

# Analytical and Traveling Wave Solutions to the Fifth Order Standard Sawada-Kotera Equation via the Generalized $\exp(-\Phi(\xi))$ -Expansion Method

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# Abstract

In this article, we propose a generalized  $\exp(-\Phi(\xi))$  -expansion method and successfully imple-

ment it to find exact traveling wave solutions to the fifth order standard Sawada-Kotera (SK) equation. The exact traveling wave solutions are established in the form of trigonometric, hyperbolic, exponential and rational functions with some free parameters. It is shown that this method is standard, effective and easily applicable mathematical tool for solving nonlinear partial differential equations arises in the field of mathematical physics and engineering.

# **Keywords**

Generalized  $\exp(-\Phi(\xi))$ -Expansion Method, Fifth Order Standard Sawada-Kotera Equation, Solitons, Periodic Wave Solutions

# **1. Introduction**

The nonlinear partial differential equations (NPDEs) are widely used as model equations to study on the dynamics of many issues concerning in the field of mathematical physics, chemistry and engineering. Therefore, it is necessary to solve them analytically or numerically. Some of the methods to solve them analytically are the

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Hirota method [1], the projective Riccati equation method [2], the homogeneous balance method [3]-[5], the tanh-function method [6]-[10], the exp-function method [11], the improved F-expansion method [12], the  $\exp(-\Phi(\xi))$ -expansion method [13]-[19], sine-cosine method [20], the modified simple equation method [21], the (G'/G)-expansion method [22] [23], the modified extended tanh-function method with Riccati equation [24] and others. In the references [13]-[19], we observe that the  $\exp(-\Phi(\xi))$ -expansion method provides only few solutions according to the nonlinear ordinary differential equation (ODE)  $\Phi'(\xi) = e^{-\Phi(\xi)} + \mu e^{\Phi(\xi)} + \lambda$ . There are a lot of previous literature, where the researchers have investigated exact travelling wave solutions to the NPDEs using the exp $(-\Phi(\xi))$ -expansion method. But there is no work to clarify why the coefficient of  $\exp(-\Phi(\xi))$  to the subsidiary equation is unity. This is the disadvantage of this method. Recently, Hafez and Akbar [25] have developed an exponential expansion method to solve NLPDEs by considering three auxiliary equations  $\Phi'(\xi) = e^{-\Phi(\xi)} + \mu e^{\Phi(\xi)} + \lambda$ ,  $\Phi'(\xi) = -\sqrt{\lambda + \mu (e^{-\Phi(\xi)})^2}$  and  $\Phi'(\xi) = -\lambda e^{-\Phi(\xi)} - \mu e^{\Phi(\xi)}$ . Form this literature, we observe that the NLPDEs has almost similar types of solutions with the help of  $\Phi'(\xi) = -\sqrt{\lambda + \mu \left(e^{-\Phi(\xi)}\right)^2} \text{ and } \Phi'(\xi) = -\lambda e^{-\Phi(\xi)} - \mu e^{\Phi(\xi)}.$  This is a good technique, but it is taking too much time for solving nonlinear evaluation equations. Recently, Hafez and Lu [26] have proposed a generalized  $\exp(-\Phi(\xi))$ -expansion method for investigating the traveling wave solutions of the nonlinear fractional partial differential equations. We have proposed this method for solving NPDEs. On the other hand, through the proposed method, not only re-derives all known solutions in a systemic way but also delivers several exact traveling wave solutions to the NPDEs found by influence of the related free parameters if the nonlinear ODE  $\Phi'(\xi) = \alpha e^{-\Phi(\xi)} + \beta e^{\Phi(\xi)} + \gamma$  is used.

