

Controllable Growth of Ni Nanocrystals Embedded in BaTiO₃/SrTiO₃ Superlattices

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ABSTRACT

 $BaTiO_3/SrTiO_3$ superlattices with embedded Ni nanocrystals (NCs) have been grown on $SrTiO_3(001)$ substrate using laser molecular beam epitaxy (L-MBE). In situ reflection high-energy electron diffraction (RHEED) was employed to investigate the process of lattice strain in the self-organization of Ni NCs and the epitaxial growth of $BaTiO_3/SrTiO_3$ superlattices. The results indicated that the strain from large lattice mismatch drove the self-organization of Ni NCs. Also, the layer-by-layer growth of $BaTiO_3/SrTiO_3$ superlattices and the island growth of Ni NCs were controllable accurately. The fine alternation of the two processes would provide a possible route to engineer controllably the nanocomposite microstructure.

Keywords: Nanocrystal; Superlattices; Self-organization

1. Introduction

Oxide artificial superlattices, especially (001) oriented BaTiO₃/SrTiO₃ superlattices (BTO/STO SLs), have been attracted more attention because their dielectricand ferroelectric properties can be enhanced by controlling the lattice strain along the polarized direction [1-3]. The dielectric matrix with the embedded metal NCs has widely potential application in nonlinear optical and electronic device [4, 5]. Therefore, several methods have been considered to fabricate the BTO and STO-based material, including sol-gel [6], reactive evaporation [7], rf-sputtering [8] and laser ablation methods [9]. However, the controllable fabrication of nanostructure remains the daunting challenge for many deposition methods. There is an excellent method referred to as self-organized growth, in which the strain would drive the three-dimensional (3D) island to form in the lattice mismatched growth process [10, 11]. Especially for the fabrication of quantum dot (QD) structures, the self-organized growth is greatly successful in semiconductor devices, such as InGaAs on GaAs [10], AlN on GaN [12]. Wu et al. have successfully fabricated Co:BTO and Ni:BTO composite film using self-organized method [13-15].

In this paper, the laser molecular beam epitaxy was used to embed the Ni NCs in the epitaxial films of Ba- $TiO_3/SrTiO_3$ superlattices (Ni NCs:BTO/STO SLs). The fabrication of the nanocomposite system is interesting

and significant for both fundamental and application aspects. Such composite films offer the combination of ferroelectric and ferromagnetic characteristics. This magnetoelectric interaction will be particularly strong in systems which simultaneously exhibit both ferromagnetism and ferroelectricity, i.e., in magnetoelectric multiferroics [16-18].

2. Experimental

The Ni NCs:BTO/STO SLs films were grown on STO (001) substrate with STO buffer layer by L-MBE. The acceleration voltage of RHEED electron gun was 25kev and the grazing incidence angle was about 1-3°. The growth process was performed in an ultrahigh-vacuum (UHV) system and the background vacuum was ~ 2 $\times 10^{-6}$ Pa. The schematic diagram of Ni NCs:BTO/STO SLs film was shown in Fig. 1. The experimental parameters were listed in detail in Table 1. Before the deposition, the substrate STO (001) was in situ annealed about 2 hours in order to remove surface contamination. At 10 Pa oxygen pressure, the sample was annealed about 20 minutes after accomplishing the BTO layer. The samples were characterized by XRD with Cu ka radiation and transmission high-resolution electron microscopy (HRTEM).

3. Results and discussion

Fig.2 shows the oscillation of the RHEED intensity and



Figure 1. A schematic diagram of Ni NCs:BTO/STO SLs film. One deposition cycle involves a number of pulses on the Ni target in ultra-high vacuum, followed by the epitaxial growth of BTO/STO SLs.

