Influence of Characteristics of Substance on Parameters of Interaction of Photons High Energy with Free Electrons

Andrey N. Volobuev, Eugene S. Petrov

Department of Physics, Samara State University, Samara, Russia E-mail: volobuev@samaramail.ru Received September 16, 2011; revised October 19, 2011; accepted November 12, 2011

Abstract

Various variants of interaction of photons high energy with free electrons in substance are investigated. It is shown, that among these variants, in substance can be observed: absorption of a photon by electron, coherent and not coherent scattering of photons, a stop electron after interaction with a photon. Dependence of change of length of a wave of a photon after interaction with electron from parameters of substance and speed of movement electron is found.

Keywords: Photons, Electrons, Interaction, Not Coherent Scattering, Compton Effect, Substance

1. Introduction

Interaction of photons and free electrons in substance it is problem known enough. One of results of such interaction, for example, is not coherent scattering. Compton effects or not coherent scattering of X-rays on electrons, poorly connected with a nucleus, at an irradiation of a paraffin surface, which was primarily observed by Compton, is well studied problem [1]. In a spectrum of scattered radiation he observed lower frequency, than frequency of falling radiation. Compton effects represent an example of the quantum nature of light since there is no classical analogue of this phenomenon which could be described Maxwell equations.

There are attempts to describe Compton effects on the basis of wave Doppler effects [2]. Thus electron it is primary considered as the mobile receiver of a wave, and then as a mobile source of a wave. However Compton criticized similar attempts [3]. We shall mark, that the scattering on motionless electron, that in strict mathematical sense it is Compton effects, in general it is not put down in the given model.

Rather detailed analysis of Compton effects is present in [4]. But the analysis [4] is directed on research of effective sections of interaction of electrons and photons. Besides the analysis of Compton effects is carried out in the assumption of is absence of influence of surrounding medium on process.

The purpose of the research is the analysis, first of all, not coherent shift of frequency of a photon and possible results, at interaction of a photon and free electron in substance, and also influence of characteristics of substance on the processes.

Let's use the method applied in [5] at research of interaction of photons of electromagnetic radiation with binding electrons.

2. Laws of Interaction of a Photon with Free Electron

Let's consider interaction moving electron with a photon, **Figure 1**. In all further transformations we believe, that the system of readout is connected to motionless nodes of a crystal lattice of substance.

The law of conservation of energy at interaction of a photon with moving electron we shall write down as:

$$E + h\nu = E' + h\nu', \qquad (1)$$

where E and E' is energy of electron before interaction with a photon, h is Planck's constant, v and v'-the frequency of a photon before and after interaction with electron.

Impulse of a photon before interaction with electron is

$$p_0 = \frac{nv}{V_{ph}} = \frac{nnv}{c}$$
, where V_{ph} is a speed of a photon in

substance, $n = \frac{c}{V_{ph}}$ —an absolute index of refraction of

substance, *c*- speed of light in vacuum.

Hence, the formula (1) will be copied as:



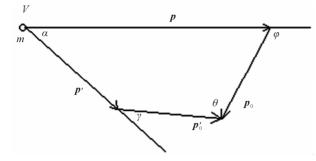


Figure 1. Direction of vectors of impulses of particles at interaction of a photon with free electron.

$$E + \frac{p_0 c}{n} = E' + \frac{p'_0 c}{n},$$
 (2)

where $p'_0 = \frac{nhv'}{c}$ is an impulse of a photon after inter-

action with electron.

The law of preservation of an impulse before interaction, **Figure 1**, looks like:

$$p + p_0 = p' + p'_0$$
, (3)

Let's transform (3) to the scalar form:

$$p_0^2 + p_0'^2 - 2p_0 p_0' \cos\theta = p^2 + p'^2 - 2pp' \cos\alpha \qquad (4)$$

where α is an angle between directions of an electron impulse before interaction with a photon, θ is an angle between impulses of falling and scattered photons.