There is an amount of article [27]-[32], where the exact solutions of the standard SK equation are analyzed. Some exact solutions of the standard SK equation have been evaluated using different methods, such as the tanh-coth method [29], generalized the tanh-coth method [30] and the exp-function method [31]. The aim of this article is to establish the advance exponential expansion method and linking the method for evaluating the exact traveling wave solutions to the standard SK equation. The advantage of the proposed method over the other existing methods is that it provides new exact traveling wave solutions to the nonlinear PDEs. This article also shows that the proposed method is very effective and standard technique rather than the  $\exp(-\Phi(\xi))$ -expansion method and exponential expansion method. The  $\exp(-\Phi(\xi))$ -expansion method is only special case of

sion method and exponential expansion method. The  $\exp(-\Phi(\zeta))$ -expansion method is only special case of the proposed advance exponential expansion method.

## 2. The Methodology

The traveling wave solutions of a nonlinear evolution equation in one space dimensional are solutions obtained by assuming

$$u(x,t) = u(\xi), \ \xi = k \, x \pm V \, t, \tag{1}$$

where *k* is the wave number is and *V* is the speed of the wave.

The method consists of writing the traveling wave solution in the following form:

$$u(\xi) = \sum_{i=0}^{N} A_i \left( e^{-\Phi(\xi)} \right)^i, \ A_N \neq 0$$
<sup>(2)</sup>

where the coefficients  $A_i$   $(0 \le i \le N)$  are constants to be determined and  $\Phi = \Phi(\xi)$  satisfies the first order nonlinear ODE:

$$\Phi'(\xi) = \alpha e^{-\Phi(\xi)} + \beta e^{\Phi(\xi)} + \gamma$$
(3)

where  $\alpha$ ,  $\beta$  and  $\gamma$  are arbitrary constants. Furthermore, if we define  $\overline{\Phi} = \Phi + A$ , with corresponding relabeling of the coefficients, then

$$\overline{\alpha} = e^{A}\alpha, \ \beta = e^{-A}\beta, \ \overline{\gamma} = \gamma$$
(4)

This substitution results in a scaling of the coefficients of the expansion for  $u(\xi)$ . Therefore one can set  $\alpha = 1$  and  $\gamma = 0$ . It is notable that the solutions of Equation (3) are given in Ref. [26].

The value of the positive integer *N* can be determined by balancing the highest order derivative term with the nonlinear term of the highest order appearing in the reducing ODE according to the traveling wave variables transformation. By substituting (1) into the reduce ODE by means of Equation (2), we can easily obtain a system of algebraic equations for  $A_i$  ( $0 \le i \le N$ ),  $\alpha$ ,  $\beta$ ,  $\gamma$ , V. With the help of symbolic computation, such as Maple, we can solve the system and obtain the values  $A_i$  ( $0 \le i \le N$ ),  $\alpha$ ,  $\beta$ ,  $\gamma$ , V. The multiple explicit solutions of a nonlinear evaluation equation will be obtained by combining the Equation (1) and the solution of Equation (3).

## 3. Application of the Method

In this section, the application of the proposed method to extract more explicit and exact traveling wave solutions to the fifth order standard SK equation has been presented. The general form of an evolution equation with same scaling properties as the KdV [29] [30] is

$$v_t + v_{xxxxx} + c v v_{xxx} + b v_x v_{xx} + a v^2 v_x = 0$$
(5)

where *a*, *b* and *c* are arbitrary nonzero real parameters and v = v(x,t) is differentiable function.

In this article, we have considered the following standard SK equation [31]:

$$u_t + u_{xxxxx} + 5u \ u_{xxx} + 5u_x u_{xx} + 5u^2 u_x = 0, \tag{6}$$

which is the special case of Equation (5). It is known that the SK equation belongs to the completely integrable hierarchy of higher-order model evolution equations and has many sets of conservation laws [32]-[34]. Here u(x,t) represents the evaluation of surface at any varied natural instances. If we set

$$u(x,t) = u(\xi) \quad \text{and} \quad \xi = x - Vt,$$
(7)

where V is the speed of traveling wave, then we can be reduced the Equation (6) in to the following ODE:

$$-V u' + u^{(5)} + 5 u u''' + 5 u' u'' + 5 u^2 u' = 0$$
(8)