 Table 1. Experimental parameters for Ni NCs:BTO/STO

 SLs films fabrication

Background vacuum	~2×10 ⁻⁶ Pa
Working vacuum	~2×10 ⁻⁵ Pa
Target	BaTiO ₃ purity >99.99% SrTiO ₃ purity >99.99% Ni purity >99.99%
Substrate	SrTiO ₃ (001)
Laser pulse frequency	1Hz for SrTiO ₃ deposition 1Hz for BaTiO ₃ deposition 2Hz for Ni deposition
The distance between the target and substrate	5 cm
Annealing condition	650 ℃ 20 min 10Pa O ₂ pressure

diffraction patterns of BTO deposited on STO surface. The diffraction sharp streaks indicate a smooth surface and good crystalline quality [19]. Compared with the unannealed STO substrate [Fig. 2(a)], the diffraction streaks become more obvious after in situ anneal at 650° C [Fig. 2(b)], indicated that the crystallization of STO surface is improved. During the growth processes, the characteristic streaks of BTO became more striking shown in Fig.1(c) and Fig.1 (d). The prefect epitaxial layers of BTO were formed on the STO surface. In this case, the intensity oscillations are obviously observed. It is concluded that the deposited BTO layers are grown with layer-by-layer growth mode.



Figure 2. Variation of RHEED intensity oscillation and patterns along the [100] azimuth of BTO on STO layer: (a) unannealed STO substrate; (b) annealed STO surface; (c) deposited 60s BTO; (d) deposited 120s BTO



Figure 3. Transition of RHEED patterns during the self-organized embedded Ni process: (a) BTO(001) surface; (b) 100 pulses Ni; (c) 500 pulses Ni; (d) 1000 pulses Ni

As indicated in Fig. 3, the self-organized Ni NCs along the [100] azimuth were deposited the new surface of BTO (001). It is obvious that the streaky patterns of BTO disappeared gradually with the increasing Ni pulses, while the spots of Ni NCs become more striking. Additionally, due to the large lattice mismatch between BTO and Ni (>10%), that results in an overall decrease in RHEED intensity. Thus, this case is considered as the formation of 3D islands. The surface-lattice parameter, which is proportional to the inverse of the distance between the different diffraction spots or streaks, can be directly measured from the RHEED patterns [19]. By measuring the vertical spacing and horizontal spacing for the dominant Ni diffraction spots, the lattice constant of Ni in cubic structure was confirmed to 0.354nm in real space. This result approaches to the lattice constant of bulk metal Ni (0.352 nm) in cubic structure [20]. Therefore, it is demonstrated that the deposited Ni maintains metallic property.



Figure 4. Variation of RHEED intensity oscillation during STO layer deposited on the Ni NCs. The insets are the RHEED patterns recorded at 10s, 40s and 120s, respectively.

After the completion of Ni NCs, the next layer is STO. As shown in Fig. 4, at the initial stage of STO deposition, the RHEED intensity decreases instantaneously because of the increasing random distribution of the incoming STO atoms on the surface. Then, the deposited STO atoms and the formed clusters are rearranged on the surface. That induces the increasing RHEED oscillated intensity. For the insets Fig. 4(a)-(d), the dominant Ni spots change to the characteristic streaks of STO with the increasing STO pulses. Meanwhile, the RHEED intensity oscillations are observed again. Those results demonstrate that the STO laver covers the geometrically irregular surface provided self-organized Ni NCs, and is grown with layer-by-layer growth mode. The accurate thickness of every individual layer is controlled by RHEED intensity oscillation which is extremely sensitive to the presence of surface atoms in the top crystalline layer [21, 22]. After the laser ended for some time, the oscillated intensity increased slowly, the lattice relaxation tended to smooth for the surface roughness. Therefore, the Ni NCs are embedded successfully in BTO/STO SLs by the alternation of 3D island Ni NCs and perfect epitaxial BTO/STO SLs. Importantly, the embedded Ni NCs can not disturb the epitaxial growth of BTO/STO SLs.