Using connection between energy and speed of elec-

tron $E = mc^2 = \frac{p}{V}c^2$, where p = mV, sizes *m* and *V*-

mass and speed of electron before interaction, we shall transform the Equation (1) to a kind:

$$\frac{pc}{V} + \frac{p_0}{n} = \frac{p'c}{V'} + \frac{p'_0}{n} \,. \tag{5}$$

In the formula (5) $E' = \frac{p'}{V'}c^2$, and V' is electron's

speed after interaction with a photon.

Let's exclude from (4) and (5) value p'. Finding from (5) value p' and inserting it in (4), we shall receive:

$$p_{0}^{2} + p_{0}^{\prime 2} - 2p_{0}p_{0}^{\prime}\cos\theta = p^{2} + \left[\frac{V^{\prime}}{c}\left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}^{\prime}}{n}\right)\right]^{2} - 2p\frac{V^{\prime}}{c}\left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}^{\prime}}{n}\right)\cos\alpha$$
(6)

Let's exclude from (6) electron's speed V' after interaction with a photon. By definition of an impulse of a particle we have:

$$p' = \frac{m_0 V'}{\sqrt{1 - \frac{{V'}^2}{c^2}}},$$
(7)

where m_0 is a rest mass of electron. Inserting (7) into (5), we shall find:

$$\frac{m_0 c}{\sqrt{1 - \frac{{V'}^2}{c^2}}} = \frac{pc}{V} + \frac{p_0}{n} - \frac{p'_0}{n} \,. \tag{8}$$

Solving (8) be relative V', we have:

$$\frac{V'}{c} = \sqrt{1 - \left(\frac{m_0 c}{\frac{pc}{V} + \frac{p_0}{n} - \frac{p'_0}{n}}\right)^2} .$$
(9)

Inserting (9) into (6), we shall find:

$$p_{0}^{2} + p_{0}^{\prime 2} - 2p_{0}p_{0}^{\prime}\cos\theta$$

$$= p^{2} + \left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}^{\prime}}{n}\right)^{2} - (m_{0}c)^{2}$$

$$-2p\left(\sqrt{\left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}^{\prime}}{n}\right)^{2} - (m_{0}c)^{2}}\right)\cos\alpha$$
(10)

The law of preservation of an impulse (3) in the scalar form can be written down differently, than (4):

$$p^{2} + p_{0}^{2} - 2pp_{0}\cos(180^{\circ} - \varphi)$$

= $p'^{2} + p_{0}'^{2} - 2p'p_{0}'\cos(180^{\circ} - \gamma)'$, (11)

or:

$$p^{2} + p_{0}^{2} + 2pp_{0}\cos\varphi = p^{\prime 2} + p_{0}^{\prime 2} + 2p^{\prime}p_{0}^{\prime}\cos\gamma, \quad (12)$$

where φ is an angle between directions of movement electron and a photon before interaction, γ - is an angle between directions electron and a photon after interaction.

As well as in the previous case, inserting value p' from (5) into (12), we have:

$$p^{2} + p_{0}^{2} + 2pp_{0}\cos\varphi = \left(\frac{V'}{c}\right)^{2} \left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}'}{n}\right)^{2} + {p_{0}'}^{2} + 2p'_{0}\left(\frac{V'}{c}\right) \left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}'}{n}\right)\cos\gamma$$
(13)

Further, using (9), we shall copy (13) as:

$$p^{2} + p_{0}^{2} + 2pp_{0}\cos\varphi = \left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}'}{n}\right)^{2} - (m_{0}c)^{2} + p_{0}'^{2} + 2p_{0}'\left(\sqrt{\left(p\frac{c}{V} + \frac{p_{0}}{n} - \frac{p_{0}'}{n}\right)^{2} - (m_{0}c)^{2}}\right)\cos\gamma$$
(14)

Copyright © 2011 SciRes.

Let's consider some special cases of transformation (10) and (14) which are not excessively cumbersome, and allow to lead the analysis of results obviously enough.