Integrating (8) with respect to  $\xi$ , we obtain

$$C - V u + u''' + 5u u'' + \frac{5}{3}u^3 = 0$$
<sup>(9)</sup>

where prime denotes derivative with respect to  $\xi$  and *C* is the constant of integration. Balancing the higher order derivative of u'''' and the nonlinear term  $u^3$  appearing in (9), we obtain N = 2. Therefore, the Equation (9) has the following form of solution:

$$u(\xi) = a_0 + a_1 e^{-\Phi(\xi)} + a_2 \left(e^{-\Phi(\xi)}\right)^2, \ a_2 \neq 0$$
(10)

By substituting (10) into (9) and using the Equation (3), we obtain a set of algebraic equations by equating the coefficients of  $(e^{-\Phi(\xi)})^i$ ,  $(i = 0, 1, 2, \dots, 6)$ . Here the set of algebraic equations are overlooked for convenience. Solving the resulting set of algebraic equations, the following sets of solutions are obtained.

Set 1: 
$$C = 52a_0\gamma^2\alpha\beta + 6\gamma^4\alpha\beta + 132\gamma^2\beta^2\alpha^2 + 96\alpha^3\beta^3 + \frac{10}{3}a_0^3 + 136a_0\alpha^2\beta^2 + 40\alpha\beta a_0^2 + a_0\gamma^4 + 5\gamma^2 a_0^2,$$
  
 $a_0 = a_0, \ a_1 = -6\alpha\gamma, \ a_2 = -6\alpha^2, \ V = 22\alpha\beta\gamma^2 + 5\gamma^2 a_0 + 76\alpha^2\beta^2 + \gamma^4 + 5a_0^2 + 40\alpha\beta a_0.$   
Set 2:  $C = 32\gamma^2\alpha^2\beta^2 - \frac{128}{3}\alpha^3\beta^3 - 8\gamma^4\alpha\beta - \frac{2}{3}\gamma^6, \ V = -8\alpha\beta\gamma^2 + \gamma^4 + 16\alpha^2\beta^2, \ a_0 = -\gamma^2 - 8\alpha\beta,$ 

 $a_1 = -12\alpha\gamma$ ,  $a_2 = -12\alpha^2$ .

When  $\alpha = 1$ , the fifth order standard SK Equation (6) has the following explicit solutions: For **Set 1**:

$$u_{l_{1}}(x,t) = a_{0} + 6\gamma \left(\frac{2\beta}{\sqrt{\Theta}\tanh\left(0.5\sqrt{\Theta}(x-Vt)\right) + \gamma}\right) - 6\left(\frac{2\beta}{\sqrt{\Theta}\tanh\left(0.5\sqrt{\Theta}(x-Vt)\right) + \gamma}\right)^{2}, \Theta > 0,$$
(11)

$$u_{1_{2}}(x,t) = a_{0} - 6\gamma \left(\frac{2\beta}{\sqrt{-\Theta}\tan\left(0.5\sqrt{-\Theta}(x-Vt)\right) + \gamma}\right) - 6\left(\frac{2\beta}{\sqrt{-\Theta}\tan\left(0.5\sqrt{-\Theta}(x-Vt)\right) + \gamma}\right)^{2}, \Theta > 0,$$
(12)

$$u_{l_3}(x,t) = a_0 - 6\gamma \left(\frac{\gamma}{\exp(\gamma(x-Vt)) - 1}\right) - 6\left(\frac{\gamma}{\exp(\gamma(x-Vt)) - 1}\right)^2, \beta = 0, \Theta > 0$$
(13)

$$u_{l_{4}}(x,t) = a_{0} + 6\gamma \left(\frac{\gamma^{2}(x-Vt)}{2\gamma(x-Vt)+4}\right) - 6\left(\frac{\gamma^{2}(x-Vt)}{2\gamma(x-Vt)+4}\right)^{2}, \gamma \neq 0, \beta \neq 0, \Theta = 0,$$
(14)

