \\ m \\ 1 \end{pmatrix} e^{J(px - ct)}$$

p and $E = \pm\sqrt{p^2 + m^2}$ are respect the momentum and energy. N is a normalization factor.

6. Conclusion

We have shown that the GC algebra belongs to a large class of Jordan algebras and have examined a few of the algebraic properties of these algebras and, like the Jordan algebra and the algebra of color, there is a very rich mathematical structure to further explore.

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