# Hybrid Post-Processing Procedure for Displacement-Based Plane Elements 

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Received August 2013


#### Abstract

In the analysis of high-rise building, traditional displacement-based plane elements are often used to get the in-plane internal forces of the shear walls by stress integration. Limited by the singular problem produced by wall holes and the loss of precision induced by using differential method to derive strains, the displacement-based elements cannot always present accuracy enough for design. In this paper, the hybrid post-processing procedure based on the Hellinger-Reissner variational principle is used for improving the stress precision of two quadrilateral plane elements. In order to find the best stress field, three different forms are assumed for the displacement-based plane elements AQ4 $\theta$ and AQ4 $\theta \lambda$ with drilling DOF. Numerical results show that by using the proposed method, the accuracy of stress solutions of these two displacement-based plane elements can be improved.


Keywords: Finite Element; Displacement-Based Plane Element; Hybrid Post-Processing Procedure

## 1. Introduction

For traditional displacement-based elements derived from the principle of minimum potential energy, the trial functions are usually assumed as a polynomial of nodal displacements, and then the internal forces or stresses can be derived by stress-strain relationship. By this process of formulation, these kinds of elements can present accurate displacement results, otherwise the construction procedures of these formulations are also easier. But in most situations, the stress results are much more important than the displacements such as the designing of high-rise buildings. The reinforcements of shear walls are depended on the internal forces integrated from stresses. For the sake of differential with displacement, the accuracy of stresses is usually lower than the displacement results. Many efforts have been made for overcoming this shortage, such as the assumed strain formulations developed by MacNeal [1], Piltner and Taylor [2].
On the other hand, the hybrid/mixed elements are developed from the modified variational principle of HuWashisu, such as the hybrid method proposed by Pian [3], Pian and Sumihara [4]. It is different from the displace-ment-based elements. To formulate these elements, the internal force field or stress field should be assumed at first. Without the differential process which is needed for
the displacement-based elements to derive strains, so they can exhibit better stress results. But they also present obvious disadvantages. For example, the construction procedures are much more complex, and the calculation of matrix inversion and condense are usually unavoidable, furthermore, the accuracy of displacement results usually are not as accurate as the displacement-based elements.

By combining the merits of both hybrid elements and displacement-based elements, a new method named hybrid post procedure was proposed by Cen [5]. In this method, the nodal displacements are derived from dis-placement-based plate elements at first, after this, a new internal force field which should satisfy the equilibrium Equation is assumed to substitute into the hybrid energy functional, by the principle of stationary, the undetermined parameters of the assumed internal force field can be solved out, and finally the new internal forces can be calculated easily. It is seemed very effective for improving the internal force accuracy of plate element. In this paper, this method is extended to improve the plane elements.

## 2. Hybrid Post-Processing Procedure

For the hybrid plate elements, the discrete energy functional can be written as follows [5,6]:

$$
\begin{align*}
\Pi_{R}^{e}= & -\frac{1}{2} \iint_{A^{e}}\{\mathbf{M}\}^{T}\left[\mathbf{D}_{b}\right]^{-1}\{\mathbf{M}\} d A \\
& -\frac{1}{2} \iint_{A^{e}}\{\mathbf{T}\}^{T}\left[\mathbf{D}_{s}\right]^{-1}\{\mathbf{T}\} d A  \tag{1}\\
& +\iint_{A^{e}}\{\mathbf{M}\}^{T}\{\mathbf{\kappa}\} d A+\iint_{A^{e}}\{\mathbf{T}\}^{T}\{\boldsymbol{\gamma}\} d A-\iint_{A^{e}} f_{z} w d A \\
& -\int_{s}\left(\bar{T}_{n} w+\bar{M}_{x n} \psi_{x}+\bar{M}_{y n} \psi_{y}\right) d s
\end{align*}
$$

where $W_{\text {exp }}$ is the work of external forces:

$$
\begin{equation*}
W_{\exp }=\iint_{A^{e}} f_{z} w d A+\int_{s}\left(\bar{T}_{n} w+\bar{M}_{x n} \psi_{x}+\bar{M}_{y n} \psi_{y}\right) d s \tag{2}
\end{equation*}
$$