3. Absorption of a Photon by Moving Free Electron

Let's we assume, that a photon falling on moving free electron, it is completely absorbed. Thus the condition $p'_0 = 0$ should satisfy. From (14) we find:

$$p^{2} + p_{0}^{2} + 2pp_{0}\cos\varphi = \left(p\frac{c}{V} + \frac{p_{0}}{n}\right)^{2} - \left(m_{0}c\right)^{2}.$$
 (15)

Carrying out transformations, we shall receive:

$$p^{2}\left[1-\left(\frac{c}{V}\right)^{2}\right]+p_{0}^{2}\left(1-\frac{1}{n^{2}}\right)$$

$$=2pp_{0}\left(\frac{c}{nV}-\cos\varphi\right)-\left(m_{0}c\right)^{2}$$
(16)

Using a formula $p = m_0 V / \sqrt{1 - \frac{V^2}{c^2}}$, we have:

$$\frac{\left(m_{0}V\right)^{2}}{\left(1-\frac{V^{2}}{c^{2}}\right)^{2}}\left[1-\left(\frac{c}{V}\right)^{2}\right]+p_{0}^{2}\left(1-\frac{1}{n^{2}}\right)$$

$$=2pp_{0}\left(\frac{c}{nV}-\cos\varphi\right)-\left(m_{0}c\right)^{2}$$
(17)

It is possible to see, that the first addend in the left part (17) is equal $-(m_0 c)^2$. Hence:

$$\cos\varphi = \frac{c}{nV} - \frac{p_0}{2p} \left(1 - \frac{1}{n^2} \right).$$
(18)

The received expression is analogue of the formula for an angle of the radiated photon, uniformly moving in medium electron in Vavilov's—Cherenkov's effect [6]. Difference is only in a sign a minus in the right part (18) that it is connected to change of a direction of a photon impulse p_0 .

From the formula (18) follows, that in vacuum, *i.e.* at n = 1, free electron cannot to absorb a photon, since $\cos \varphi > 1$. But in substance, at a condition:

$$\frac{c}{nV} - \frac{p_0}{2p} \left(1 - \frac{1}{n^2} \right) \le 1,$$
(19)

such process is possible.

Let's consider in more detail a condition of absorption of a photon by electron. We shall transform:

$$\frac{p_0}{p} = \frac{nh\nu}{mVc} \approx \frac{nh\nu}{m_0Vc} = \frac{n\Lambda\nu}{V} = \frac{n\Lambda c}{Vn\lambda} = \frac{\Lambda}{\lambda}\frac{c}{V}, \quad (20)$$

where size $\Lambda = \frac{h}{m_0 c} = 0.002426$ nm is Compton's length

of an electron wave. It is accepted also, that speed of electron is much less than speed of light in vacuum, so $m \approx m_0$. Hence, the formula (19) becomes:

$$\frac{\Lambda}{\lambda} \left(1 - \frac{1}{n^2} \right) \ge \frac{2}{n} \,. \tag{21}$$

For $0 < n \le 1$ an inequality (19) and (21) it is not correct. If an index of refraction n > 1, we receive:

$$\lambda \le \frac{\Lambda}{2} \left(n - \frac{1}{n} \right). \tag{22}$$

The formula (22) shows, that absorption of a photon, moving in substance with an index of refraction n > 1 free electron it is possible, probably, only for X-ray radiation. And, the more a index of refraction, the bigger length of a wave can to absorb electron.

For lines of the substances resulted in [7], index of refraction for X-ray radiation is less unit, and it is very close to unit. The vacuum for x-ray radiation is in comparison with these substances optically is large dense medium. For X-ray radiation absorption of photons by electron according to (22) is impossible. However, there are no physical bases to deny existence of substances with a index of refraction to more unit for X-ray radiation. Free electrons in such substances could absorb photons of X-ray radiation.

4. Scattering of a Photons on Free Electron

1) Coherent scattering

Coherent scattering of a photon on moving free electron arises, if the condition $p'_0 = p_0$ is observed. The formula (5) will be transformed to a kind p = p', *i.e.* electron changes only a direction of the movement. In this case the Equation (4) have form:

$$p_0 \sin \frac{\theta}{2} = p \sin \frac{\alpha}{2}.$$
 (23)

The Equation (12) results to a condition:

$$\cos\varphi = \cos\gamma \,. \tag{24}$$

The condition (24) shows, there are $\varphi = \gamma$ at coherent scattering. The given condition actually is the law of reflection of the photon, similar to the law of reflection of light from mirror interface of two media.