In order to explain the formation mechanism of 3D island Ni NCs, the changes in the $(a_{Ni}-a_{BTO})/a_{BTO}$ as a function of the pulses, using the spot spacing for Ni NCs and the streak spacing for BTO, were calculated (Fig. 5). Initially, the Ni NCs retain an in-plane lattice parameter identical to that of the BTO layer, and then recover to the value of bulk Ni rapidly. Owing to the large lattice mismatch between Ni and BTO (>10%), Ni NCs were subjected to considerable in-plane tensile strain. Simultaneously, this in-plane straining would shorten the out-of-plane lattice parameter of Ni NCs, leading to the structural transition from cubic to tetragonal. With the deposition being continued, the increasing strain energy

was relieved by the transformation from 2D to 3D islands as the thickness exceeded the critical thickness h_c , in which the lattice was relaxed laterally to the unstrained value, and the critical thickness is in inverse proportional to the lattice mismatch, according to the Matthews-Blakeslee formula. The strain from lattice mismatch acts as a source of driving force for the self-organization of Ni NCs. Therefore such metal NCs growth route is feasible especially in the large lattice mismatch system.



Figure 5. The in-plane lattice relaxation during the self-organization of Ni NCs. Inset an expanded view of the first 90 pulses



Figure 6. The cross-sectional HRTEM image of Ni NCs:BTO/STO SLs nanocomposition film

Fig. 6 is the cross-section HRTEM image of Ni NCs:BTO/STO SLs composition films. The irregular interface of strained Ni NC layers is alternated with 7 times. We estimate the sizes of Ni NCs about 5nm. As shown in Fig. 6, the perfect layer-by-layer epitaxia growth of STO and BTO layers is performed. The separate BTO layer and STO layer have the uniform thickness of 8 nm and 10 nm accordingly. The thickness of

every separate layer is accurately controlled by RHEED intensity oscillation which is extremely sensitive to the presence of surface atoms in the top crystalline layer [3]. The individual BTO layer and STO layer have the uniform thickness of 14nm and 21nm, respectively. According to the results of RHEED monitored, the growth rate of STO and BTO are estimated to be 0.037 ML/s and 0.043 ML/s, where 1 monolayer (ML) corresponds to a layer thickness of 0.781 nm (STO) and 0.806 nm (BTO), the double value of the c-axis lattice constant of bulk STO (0.3905 nm) and BTO (0.4029 nm) [5], respectively. So it is concluded that one period of RHEED intensity oscillation corresponds to the growth of two STO unit cells and two BTO unit cells, respectively. We achieve atomic-level control for the growth of the nanocomposite film.



Figure 7. XRD θ -2 θ scans of the composite films on STO (001) substrate; the inset is an expand view around the BTO (002) peak

Fig. 7 shows the X-ray diffraction θ -2 θ scans of the composite films grown on STO (001) substrate. The characteristic peaks of BTO and STO show that the epitaxial growth is satisfying along basically *c*-axis with the lower diffraction angles. Using STO single-crystal substrate, we have obtained very high quality epitaxial films. Magnifying the BTO (002) peak for the inset, a weak peak is observed beside the leading peak. We suppose that the weak peak is Ni (111). The crystallographic orientation with Ni (111) takes priority of the fcc metal structure. The appearance of the Ni (111) is supposed to utmost compressed plane in order to get minimum surface energy and remain metallic property, so as to give the flat surface to the vicinal separation layer. This way reduces extremely the effect of the strained distortion by the large lattice mismatch between the different interfaces. The results of the XRD characterization are in good agree with RHEED analysis.

4. Conclusion

In summary, Ni NCs:BaTiO₃/SrTiO₃ superlattices were

fabricated by L-MBE and controlled accurately by in situ monitoring RHEED. With the self-organized method, the alternate growth of Ni NCs was embedded successfully in BaTiO₃/SrTiO₃ superlattices and could not disturb the epitaxial growth of BaTiO₃/SrTiO₃ superlattices, in good agree with the analysis of XRD and HRTEM. The results were concluded that the layer-by-layer growth of Ba-TiO₃/SrTiO₃ superlattices and the three-dimensional island growth of Ni NCs. The strain in the composite films from lattice mismatch acts as a source of driving force for the self-organization of Ni NCs. Therefore such self-organized metal NCs growth route is propitious to fabricate three-dimensional quantum dot devices.

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