$$u_{1_5}(x,t) = a_0 - 6\left(\frac{1}{x - Vt}\right)^2, \gamma = 0, \beta = 0.$$
(15)

where  $V = 22\beta\gamma^2 + 5\gamma^2a_0 + 76\beta^2 + \gamma^4 + 5a_0^2 + 40\beta a_0$  and  $a_0$  are arbitrary constants. For **Set 2**:

$$u_{2_{1}}\left(\xi\right) = a_{0} + 12\gamma \left(\frac{2\beta}{\sqrt{\Theta}\tanh\left(0.5\sqrt{\Theta}\left(x - Vt\right)\right) + \gamma}\right) - 12 \left(\frac{2\beta}{\sqrt{\Theta}\tanh\left(0.5\sqrt{\Theta}\left(x - Vt\right)\right) + \gamma}\right)^{2}, \Theta > 0, \tag{16}$$

$$u_{2_{2}}\left(\xi\right) = a_{0} - 12\gamma \left(\frac{2\beta}{\sqrt{-\Theta}\tan\left(0.5\sqrt{-\Theta}\left(x-Vt\right)\right)+\gamma}\right) - 12\left(\frac{2\beta}{\sqrt{-\Theta}\tan\left(0.5\sqrt{-\Theta}\left(x-Vt\right)\right)+\gamma}\right)^{2}, \Theta > 0, \quad (17)$$

$$u_{2_3}(x,t) = a_0 - 12\gamma \left(\frac{\gamma}{\exp(\gamma(x-Vt)) - 1}\right) - 12 \left(\frac{\gamma}{\exp(\gamma(x-Vt)) - 1}\right)^2, \beta = 0, \Theta > 0$$
(18)

$$u_{2_{4}}(x,t) = a_{0} + 12\gamma \left(\frac{\gamma^{2}(x-Vt)}{2\gamma(x-Vt)+4}\right) - 12 \left(\frac{\gamma^{2}(x-Vt)}{2\gamma(x-Vt)+4}\right)^{2}, \gamma \neq 0, \beta \neq 0, \Theta = 0,$$
(19)

$$u_{2_{5}}(x,t) = -12\left(\frac{1}{x-Vt}\right)^{2}, \gamma = 0, \beta = 0.$$
 (20)

where  $V = -8\beta \gamma^2 + \gamma^4 + 16\beta^2$  and  $a_0 = -\gamma^2 - 8\alpha\beta$  are arbitrary constants.

When  $\gamma = 0$ , the another type of explicit solutions to the fifth order standard SK Equation (6) are obtained in the followings form:

$$u_{3_1}(x,t) = a_0 + 6\alpha\beta \tanh^2 \left[ \sqrt{-\alpha\beta} \left\{ x - \left( 5a_0^2 + 76\alpha^2\beta^2 + 40a_0\alpha\beta \right) t \right\} \right], \alpha\beta < 0$$
<sup>(21)</sup>

$$u_{3_2}(x,t) = a_0 - 6\alpha\beta \tan^2 \left[ \sqrt{\alpha\beta} \left\{ x - \left( 5a_0^2 + 76\alpha^2\beta^2 + 40a_0\alpha\beta \right) t \right\} \right], \alpha\beta > 0$$
<sup>(22)</sup>

$$u_{3_3}(x,t) = a_0 + 6\alpha\beta \operatorname{coth}^2 \left[ \sqrt{-\alpha\beta} \left\{ x - \left( 5a_0^2 + 76\alpha^2\beta^2 + 40a_0\alpha\beta \right) t \right\} \right], \alpha\beta < 0$$
<sup>(23)</sup>

$$u_{3_{4}}(x,t) = a_{0} - 6\alpha\beta \cot^{2} \left[ \sqrt{\alpha\beta} \left\{ x - \left( 5a_{0}^{2} + 76\alpha^{2}\beta^{2} + 40a_{0}\alpha\beta \right) t \right\} \right], \alpha\beta > 0$$
(24)