Formulas (23) and (24) also show, that at coherent scattering the quadrangle of impulses on **Figure 1** to become symmetric, **Figure 2**.

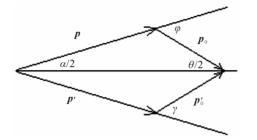


Figure 2. Direction of vectors of impulses of particles at coherent scattering of a photon on free electron.

2) Compton effect

Parameters of Compton effect are provided that electron is motionless [6], *i.e.* p = 0, V = 0,

$$\frac{pc}{V} = m_0 c \bigg/ \sqrt{1 - \frac{V^2}{c^2}} = m_0 c$$

In this case the formula (10) becomes:

$$p_0^{2} + p_0^{\prime 2} - 2p_0 p_0^{\prime} \cos \theta = \left(m_0 c + \frac{p_0}{n} - \frac{p_0^{\prime}}{n}\right)^2 - \left(m_0 c\right)^2.$$
(25)

Carrying out the further transformations (25), we find:

$$p_{0} - p_{0}$$

$$= \frac{n}{2m_{0}c} \left[\left(p_{0}^{2} + p_{0}^{\prime 2} \right) \left(1 - \frac{1}{n^{2}} \right) + 2p_{0}p_{0}^{\prime} \left(\frac{1}{n^{2}} - \cos\theta \right) \right].$$
(26)

Adding and subtracting unit in brackets of last addend (26), we shall receive:

$$p_0 - p'_0 = \frac{n}{2m_0c} \left(p_0 - p'_0 \right)^2 \left(1 - \frac{1}{n^2} \right) + \frac{np_0p'_0}{m_0c} \left(1 - \cos\theta \right).$$
(27)

Let's designate a difference of impulses falling and scattered photons on electron in vacuum [8]:

$$\Delta p_{vac} = \frac{p_0 p'_0}{m_0 c} \left(1 - \cos\theta\right). \tag{28}$$

Thus, the formula (27) can be written down as:

$$\Delta p_0 = \frac{n}{2m_0 c} \Delta p_0^2 \left(1 - \frac{1}{n^2} \right) + n \Delta p_{vac} , \qquad (29)$$

where $\Delta p_0 = p_0 - p'_0$ a difference of impulses falling and scattered photons on electron in substance.

Let's designate dimensionless differences of photons $\Delta p_0 = \Delta \Delta v_0$

impulses $\eta = \frac{\Delta p_0}{m_0 c} = \frac{\Lambda \Delta v_0}{c}$ in substance and

$$\eta_{vac} = \frac{\Delta p_{vac}}{m_0 c} = \frac{\Delta \Delta v_{vac}}{c}$$
 in vacuum. Values Δv_0 and

 Δv_{vac} is change of frequency at scattering of a photon,

Copyright © 2011 SciRes.

accordingly, in substance and in vacuum.

In new designations the Equation (29) can be copied:

$$\frac{1}{2}n\eta^{2}\left(1-\frac{1}{n^{2}}\right)-\eta+n\eta_{vac}=0.$$
 (30)

Solving a quadratic Equation (30), we find:

$$\eta = \frac{n}{n^2 - 1} \left[1 - \sqrt{1 - 2(n^2 - 1)\eta_{vac}} \right].$$
(31)

The sign plus before a root is unacceptable, since in this case differences of photons impulses in substance and in vacuum have a different directivity of change. As it will be from the further, to use in (31) expansion of a

root as $\sqrt{1-x} \approx 1 - \frac{1}{2}x$ is not informatively for the analysis of influence of characteristics of substance on not coherent scattering of a photon.