Assume the new internal force field as follows:

$$
\begin{align*}
& \{\mathbf{M}\}=\left[\mathbf{P}_{M}\right]\left\{\boldsymbol{\alpha}_{M}\right\} \\
& \{\mathbf{T}\}=\left\{\begin{array}{l}
T_{x} \\
T_{y}
\end{array}\right\}=\left\{\begin{array}{l}
M_{x, x}+M_{x y, y} \\
M_{x y, x}+M_{y, y}
\end{array}\right\}=\left[\mathbf{P}_{T}\right]\left\{\boldsymbol{\alpha}_{M}\right\} \tag{3}
\end{align*}
$$

where

$$
\begin{align*}
& \{\mathbf{M}\}=\left[\begin{array}{lll}
M_{x} & M_{y} & M_{x y}
\end{array}\right]^{T} \\
& {\left[\mathbf{P}_{M}\right]=\left[\begin{array}{cccccccccccc}
1 & \xi & \eta & \xi \eta & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 1 & \xi & \eta & \xi \eta & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & \xi & \eta & \xi \eta
\end{array}\right]} \\
& {\left[\mathbf{P}_{T}\right]=\left[\begin{array}{cccccccc}
0 & j_{11} & j_{12} & j_{11} \eta+j_{12} \xi & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & j_{21} & j_{22} & j_{21} \eta+j_{22} \xi
\end{array}\right.}  \tag{4}\\
& \left.\begin{array}{llll}
0 & j_{21} & j_{22} & j_{21} \eta+j_{22} \xi
\end{array}\right] \\
& \left.\begin{array}{llll}
0 & j_{11} & j_{12} & j_{11} \eta+j_{12} \xi
\end{array}\right] \\
& \left\{\boldsymbol{\alpha}_{M}\right\}=\left[\begin{array}{lllll}
\alpha_{1} & \alpha_{2} & \cdots & \alpha_{11} & \alpha_{12}
\end{array}\right]^{T}
\end{align*}
$$

and $j_{11}, j_{12}, j_{21}, j_{22}$ are components of the inversion matrix of Jacobian .
After substituting Equation (3) into Equation (1), the new form of energy functional can be written as:

$$
\begin{align*}
\Pi_{R}^{e} & =-\frac{1}{2}\left\{\boldsymbol{\alpha}_{M}\right\}^{T}\binom{\iint_{A^{e}}\left[\mathbf{P}_{M}\right]^{T}\left[\mathbf{D}_{b}\right]^{-1}\left[\mathbf{P}_{M}\right] d A}{+\iint_{A^{e}}\left[\mathbf{P}_{T}\right]^{T}\left[\mathbf{D}_{s}\right]^{-1}\left[\mathbf{P}_{T}\right] d A}\left\{\boldsymbol{\alpha}_{M}\right\} \\
& +\left\{\boldsymbol{\alpha}_{M}\right\}^{T} \iint_{A^{e}}\left(\left[\mathbf{P}_{M}\right]^{T}\left[\mathbf{B}_{b}\right]\right.  \tag{5}\\
& \left.+\left[\mathbf{P}_{T}\right]^{T}\left[\mathbf{B}_{s}\right]\right) d A \cdot\{\mathbf{q}\}^{e}-W_{\exp }
\end{align*}
$$

By the principle of stationary:

$$
\begin{equation*}
\frac{\partial \Pi_{R}^{e}}{\partial\left\{\boldsymbol{\alpha}_{M}\right\}}=0 \tag{6}
\end{equation*}
$$

the parameters of the assumed internal force field can be written as:

$$
\begin{equation*}
\left\{\boldsymbol{\alpha}_{M}\right\}=-\left[\mathbf{K}_{M M}\right]^{-1}\left[\mathbf{K}_{M q}\right]\{\mathbf{q}\}^{e} \tag{7}
\end{equation*}
$$