$$u_{3_{5}}(x,t) = a_{0} - 6 \left( \frac{1}{x - \left(5a_{0}^{2} + 76\alpha^{2}\beta^{2} + 40a_{0}\alpha\beta\right)t} \right)^{2}, \alpha = 0, \beta > 0$$
<sup>(25)</sup>

$$u_{4_{1}}(x,t) = -8\alpha\beta + 12\alpha\beta \tanh^{2} \left[ \sqrt{-\alpha\beta} \left( x - 16\alpha^{2}\beta^{2}t \right) \right], \alpha\beta < 0$$
(26)

$$u_{4_2}(x,t) = -8\alpha\beta - 12\alpha\beta\tan^2\left[\sqrt{\alpha\beta}\left(x - 16\alpha^2\beta^2t\right)\right], \alpha\beta > 0$$
<sup>(27)</sup>

$$u_{4_3}(x,t) = -8\alpha\beta + 12\alpha\beta^2 \operatorname{coth}^2 \left[\sqrt{-\alpha\beta} \left(x - 16\alpha^2\beta^2 t\right)\right], \alpha\beta < 0$$
(28)

$$u_{4_4}(x,t) = -8\alpha\beta - 12\alpha\beta\cot^2\left[\sqrt{\alpha\beta}\left(x - 16\alpha^2\beta^2t\right)\right], \alpha\beta > 0$$
<sup>(29)</sup>

From the solutions (11)-(29), we observe that the proposed advance exponential expansion method scheme relating to the subsidiary equation gives more valuable explicit form solutions to the standard SK equation.

### 4. Comparison between Matinfar et al. [31] Solutions and Our Obtained Solutions

Matinfar et al. [31] have investigated the solutions of the fifth order standard Sawada-Kotera equation using the exp-function method to the following forms

$$u_1(x,t) = \frac{\frac{1}{4}e^{(-x+1t)} + 4 + e^{(x-1t)}}{\frac{1}{4}e^{(-x+1t)} + 1 + e^{(x-1t)}}, \quad u_2(x,t) = \frac{-\frac{1}{4}e^{(-x+t)} + 5 - e^{(x-t)}}{\frac{1}{4}e^{(-x+t)} + 1 + e^{(x-t)}}.$$

On the other hand, twenty solutions including Matinfar *et al.* [31] solutions are constructed of this equation by applying the generalized expansion method. It is noteworthy to observe that the obtained solutions are new, realistic and applicable to practical problems. The auxiliary equation used in this paper is different, so obtained solutions is also different. It can be shown that this method is easier than other methods and applicable to other nonlinear evolution equations.

## **5. Physical Explanations**

This section represents the physical explanations of the above determined solutions of the standard SK equations. The fifth order standard SK equation is an important model equation that appeared for describing several types of wave propagation such gravity-capillary waves, the propagation of shallow water waves over a flat surface and magneto-sound wave propagation in plasmas. It contains nonlinear terms and dispersive terms, so the solitary wave solutions of this equation have investigated. The solitary wave is a wave which propagates without any temporal evolution in shape or size when viewed in the reference frame moving with the group velocity of the wave. From the above determined solutions, we observe that the advance  $\exp(-\Phi(\xi))$ -expansion method are given not only more new multiple explicit solutions but also many types of exact traveling wave solutions as well as non-traveling wave solutions. That is, the solution of the Equation (7) according to the auxiliary nonlinear ordinary differential equations gives various types of traveling wave solutions, such as the solitary wave solution of bell type, the solitary wave solution of cuspon type, the solitary wave solutions by taking the particular values of the additional free parameters. The obtaining new exact traveling wave solutions of the standard SK equation are shown in Figures 1-7.