Let's find in an obvious form change of length of photons wave at their not coherent scattering. We shall transform values:

$$\eta = \frac{\Lambda \Delta v_0}{c} \approx -\frac{\Lambda}{c} dv_0 = -\frac{\Lambda}{c} d\left(\frac{V_{\phi}}{\lambda}\right), \qquad (32)$$
$$= -\frac{\Lambda}{n} d\left(\frac{1}{\lambda}\right) = \frac{\Lambda}{n\lambda^2} d\lambda \approx \frac{\Lambda}{n\lambda^2} \Delta\lambda, \qquad (32)$$
$$\eta_{vac} = \frac{\Lambda \Delta v_{vac}}{c} \approx -\frac{\Lambda}{c} dv_{vac} = -\frac{\Lambda}{c} d\left(\frac{c}{\lambda_{vac}}\right), \qquad (33)$$
$$= \frac{\Lambda}{\lambda_{vac}^2} d\lambda_{vac} \approx \frac{\Lambda}{\lambda_{vac}^2} \Delta\lambda_{vac}$$

where λ and λ_{vac} is a lengths of photons waves at it scattering on electrons in substance and in vacuum. Signs a minus are connected by that values $\Delta v_0 = v - v'$ and $\Delta v_{vac} = v_{vac} - v'_{vac}$ used positive. Besides it is designated $\Delta \lambda = \lambda' - \lambda \approx d\lambda$ and $\Delta \lambda_{vac} = \lambda'_{vac} - \lambda_{vac} \approx d\lambda_{vac}$ where dashes mean lengths of photons waves after interaction with electron.

Using determination of an index of refraction as $n = \frac{\lambda_{vac}}{\lambda}$ [7], we shall transform (33) to a form:

$$\eta_{vac} \approx \frac{\Lambda}{n^2 \lambda^2} \Delta \lambda_{vac} \,, \tag{34}$$

Inserting (32) and (34) into (31), we'll find:

$$\Delta \lambda = \frac{n^2}{n^2 - 1} \frac{\lambda^2}{\Lambda} \left[1 - \sqrt{1 - 2\frac{\left(n^2 - 1\right)}{n^2} \frac{\Lambda}{\lambda^2}} \Delta \lambda_{\text{вак}} \right].$$
(35)

In expansion of a root in the formula (35) up to the first order we have $\Delta \lambda = \Delta \lambda_{vac}$, therefore it is necessary to execute the expansion up to the second order

$$\sqrt{1-x} \approx 1 - \frac{1}{2}x - \frac{1}{8}x^2$$
. In result we shall find:

$$\Delta \lambda \approx \Delta \lambda_{vac} + \frac{1}{2} \left(\frac{n^2 - 1}{n^2} \right) \frac{\Lambda}{\lambda^2} \Delta \lambda_{vac}^2 \,. \tag{36}$$

where $\Delta \lambda_{vac} = \Lambda (1 - \cos \theta)$, for example [1].

Thus, in the second order of expansion (35) change of length of a wave at not coherent scattering depends from the length of a wave.

Let's note, that in X-ray area of a spectrum of radiation an index of refraction is n < 1. Therefore, according to (36) $\Delta \lambda < \Delta \lambda_{varc}$.

Let's designate, following [7], unit decrement of an index of refraction $\delta = 1 - n$. Then the formula (36) will be transformed to a form:

$$\Delta \lambda \approx \Delta \lambda_{vac} - \frac{1}{2} \frac{\delta}{\lambda^2} \left(\frac{n+1}{n^2} \right) \Delta \Delta \lambda_{vac}^2 .$$
 (37)

The index of refraction of substances in a X-ray range is close to unit [7], therefore the formula (37) can be copied as follows:

$$\frac{\Delta\lambda_{vac} - \Delta\lambda}{\Delta\lambda_{vac}} = \frac{\delta}{\lambda^2} \Lambda \Delta\lambda_{vac} \,. \tag{38}$$

The formula (38) shows, that in a X-ray range influence of substance on parameters of not coherent scattering is very insignificant and at practical calculations it is possible to use classical formula of Compton

 $\Delta \lambda \approx \Delta \lambda_{\text{sax}} = \Lambda (1 - \cos \theta)$. The opportunity of use of Compton formula in substance mathematical is determined by that expansion (35) up to the first order parameters of substance is absent.