where

$$
\begin{align*}
{\left[\mathbf{K}_{M M}\right]=} & -\iint_{A^{e}}\left(\left[\mathbf{P}_{M}\right]^{T}\left[\mathbf{D}_{b}\right]^{-1}\left[\mathbf{P}_{M}\right]\right. \\
& \left.+\left[\mathbf{P}_{T}\right]^{T}\left[\mathbf{D}_{s}\right]^{-1}\left[\mathbf{P}_{T}\right]\right) d A  \tag{8}\\
{\left[\mathbf{K}_{M q}\right]=} & \iint_{A^{e}}\left(\left[\mathbf{P}_{M}\right]^{T}\left[\mathbf{B}_{b}\right]+\left[\mathbf{P}_{T}\right]^{T}\left[\mathbf{B}_{s}\right]\right) d A
\end{align*}
$$

Substitute Equation (7) into Equation (3), the new internal force can be obtained.

Compared with plate elements, it is much easier to extend this method to plane elements. The according hybrid discrete energy functional of plane elements can be written as [6]:

$$
\begin{equation*}
\Pi_{R}^{e}=\iint_{V^{e}}\left[\{\boldsymbol{\sigma}\}^{T} D\{\mathbf{u}\}-\frac{1}{2}\{\boldsymbol{\sigma}\}^{T}[\mathbf{D}]^{-1}\{\boldsymbol{\sigma}\}\right] d \mathrm{~V} \tag{9}
\end{equation*}
$$

Similar to Equation (4), a new form of stress field of displacement-based plane element can be assumed as:

$$
\begin{equation*}
\{\boldsymbol{\sigma}\}=\left[\mathbf{P}_{\sigma}\right]\left\{\boldsymbol{\alpha}_{\sigma}\right\} \tag{10}
\end{equation*}
$$

Substitute Equation (10) into Equation (9), the parameters $\alpha_{\mathrm{i}}$ in $\left[\mathbf{P}_{\sigma}\right]$ can be obtained by using the principle of stationary. The according matrix named $\left[\mathbf{K}_{\sigma \sigma}\right]$ and $\left[\mathbf{K}_{\sigma q}\right]$ for plane elements are as follows:

$$
\begin{align*}
& {\left[\mathbf{K}_{\sigma \sigma}\right]=-\iint_{A^{e}}\left[\mathbf{P}_{\sigma}\right]^{T}[\mathbf{D}]^{-1}\left[\mathbf{P}_{\sigma}\right] t d A} \\
& {\left[\mathbf{K}_{\sigma q}\right]=\iint_{A^{e}}\left[\mathbf{P}_{\sigma}\right]^{T}\left[\mathbf{B}_{b}\right] t d A} \tag{11}
\end{align*}
$$

With different matrix of $\left[\mathbf{P}_{\sigma}\right]$, the stress field will be different too. This new method will be used to try to improve the stress accuracy of plane element named $\mathrm{AQ} 4 \theta$ and $\mathrm{AQ} 4 \theta \lambda$ [7].

## 3. Introduction of $A Q 4 \theta$ and $A Q 4 \theta \lambda$

$\mathrm{AQ} 4 \theta$ and $\mathrm{AQ} 4 \theta \lambda$ are two plane elements with drilling DOF formulated using quadrilateral area coordinate methods presented by Long et al. [8,9]. They have the merits of high accuracy and robust against mesh distortions. After having been programmed into the software for the analysis of high-rise buildings, the results show that better accuracy is still needed for the stress of shear walls with holes.

The definition of DOF of these two elements is:

$$
\{\mathbf{q}\}^{e}=\left[\begin{array}{llll}
\left\{\mathbf{q}_{1}\right\}^{T} & \left\{\mathbf{q}_{2}\right\}^{T} & \left\{\mathbf{q}_{3}\right\}^{T} & \left\{\mathbf{q}_{4}\right\}^{T} \tag{12}
\end{array}\right]^{T}
$$

where

$$
\left\{\mathbf{q}_{i}\right\}=\left[\begin{array}{lll}
u_{i} & v_{i} & \theta_{i} \tag{13}
\end{array}\right]^{T}(i=1,2,3,4)
$$

and $\theta_{i}$ is the additional rigid rotation at element node.
The displacements of element are as follows:

$$
\begin{equation*}
\{\mathbf{u}\}=\left\{\mathbf{u}^{0}\right\}+\left\{\mathbf{u}^{\theta}\right\} \tag{14}
\end{equation*}
$$

where $\left\{\mathbf{u}^{0}\right\}$ is a polynomial about $u_{i}$ and $v_{i},\left\{\mathbf{u}^{\theta}\right\}$ is
the additional displacement field induced only by the rigid rotation at nodes denoted as $\theta_{i}$.