**Figure 1** displays the exact conidal soliton like wave profile according to the solution  $u_{l_1}(x,t)$  with the fixed values of  $a_0 = 0$ ,  $\gamma = 1$ ,  $\beta = 0.01$  and  $-6 \le x, t \le 6$ . **Figure 2** presents the exact soliton type wave profile corresponding to  $u_{3_3}(x,t)$  with fixed parameters  $a_0 = 1$ ,  $\alpha = -0.5$ ,  $\beta = 0.5$ , and  $-5 \le x, t \le 5$ . **Figure 3** shows the exact cuspon type (that is the peak amplitude of soliton is narrow) wave profile corresponding to  $u_{3_1}(x,t)$  with fixed parameters  $a_0 = 5$ ,  $\alpha = -1$ ,  $\beta = 1$  and  $-5 \le x, t \le 5$ . **Figure 4** exhibits the exact compacton type wave profile corresponding to the same solution  $u_{3_1}(x,t)$  with  $a_0 = 0.5$ ,  $\alpha = -0.01$ ,  $\beta = 0.05$  and  $-5 \le x, t \le 5$ . **Figure 5** represents the exact solutions corresponding to the rational function solution  $u_{1_4}(x,t)$  for the fixed values of  $a_0 = 0$ ,  $\gamma = 2$ ,  $\beta = 1$  and  $-6 \le x, t \le 6$ . **Figure 6** shows the exact non-topological bell nature solitary wave profile corresponding to  $u_{1_4}(x,t)$  for  $\alpha = -0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . **Figure 7** displays the exact periodic traveling wave profile corresponding to  $u_{4_2}(x,t)$  with fixed parameters  $\alpha = 0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . Figure 8 represents the exact solitary wave profile of singular soliton type are profile to  $u_{4_3}(x,t)$  with fixed parameters  $\alpha = -0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . Figure 8 represents the exact solitary wave profile of singular soliton type corresponding to  $u_{4_3}(x,t)$  with fixed parameters  $\alpha = -0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . Figure 8 represents the exact solitary wave profile of solitor solitary wave profile corresponding to  $\alpha = -0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . Figure 8 represents the exact solitary wave profile of singular soliton type corresponding to  $u_{4_3}(x,t)$  with fixed parameters of  $\alpha = -0.5$ ,  $\beta = 0.5$  and  $-5 \le x, t \le 5$ . Figure 8 represents the exact solitary wave profile of singular soliton type corresponding to  $u_{4_3}(x,t)$  with fixed parameteric values of  $\alpha = -0.5$ ,  $\beta = 0.$ 



Figure 1. Surface plot of exact soliton like solitary wave solution.



Figure 2. Surface plot of exact multiple soliton like solitary wave solution.



Figure 3. Surface plot of exact cuspon like solitary wave solution.



Figure 4. Surface plot of exact compacton like solitary wave solutions.



Figure 5. Surface plot of soliton type solitary wave solution.



Figure 6. Surface plot of exact bell like solitary wave solution.



Figure 8. Surface plot of exact singular soliton type solitary wave solution.

 $-5 \le x, t \le 5$ . Thus we should point out the theoretical as well as numerical investigation that are presented in this manuscript are useful in investigating the nonlinear traveling wave propagation in any natural varied instance according to the variation of the parameters.

## 6. Conclusion

In this article, the generalized  $\exp(-\Phi(\xi))$ -expansion method has been proposed and successfully applied to find the new exact traveling wave solutions to the fifth order standard SK equation. This method not only provides some pre-published solutions extracted by other methods but can also pick up what we accept as true to be new solutions missed by other research scholar. The obtained solutions in this article provide that the proposed method is powerful, and straightforward and easily applicable for solving nonlinear evaluation equations arise in mathematical physics. The paper also shows that the proposed method is very standard mathematical tool to find the traveling wave solutions of nonlinear evaluation equations rather than other existing method, particularly the exponential expansion method and the  $\exp(-\Phi(\xi))$ -expansion method. We have applied this method to the fifth order standard SK equation, but it can be further applied to other evolution equations to construct new exact solutions.

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