Proceeding from the classical theory of a Lorentz dispersion [7], we have in a X-ray range dependence

 $\frac{\delta}{\lambda^2} = 2,7 \cdot 10^{11} \frac{\rho Z}{A} \left[m^{-2} \right], \text{ where } \rho \text{ is density of sub-$

stance in kg/m³, Z is charging number, A is mass number. Using $\Lambda = 2,426 \cdot 10^{-12}$ m, we find:

$$\varepsilon = \frac{\Delta \lambda_{vac} - \Delta \lambda}{\Delta \lambda_{vac}} = 0,655 \frac{\rho Z}{A} \Delta \lambda_{vac} \,. \tag{39}$$

Relative change of shift of wave length ε in substance in comparison with vacuum at not coherent scattering in a X-ray range of a spectrum of electromagnetic radiation is proportional to density of substance, its charging number and in inverse proportion to nuclear mass of substance. It also linearly depends from Compton's shift of length of a wave in vacuum.

5. The Stop of Free Electron after It Interaction with Photon

Let's find conditions at which free electron after interac-

Copyright © 2011 SciRes.

tion with a photon in substance will be stop, *i.e.* p' = 0. In this case the Equation (5) gets a form:

$$\frac{pc}{V} + \frac{p_0}{n} = m_0 c + \frac{p'_0}{n} \,. \tag{40}$$

From the Equation (40) we'll find:

$$p = \left(nm_0 c - \Delta p_0\right) \frac{V}{nc},\tag{41}$$

where, as before, $\Delta p_0 = p_0 - p'_0$. Formulas (4) and (12) change:

$$p_0^{2} + p_0^{\prime 2} - 2p_0 p_0^{\prime} \cos\theta = p^2, \qquad (42)$$

$$p^{2} + p_{0}^{2} + 2pp_{0}\cos\varphi = {p_{0}^{\prime}}^{2}.$$
 (43)

Using $p'_0 = p_0 - \Delta p_0$ the Equation (43) it is possible to transform to a form:

$$p^{2} + \Delta p_{0} \left(2p_{0} - \Delta p_{0} \right) + 2pp_{0} \cos \varphi = 0.$$
 (44)

Solving in common (42) and (43) we shall find:

$$p_0 + p\cos\varphi - p'_0\cos\theta = 0, \qquad (45)$$

that is equivalent:

$$p_0 (1 - \cos \theta) + p \cos \phi + \Delta p_0 \cos \theta = 0.$$
 (46)

Finding from (46) value p_0 and inserting it in (44), we have:

$$p^{2} - \Delta p_{0}^{2} - \frac{2p^{2}\cos^{2}\varphi}{(1-\cos\theta)} - \Delta p_{0}\frac{2p\cos\theta\cos\varphi}{(1-\cos\theta)}$$

$$-\Delta p_{0}\frac{2p\cos\varphi}{(1-\cos\theta)} - \Delta p_{0}^{2}\frac{2\cos\theta}{(1-\cos\theta)} = 0.$$
(47)

Transform (47), we'll find:

$$\Delta p_0^2 - 2\Delta p_0 p \cos \varphi - p^2 \left(1 - \frac{2\cos^2 \varphi}{(1 - \cos \theta)} \right) \frac{(1 - \cos \theta)}{(1 + \cos \theta)} = 0.$$
(48)

Solving a quadratic Equation (48), we find:

$$\frac{\Delta p_0}{p} = \cos\varphi + \sqrt{\left(\cos\varphi\right)^2 + \frac{\left(1 - \cos\theta\right)}{\left(1 + \cos\theta\right)} \left(1 - \frac{2\cos^2\varphi}{\left(1 - \cos\theta\right)}\right)}.$$
(49)

As it will be shown below, before a root it is necessary to take a sign plus.

After simple transformations of the formula (49), we shall find:

$$\frac{\Delta p_0}{p} = \frac{\cos\left(\varphi - \frac{\theta}{2}\right)}{\cos\frac{\theta}{2}}.$$
(50)

The formula (50) shows, that at a head-on interaction

JMP

of a photon and moving electron, *i.e.* at $\varphi = 180^{\circ}$ in case if electron after interaction stops, a direction of an scattered photon equiprobably in all directions. The sign plus before a root in the formula (49) is connected with geometrical interpretation of the formula (50), **Figure 3**.

