

# Molecular Dynamics Simulation of Grain Refinement in a Polycrystalline Material under Severe Compressive Deformation

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## Abstract

Grain refinement in a polycrystalline material resulting from severe compressive deformation was simulated using molecular dynamics. A simplified model with four square grains surrounded by periodic boundaries was prepared, and compressive deformation was imposed by shortening the length in the y direction. The model first deformed elastically, and the compressive stress increased monotonically. Inelastic deformation was then initiated, and the stress decreased drastically. At that moment, dislocation or slip was initiated at the grain boundaries or triple junction and then spread within the grains. New grain boundaries were then generated in some of the grains, and sub-grains appeared. Finally, a microstructure with refined grains was obtained. This process was simulated using two types of grain arrangements and three different combinations of crystal orientations. Grain refinement generally proceeded in a similar fashion in each scenario, whereas the detailed inelastic deformation and grain refinement behavior depended on the initial microstructure.

## **Keywords**

Molecular Dynamics, Polycrystalline Material, Microstructure, Grain Boundary, Grain Refinement, Severe Plastic Deformation, Computer Simulation

## **1. Introduction**

The microstructures of metallic materials determine the mechanical properties. Thus, controlling the microstructure is important in materials processing. Heat treatments such as annealing, quenching, and tempering have long been used for this purpose, and various advanced techniques have been developed. These processes result in phase transformation along with grain refinement or coarsening, allowing both the phase and grain size to be controlled. However, since heat treatment requires large equipment and significant energy inputs, alternative methods are desired. One such alternative is mechanical processing, and various mechanical processing methods have been developed [1] [2], including equal-channel angular pressing [3] [4], accumulative roll bonding [5] [6], and asymmetric rolling [7]. In all these processes, the material is subjected to severe plastic deformation (SPD), and the microstructure is crushed by compressive load. Consequently, small crystal grains with sizes on the order of sub-microns or nanometers are obtained, generally resulting in remarkable improvements in material strength. However, the mechanism underlying the change in microstructure is very complicated, and the optimal conditions for microstructural control have only been empirically explored.

To clarify the mechanism of grain refinement and develop more effective refinement processes, many types of investigations have been conducted. For example, electron backscatter diffraction (EBSD) method can be used to study the polycrystalline microstructure [8] [9]. However, *in situ* observation of dynamic behavior remains difficult, and computer simulation is indispensable for this purpose. Molecular dynamics (MD), which treats the motions of all individual atoms in the system considered, can effectively capture temporal changes in microstructure. To date, MD simulations have been used to clarify the fundamental mechanism of plastic deformation, including the generation and motion of dislocations, grain-boundary migration, changes in crystal orientation, and their relationships and interactions [10]-[11]. The author has also used MD simulations to investigate the atomistic mechanisms of phase transformation and plastic deformation. For example, an MD study of grain boundary stability revealed that the grain shapes change to reach a stable structure [16]. In a model of a polycrystalline shape-memory alloy, deformation was found to occur individually in every grain [17]. Transformation-induced plasticity was also simulated, and the dependence of strain on applied load was discussed [18]. Following these simulations, this study focuses on changes in microstructure resulting from severe deformation. Prior to this study, we reported on changes in microstructure under tensile load, during which grain-boundary migration and changes in grain shape were observed [19]. To study a more practical scenario for engineering applications, we evaluate severe compression in this study [20] and simulate the grain refinement process. This paper also presents a detailed investigation of the initiation of plastic deformation in a polycrystalline material.

# 2. Simulation Method, Model, and Conditions

## 2.1. Fundamental Equations

Classical MD method is applied in this study. Newton's equation of motion for all atoms in the system is solved numerically using the following fundamental equation:

$$\ddot{\boldsymbol{r}}_{i} = \frac{1}{m_{i}} \boldsymbol{F}_{i} = \frac{1}{m_{i}} \sum_{j} \boldsymbol{f}_{ij}, \quad \boldsymbol{f}_{ij} = -\frac{\mathrm{d}\phi_{ij}}{\mathrm{d}r} \frac{\boldsymbol{r}_{ij}}{|\boldsymbol{r}_{ij}|}, \quad (1)$$

where,  $m_i$  and  $r_i$  are the mass and position vector of the *i*-th atom, respectively,  $F_i$  is the force acting on the *i*-th atom,  $f_{ij}$  is the force acting on the *i*-th atom from the *j*-th atom,  $\phi_{ij}$  is the interatomic potential between the *i*-th and the *j*-th atoms, and  $r_{ij}$  is the position vector drawn from the *i*-th to the *j*-th atoms. The following Lennard-Jones function, which is often used for general investigation of fcc materials, is applied in this study:

$$\phi = 4D((R/r)^{12} - (R/r)^6), \qquad (2)$$

where, D and R are energy and length parameters, respectively. This function is standardized in the non-dimensional form of

$$\phi^* = r^{*-12} - r^{*-6}, \qquad (3)$$

where  $\phi^* = \phi/(4D)$  and  $r^* = r/R$ . All other related physical quantities such as stress and temperature are also non-dimensionalized using these standards. The values presented in this study are in the non-dimensional ones, while the asterisk notion is omitted hereafter.