In order to determine the displacement field $\left\{\mathbf{u}^{0}\right\}$, the shape functions in reference [10] is used, they are as follows:

$$
\begin{align*}
& u^{0}=\sum_{i=1}^{4} N_{i}^{0} u_{i}  \tag{15}\\
& v^{0}=\sum_{i=1}^{4} N_{i}^{0} v_{i}
\end{align*}
$$

where:

$$
\begin{align*}
& N_{i}^{0}=-\frac{1}{2} g_{i+2}+L_{i}+L_{i+1}+\xi_{i} \eta_{i} g_{i+2} P  \tag{16}\\
& (i=1,2,3,4)
\end{align*}
$$

and

$$
\begin{aligned}
& P= \\
& \frac{3\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right)-\left(g_{2}-g_{3}\right)\left(L_{3}-L_{1}\right)-\left(g_{1}-g_{2}\right)\left(L_{4}-L_{2}\right)-\frac{1}{2}\left(g_{2} g_{4}-g_{1} g_{3}\right)(17)}{1+g_{1} g_{3}+g_{2} g_{4}}
\end{aligned}
$$

Assume the rotational displacement field as polynomials of quadrilateral area coordinates:

$$
\begin{align*}
u_{\theta}= & \alpha_{1}+\alpha_{2}\left(L_{3}-L_{1}\right)+\alpha_{3}\left(L_{4}-L_{2}\right) \\
& +\alpha_{4}\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right) \\
& +\alpha_{5} L_{1} L_{3}+\alpha_{6} L_{2} L_{4} \\
v_{\theta}= & \beta_{1}+\beta_{2}\left(L_{3}-L_{1}\right)+\beta_{3}\left(L_{4}-L_{2}\right)  \tag{18}\\
& +\beta_{4}\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right) \\
& +\beta_{5} L_{1} L_{3}+\beta_{6} L_{2} L_{4}
\end{align*}
$$

with the conforming Equations as follows, the rotational displacement can be obtained.

$$
\begin{align*}
& \sum_{i=1}^{4}\left(u_{\theta}^{i}-\bar{u}_{\theta}^{i}\right)=0 \\
& \sum_{i=1}^{4} \xi_{i} \eta_{i}\left(u_{\theta}^{i}-\bar{u}_{\theta}^{i}\right)=0  \tag{19}\\
& \int_{l_{y}}\left(u_{\theta}-\bar{u}_{\theta}\right) d \bar{s}=0 \quad(i j=12,23,34,41)
\end{align*}
$$

where

$$
\begin{align*}
& \left\{\begin{array}{l}
\bar{u}_{\theta 12} \\
\bar{v}_{\theta 12}
\end{array}\right\}=\left\{\begin{array}{l}
-b_{4} \\
-c_{4}
\end{array}\right\} L_{1} L_{3}\left(\frac{1}{g_{1} g_{2}^{2}} L_{1} \theta_{1}-\frac{1}{g_{1}^{2} g_{2}} L_{3} \theta_{2}\right) \\
& \left\{\begin{array}{l}
\bar{u}_{\theta 23} \\
\bar{v}_{\theta 23}
\end{array}\right\}=\left\{\begin{array}{l}
-b_{1} \\
-c_{1}
\end{array}\right\} L_{2} L_{4}\left(\frac{1}{g_{2} g_{3}^{2}} L_{2} \theta_{2}-\frac{1}{g_{2}^{2} g_{3}} L_{4} \theta_{3}\right)  \tag{20}\\
& \left\{\begin{array}{l}
\bar{u}_{\theta 34} \\
\bar{v}_{\theta 34}
\end{array}\right\}=\left\{\begin{array}{l}
-b_{2} \\
-c_{2}
\end{array}\right\} L_{1} L_{3}\left(\frac{1}{g_{3} g_{4}^{2}} L_{3} \theta_{3}-\frac{1}{g_{3}^{2} g_{4}} L_{1} \theta_{4}\right) \\
& \left\{\begin{array}{l}
\bar{u}_{\theta 41} \\
\bar{v}_{\theta 41}
\end{array}\right\}=\left\{\begin{array}{l}
-b_{3} \\
-c_{3}
\end{array}\right\} L_{2} L_{4}\left(\frac{1}{g_{4} g_{1}^{2}} L_{4} \theta_{4}-\frac{1}{g_{4}^{2} g_{1}} L_{2} \theta_{1}\right)
\end{align*}
$$