Let's marked, that in case if electron after interaction with a photon does not stop, the formula (50) according to the similar simple geometrical analysis **Figure 1**, has more a general form:

$$\left|\Delta p_0 \cos\frac{\theta}{2}\right| = \left|p' \cos\left(\gamma - \frac{\theta}{2}\right) - p \cos\left(\varphi - \frac{\theta}{2}\right)\right|.$$
(51)

At use (41) formula (50) will be transformed to a form:

$$\Delta p_{0} = \frac{\cos\left(\varphi - \frac{\theta}{2}\right)}{\cos\frac{\theta}{2}} \frac{\left(m_{0}V\right)}{\left[1 + \frac{\cos\left(\varphi - \frac{\theta}{2}\right)}{\cos\frac{\theta}{2}}\frac{V}{nc}\right]}.$$
 (52)

Similarly (32), we shall pass from change of an impulse of photons Δp_0 to the shift of lengths of waves:

$$\frac{\Delta p_0}{m_0 c} = \frac{\Lambda \Delta v_0}{c} \approx -\frac{\Lambda}{c} dv_0 = \frac{\Lambda}{n\lambda^2} \Delta \lambda .$$
 (53)

Hence, from (52) we shall find:

$$\Delta \lambda = \frac{n\lambda^2 \Delta p_0}{\Lambda m_0 c} = \frac{\lambda^2}{\Lambda} \frac{\cos\left(\varphi - \frac{\theta}{2}\right)}{\cos\frac{\theta}{2}} \frac{nV/c}{\left[1 + \frac{\cos\left(\varphi - \frac{\theta}{2}\right)}{\cos\frac{\theta}{2}} \frac{V}{nc}\right]}.$$
(54)

The formula (54) is a special case of not coherent scattering when electron on which there is a scattering, after interaction with a photon stops. Energy of movement electron is transferred an scattered photon, therefore its frequency increases, and the length of a wave

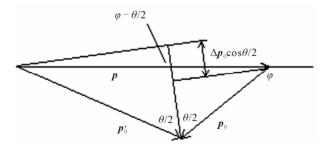


Figure 3. Geometrical interpretation of interaction of a photon with stopping electron.

Copyright © 2011 SciRes.

decreases on value $\Delta \lambda$.

As considered the case when electron stops, an angle is 90° < φ < 270°, **Figure 3**. Average value of angle φ takes place in case of a head-on interaction of X-ray photon ($n \approx 1$) and electron, therefore further we shall be limited to this case. At $\varphi = 180^\circ$, we have:

$$\Delta \lambda = -\frac{\lambda^2}{\Lambda} \frac{V/c}{\left(1 - \frac{V}{c}\right)} \,. \tag{55}$$

The sign a minus specifies on reduction of length of a wave after interaction of a photon with stopping electron.

Let's note, that the formula (55), really, can lead to idea - describe Compton effect by Doppler's shift of length of a wave [2]. However, dependence of shift of length of a wave from angle θ in the formula (54) differs from its in [2] which, proceeding from Doppler effect is received as $\Delta \lambda \sim (1-\cos \theta)$.

Change of shift of lengths of waves due to movement of electrons at not coherent scattering, taking into account (55), is equal:

$$\delta \lambda \approx \Delta \lambda_{\scriptscriptstyle BAK} - \frac{\lambda^2}{\Lambda} \frac{V/c}{\left(1 - \frac{V}{c}\right)} = \Lambda \left[1 - \cos\theta - \left(\frac{\lambda}{\Lambda}\right)^2 \frac{V/c}{\left(1 - \frac{V}{c}\right)} \right].$$
(56)

On **Figure 4** the curve of dependence of shift $\Delta \lambda$ for not coherent line of scattered radiation from angle θ [7] is shown. The experimental points received for carbon at length of a wave of radiation $\lambda = 0.71 \cdot 10^{-11}$ m are shown also. From the analysis of curves follows, that at angle θ close to 180° rather big reduction of shift $\Delta \lambda$ is observed. This reduction can be explained only by movement of electron, since from the formula (37) cannot get

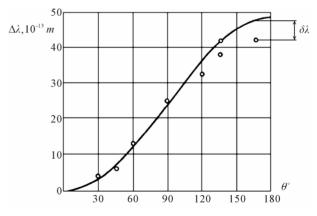


Figure 4. Experimental dependence of shift of not coherent line of scattered X-ray radiation from an angle of scattering of photons θ . Continuous line is calculation under formula of Compton $\Delta \lambda_{vac} = \Lambda (1 - \cos \theta)$.