## 2.2. Simulation Model

The simulation model is illustrated in **Figure 1**. The calculation model is divided into four domains and arranged in two different ways (Types A and B), as shown in **Figure 1(a)** and **Figure 1(b)**. Compression is imposed along the *y* axis, and Types A and B differ based on whether the straight grain boundary traversing the model is perpendicular (Type A) or parallel (Type B) to the loading direction. The model is three-dimensional in x-y-z space, and periodic boundary conditions are applied in all directions. The atoms are set on the fcc lattice points so that the (001) plane is on the x-y plane, and the [100] direction is along the *x* axis. The atoms are then rotated around the *z* axis, and the rotation angles are provided to the four grains individually. Consequently, grain boundaries are created as the misfit of the crystal arrangement at the domain boundaries. Note that all grain boundaries are perpendicular to the x-y plane, and no



Figure 1. Illustration of the simulation model: (a) Type A; (b) Type B; (c) MD model.

boundaries parallel to the x-y plane are considered; *i.e.* all grains correspond to infinitely long columnar grains because of the periodic boundary in the z direction. Figure 1(c) shows a sample of the atomic arrangement colored based on potential energy; the grain boundaries in Figure 1(c) are depicted by green or red atoms.

Compression is imposed by shortening the length in the *y* direction  $(L_y)$  at a constant rate. At the same time, the stress components  $\sigma_{xx}$  and  $\sigma_{zz}$  are kept equal to zero by adjusting the *lengths* in the *x* and *z* directions  $(L_x \text{ and } L_x \text{ respectively})$ . Shear deformation is not considered, and the model retains its rectangular parallelepiped form.

The model size is set as  $50 \times 50 \times 5$  fcc unit cells in the *x*, *y*, and *z* directions, respectively. Some atoms were added or removed around the grain boundary area when constructing the initial polycrystalline arrangement. The rotation angles of Grain *k* (*k* = I, II, III or IV),  $\theta_{k^0}$  are set as shown in **Table 1**. In **Table 1**, the letters in the "Model ID" column represent the grain arrangement (Type A or B), and the number identifies the combination of the rotation angles of the four grains. The angles are selected so that the structure of the grain boundary is regularly ordered. For example, an illustration for the case of  $\theta = \tan^{-1}(1/2)$  is shown in the insert image in **Table 1**. Consequently, some of the boundaries in the model become coincident site lattice (CSL) boundaries. In Models A1 and B1, for example, the grain boundaries between Grains I and II and between Grains III and IV are typical CSL boundaries indicated by the low-index sigma values.

## 2.3. Simulation Conditions

The initially provided atomic positions are not stable sites, especially around the grain boundaries. Therefore, the configuration is relaxed for 15,000 timesteps under a stress-free condition while keeping temperature constant at T = 0.05 in non-dimensional value. Compressive load is then imposed in the *y* direction by decreasing the edge length  $L_y$  at a constant rate of  $\Delta L_y/\tau_s$ , meaning that the total decrease in length  $(\Delta L_y)$  is imposed in the compression period ( $\tau_s$ , in timesteps). In this study, the values are set as  $\Delta L_y = 25.0$  and  $\tau_s = 50,000$ . The length before compression,  $L_y^0$ , is approximately 73.0, although differs slightly for each model. Hence the total compressive strain is approximately 0.34. In the practical SPD processes, the stress state is more complicated including shear, and the order of applied strain reaches 1 by repeating the unit process for several times. In this

Table 1. Rotation angles in Grains I to IV.

Model ID	Grain No.				1
	Ι	II	III	IV	
A1, B1	+1/2	-1/2	+1/3	-1/3	The second
A2, B2	+1/2	-1/3	+1/3	-1/2	Illustration of the
A3, B3	+1/2	+1/3	-1/2	-1/3	case for $x = 1/2$

\*values represent x of  $\theta = \tan^{-1} x$ .

paper, a single process with uniaxial compression is assumed for modeling the fundamental process. Temperature is kept constant at T = 0.05 using velocity-scaling method. Following the compression period, the model is unloaded by releasing the length restriction in the *y*-direction, and the calculation is continued for 5000 more timesteps.

## 3. Simulation Results for Model A1

**Figure 2** shows the variation in the configuration of atoms. Each image in **Figure 2** is a snapshot of atoms projected onto the *x*-*y* plane, and the color indicates the potential energy of each atom. The atoms in blue possess low energy, while those in green, yellow, and red have higher energy. Consequently, the atoms inside the crystal grains are colored in blue, and those on the grain boundaries are depicted in green. **Figure 3(a)** and **Figure 3(b)** show the variation in compressive stress in the loading direction and the stress-strain curve, respectively. The stress variation in **Figure 3(a)** is divided into four stages: Stage I, initial relaxation and monotonic increase; Stage II, stress peak and drastic decrease; Stage III, repeating gradual increases and sudden decreases; and Stage IV, unloading. The results in each stage are detailed in the following sections.