and $b_{i}=y_{i+1}-y_{i+2}, \quad c_{i}=x_{i+2}-x_{i+1}$

Then the stiffness matrix of element $\mathrm{AQ} 4 \theta$ can be solved out easily, and the strain matrix is:

$$
\begin{equation*}
\{\varepsilon\}=\left[\boldsymbol{B}_{q}\right]\{\boldsymbol{q}\}^{e} \tag{21}
\end{equation*}
$$

where

$$
\left[\boldsymbol{B}_{q}\right]=\left[\begin{array}{llll}
\boldsymbol{B}_{1} & \boldsymbol{B}_{2} & \boldsymbol{B}_{3} & \boldsymbol{B}_{4} \tag{22}
\end{array}\right]
$$

and

$$
\left[\boldsymbol{B}_{i}\right]=\left[\begin{array}{ccc}
\frac{\partial N_{i}^{0}}{\partial x} & 0 & \frac{\partial N_{u \theta i}}{\partial x}  \tag{23}\\
0 & \frac{\partial N_{i}^{0}}{\partial y} & \frac{\partial N_{v \theta_{i}}}{\partial y} \\
\frac{\partial N_{i}^{0}}{\partial y} & \frac{\partial N_{i}^{0}}{\partial x} & \frac{\partial N_{u \theta i}}{\partial y}+\frac{\partial N_{v \theta_{i}}}{\partial x}
\end{array}\right]
$$

After substituting Equation (21) into the stress-strain relationship, the stress of $\mathrm{AQ} 4 \theta$ can be obtained.
$\mathrm{AQ} 4 \theta$ is an improved element based on $\mathrm{AQ} 4 \theta$ by adding a displacement field which is induced by internal parameters.

The additional displacement fields mentioned above are as follows:

$$
\begin{align*}
u_{\lambda}= & \alpha_{1}+\alpha_{2}\left(L_{3}-L_{1}\right)+\alpha_{3}\left(L_{4}-L_{2}\right) \\
& +\alpha_{4} L_{1} L_{3}+\lambda_{1} L_{2} L_{4}+\lambda_{2}\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right) \\
v_{\lambda} & =\beta_{1}+\beta_{2}\left(L_{3}-L_{1}\right)+\beta_{3}\left(L_{4}-L_{2}\right)  \tag{24}\\
& +\beta_{4} L_{1} L_{3}+\lambda_{1}^{\prime} L_{2} L_{4}+\lambda_{2}^{\prime}\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right)
\end{align*}
$$

where $\lambda_{1}, \lambda_{1}^{\prime}, \lambda_{2}, \lambda_{2}^{\prime}$ are the internal parameters.
In order to solve the undetermined parameters, the conforming Equations are taken as:

$$
\begin{equation*}
\int_{l_{i j}}\left\{\mathbf{u}_{\lambda}\right\} d s=0 \tag{25}
\end{equation*}
$$