such big reduction $\Delta \lambda$. Delay of movement of electron at interaction with a photon results to increase of energy of a scattered photon, *i.e.* reduction value $\Delta \lambda$ on size $\delta \lambda$. The estimation which has been carried out under the formula (56) shows, that value $\delta \lambda \approx 7 \cdot 10^{-13}$ m observable in experiment at an angle $\theta \approx 180^{\circ}$ arises at initial speed of electron $V = 0.6 \cdot 10^{6} \frac{\text{m}}{\text{m}}$. This speed is close

tial speed of electron $V = 0.6 \cdot 10^6 \frac{\text{m}}{\text{s}}$. This speed is close

to speed to Fermi for electron which can be analogue of speed of thermal movement of electronic gas [9].

Calculation under the formula (56), obviously, has estimated character. First, there is thermal distribution Fermi-Dirac of speeds of electronic gas in substance, therefore photons are necessary for a stop of electrons must be various energy. Second, the head-on interaction of photons and electrons is enough rare. Thirdly, the stop of electron after interaction with a photon also is occurrence enough rare. In a general change of an impulse of a photon at not coherent scattering is determined by the law (51). Strictly speaking stop in electronic gas two electrons can only with various directions spin. It follows from a principle of veto of Pauli. Hence, there should be a statistical distribution of size $\delta \lambda$ which is stochastic ways form of its numerical value. However, taking into account, that the head-on interaction of a photon and electron is observed, when angle φ achieves average value 180° in considered process (90° < ϕ < 270°), and also close to speed of electron under the formula (56) to speeds of Fermi, it is possible to assume that an estimation of size $\delta \lambda$ under this formula is allowable at angle θ close to 180°.

6. Conclusions

At interaction of a photon with free electron in substance various results are possible: absorption of a photon, coherent and not coherent scattering, a stop of electron.

Thus absorption of a photon in a X-ray range probably only in substance with an index of refraction is more Change of length of a wave at not coherent scattering of X-ray radiation in substance in the second order depends from the length of a wave, as against not coherent scattering in vacuum. Influence of parameters of substance: density, charging and mass numbers on shift of lengths of waves of falling and scattered radiations very insignificantly.

Observable reduction of shift of lengths of waves at not coherent scattering can be explained only by movement of free electrons, interacting with photons. At a head-on interaction of a photon and moving electron in case electron after interaction stops, a direction of a scattered photon equiprobably in all directions.

7. References

- [1] E. H. Wichmann, "Quantum Physics," Science, Moscow, 1977, p. 156.
- [2] R. Kidd and J. Ardini, A. Anton, "Compton Effect as a Double Doppler Shift," *American Journal of Physics*, Vol. 53, No.7, 1985, p. 641. doi:10.1119/1.14274
- [3] A. H. Compton, "Secondary Radiations Produced by X-Rays," The National Research Council of the National Academy of Sciences, Washington, DC, 1922, pp. 18-20.
- [4] V. B. Berestetsky, E. M. Lifshits and L. P. Pitaevskij, "Quantum Electrodynamics," Science, Moscow, Vol. 4, 1989, p. 402.
- [5] A. N. Volobuev and E. S. Petrov, "Angular Distribution of Photoelectrons during Irradiation of Metal Surface by Electromagnetic Wave," *Journal of Modern Physics*, Vol. 2, No. 8, 2011, pp. 780-786. doi:10.4236/jmp.2011.28091
- [6] V. G. Levich, "The Course of Theoretical Physics," *PhysMathGiz*, Moscow, Vol. 1, 1962, pp. 246-302.
- [7] M. A. Blohin, "Physics of Roentgen Rays," State Publishing House, Moscow, 1957, pp. 192-301.
- [8] A. N. Matveev, "Nuclear Physics," The Higher School, Moscow, 1989, p. 27.
- [9] N. Ashcroft and N. Mermin, "Solid-State Physics," Holt, Rinehart and Winston, New York, 1976, p. 49.