Substitute Equation (24) into (25), the shape functions of $\left\{\mathbf{u}_{\lambda}\right\}$ can be formulated out, they are:

$$
\begin{align*}
N_{\lambda 1}= & \frac{\left(6 g_{1} g_{2} g_{3} g_{4}-g_{1} g_{3}-g_{2} g_{4}\right)}{6\left(1-g_{1} g_{4}-g_{2} g_{3}\right)} \\
& +\frac{g_{1}-g_{2}}{6}\left(L_{3}-L_{1}\right) \\
& +\frac{\left(g_{1}-g_{4}\right)\left(g_{1} g_{4}+g_{2} g_{3}\right)}{6\left(1-g_{1} g_{4}-g_{2} g_{3}\right)}\left(L_{4}-L_{2}\right)  \tag{26}\\
& +\frac{g_{1} g_{4}+g_{2} g_{3}}{1-g_{1} g_{4}-g_{2} g_{3}} L_{1} L_{3}+L_{2} L_{4} \\
N_{\lambda 2}= & \frac{1}{3}\left(g_{3}-g_{2}\right)\left(L_{3}-L_{1}\right)-2\left(g_{1}-g_{2}\right)\left(L_{4}-L_{2}\right) \\
& +\left(L_{3}-L_{1}\right)\left(L_{4}-L_{2}\right)
\end{align*}
$$

For element $\mathrm{AQ} 4 \theta \lambda, N_{\lambda 1}$ is taken as the final internal shape function, then:

$$
\begin{equation*}
\left\{\mathbf{u}_{\lambda}\right\}=\left[\mathbf{N}_{\lambda}\right]\{\lambda\} \tag{27}
\end{equation*}
$$

where

$$
\begin{align*}
& \{\boldsymbol{\lambda}\}=\left[\begin{array}{ll}
\lambda_{1} & \lambda_{1}^{\prime}
\end{array}\right]^{T} \\
& {\left[\mathbf{N}_{\lambda}\right]=\left[\begin{array}{cc}
N_{\lambda 1} & 0 \\
0 & N_{\lambda 1}
\end{array}\right]} \tag{28}
\end{align*}
$$

Through the condense calculation, stiffness matrix of $\mathrm{AQ} 4 \theta \lambda$ can be written as:

$$
\begin{equation*}
[\boldsymbol{k}]^{e}=\left[\boldsymbol{k}_{q q}\right]-\left[\boldsymbol{k}_{\lambda q}\right]^{T}\left[\boldsymbol{k}_{\lambda \lambda}\right]^{-1}\left[\boldsymbol{k}_{\lambda q}\right] \tag{29}
\end{equation*}
$$

where

$$
\begin{align*}
& {\left[\mathbf{k}_{q q}\right]=\iint\left[\mathbf{B}_{q}\right]^{T}[\mathbf{D}]\left[\mathbf{B}_{q}\right] t d A} \\
& {\left[\mathbf{k}_{\lambda \lambda}\right]=\iint\left[\mathbf{B}_{\lambda}\right]^{T}[\mathbf{D}]\left[\mathbf{B}_{\lambda}\right] t d A}  \tag{30}\\
& {\left[\mathbf{k}_{\lambda q}\right]=\iint\left[\mathbf{B}_{\lambda}\right]^{T}[\mathbf{D}]\left[\mathbf{B}_{q}\right] t d A}
\end{align*}
$$

and $\left[\mathbf{B}_{\lambda}\right]$ is the strain matrix of additional displacement field. The according stress fields are as follows:

$$
\begin{equation*}
\left\{\boldsymbol{\varepsilon}_{\lambda}\right\}=\left[\mathbf{B}_{\lambda}\right]\{\lambda\} \tag{31}
\end{equation*}
$$

Combined Equation (21) with (31), the stress field of $\mathrm{AQ} 4 \theta \lambda$ can be obtained.

## 4. Hybrid Post-Processing of AQ4 $\theta$ and AQ4日

In the formulating of stress in elements $\mathrm{AQ} 4 \theta$ and $\mathrm{AQ} 4 \theta \lambda$, the strain matrix $\left[\mathbf{B}_{\mathrm{q}}\right]$ and $\left[\mathbf{B}_{\lambda}\right]$ are the differential results of displacements to coordinates. The differential calculation lowered the accurate order of stress. To avoid the differential, Hybrid Post-processing procedures can be used to improve the stress accuracy. Based on this theory, three forms of stress fields are presented for these two displacement-based elements.

### 4.1. Stress Field I

The first form of stress field is assumed as Equation (32). It is the stress field of a hybrid element developed by Pian and Wu [6]:

$$
\{\boldsymbol{\sigma}\}=\left[\begin{array}{ccccc}
1 & & & a_{1}^{2} \eta & a_{3}^{2} \xi  \tag{32}\\
& 1 & & b_{1}^{2} \eta & b_{3}^{2} \xi \\
& & 1 & a_{1} b_{1} \eta & a_{3} b_{3} \zeta
\end{array}\right]\left\{\begin{array}{c}
\beta_{1} \\
\vdots \\
\beta_{5}
\end{array}\right\}
$$

where

$$
\begin{array}{lll}
a_{1}=\frac{1}{4} \sum_{i=1}^{4} \xi_{i} x_{i} & a_{2}=\frac{1}{4} \sum_{i=1}^{4} \xi_{i} \eta_{i} x_{i} & a_{3}=\frac{1}{4} \sum_{i=1}^{4} \eta_{i} x_{i}  \tag{33}\\
b_{1}=\frac{1}{4} \sum_{i=1}^{4} \xi_{i} y_{i} & b_{2}=\frac{1}{4} \sum_{i=1}^{4} \xi_{i} \eta_{i} y_{i} & b_{3}=\frac{1}{4} \sum_{i=1}^{4} \eta_{i} y_{i}
\end{array}
$$



Figure 1. Patch test.


Figure 2. Cook's skew beam.
Table 1. Stress at point A and B of Cook's Beam of AQ4 $\boldsymbol{\theta}$.

| Method | $\sigma_{\text {Amax }}$ |  |  | $\sigma_{\text {Bmin }}$ |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $2 \times 2$ | $4 \times 4$ | $8 \times 8$ | $2 \times 2$ | $4 \times 4$ | $8 \times 8$ |
| Field I | 0.1791 | 0.2261 | 0.2338 | -0.1700 | -0.1929 | -0.2002 |
| Field II | 0.1951 | 0.2298 | 0.2348 | -0.1942 | -0.1933 | -0.2010 |
| Field III | 0.1914 | 0.2240 | 0.2319 | -0.1769 | -0.1938 | -0.2009 |
| Source Val. | 0.1917 | 0.2241 | 0.2377 | -0.1877 | -0.1939 | -0.2060 |
| Ref. Val. |  | 0.2362 |  |  | -0.2023 |  |

Table 2. Stress at point A and B of Cook's Beam of AQ4 $\theta \lambda$.

| Method | $\sigma_{\text {Amax }}$ |  |  |  |  | $\sigma_{\text {Bmin }}$ |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $2 \times 2$ | $4 \times 4$ | $8 \times 8$ | $2 \times 2$ | $4 \times 4$ | $8 \times 8$ |  |  |
|  | 0.1913 | 0.2271 | 0.2342 | -0.1748 | -0.1919 | -0.2009 |  |  |
|  | 0.2145 | 0.2358 | 0.2364 | -0.2084 | -0.2032 | -0.2027 |  |  |
|  | 0.2147 | 0.2358 | 0.2364 | -0.2092 | -0.2033 | -0.2027 |  |  |
|  | 0.2498 | 0.2338 | 0.2358 | -0.1729 | -0.1896 | -0.2018 |  |  |
|  |  | 0.2362 |  |  | -0.2023 |  |  |  |

the Hellinger-Reissner variational principle, hybrid postprocess procedure can take advantage of the merits of
these two kinds of elements to establish the relationship between the displacement and the stress or internal force fields. In this paper, based on this theory, three forms of stress fields are used to improve the stress of plane elements with drilling DOF. Through the numerical results, for element AQ4日, only the second form of stress field is effective, but for $\mathrm{AQ} 4 \theta \lambda$, except for the first form of stress, the other two forms can present better results than the source elements. It is proved that the method of hybrid post-process procedure is workable.

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