#### 3.1. Stage I—Elastic Deformation

The atomic configuration is initially relaxed under a stress-free condition. As shown in **Figure 2(a)**, the shapes and sizes of the crystal grains are almost maintained, although the grains change slightly in shape from perfect square to slightly rounded squares as a result of the blunting of the angles at triple junctions. We can also see that every grain boundary has a different structure; continuous straight boundaries are observed between Grains I/II, III/IV, I/IV, and II/III, whereas discrete boundaries are found between Grains I/III and II/IV.



Figure 2. Variation in microstructure under compression for Model A1.



Figure 3. Variation of compressive stress (a) and stress-strain curve (b) for Model A1.

An external compressive load is imposed starting at the 15,000th timestep. The length in the y direction decreases, but the atomic configuration is not affected (**Figure 2(b)**). Since the compressive stress increases monotonically (**Figure 3(b)**), deformation in this stage is regarded as elastic deformation.

## 3.2. Stage II—Yielding

At the 24000th timestep, three regions indicated by different color densities, marked by circles in **Figure 2(c)**, are generated. The color is still the same, indicating that the energy has not changed at this moment. Therefore, the atoms are considered to retain regularity in the local arrangement. The difference in appearance is attributed to the change in orientation. The original grains are all arranged so that the (001) plane is on the *x*-*y* plane, and the atoms on multiple layers in the *z* direction overlap completely. Consequently, the snapshot appears to have discrete dots with large intervals. When the orientation of the crystals (*i.e.*, the [001] direction) is dislocated from the *z* axis, all atoms on the accumulated layers do not overlap and are individually depicted in the snapshots. Therefore, a change in orientation is considered to have occurred in the dense blue region of **Figure 2(c)**. Hereafter, we refer to such a region as the deformation region. All the deformation regions are generated at grain boundaries; one is at the triple junction, and the others are on the grain boundaries.

These regions spread rapidly inside the grains, and other deformation regions are successively generated. Finally, the deformation regions occupy the entire model by the 26,000th timestep, as shown in **Figure 2(d)**. At the same time, the stress decreases drastically, as shown in **Figure 3**; hence, the behavior can be regarded as yielding. The decrease in stress is drastic but not momentary; it takes approximately 2000 timesteps from start to end. The detailed changes in atomic configuration along with a close-up image of the stress curve are shown in **Figure 4**. No specific change in orientation or deformation is observed at the 20,000th timestep. Some small indications (marked by circles) start appearing at the 22,000th timestep; however, these regions do not spread quickly. The stress value continues to increase, and no specific variation is observed. At the 23,000th



Figure 4. Detailed variation in microstructure around the first stress peak for Model A1.

timestep, another initiation is observed at the left-bottom corner of Grain III (marked with a red circle), and the stress stops increasing at this moment. This deformation region spreads quickly, as shown in **Figure 4(e)** to **Figure 4(h)**, and the stress begins to decrease. Deformation regions originated from other points are also activated and extend to the entire domain by the 26,000th timestep, and the stress value is minimized at this point.

In the above process, two types of deformation regions are observed: linear deformation regions like those marked with yellow circles in **Figure 4(c)**; and planar regions like the one marked with the red circle in **Figure 4(e)**. These deformation modes correspond to dislocation and slip, respectively. Since the periodic boundary condition in the depth (z) direction may affect the behavior, more detailed discussion is reserved for future work.

## 3.3. Stage III—Plastic Deformation and Grain Refinement

At the completion of the extension of the deformation region at the 26,000th timestep, some defects or mismatches in crystal orientation have been introduced. For instance, in Grain II, a vertical grain boundary is generated, as shown in **Figure 2(d)**. In other grains, as well as similar new boundaries, transversal oblique lines are observed. These misalignments in crystal structure caused by abrupt deformation are relieved by the 30,000th timestep (**Figure 2(e)**). In this time duration, the stress value again increases and reaches the second peak value followed by a drastic decrease. The snapshot after the second peak at the 36000th timestep (**Figure 2(f)**) shows significantly distorted grain shapes. In particular, the original vertical grain boundaries between Grains I/II and III/IV are continuously curved and become S-shaped, as shown in **Figure 2(g)**. In addition, several new boundaries and defects are generated in all grains. In the stress curve, no notable peak like the first and second peaks is observed; instead, a gradual increase with intermittent increases and decreases is found. Finally, the model is divided into many fine grains consisting of various sizes and shapes of grains, as shown in **Figure 2(k)**.

## 3.4. Stage IV—Unloading

The addition of compressive load stops at the 65,000th timestep, and the stress in the y direction is released to zero. The edge length in the y direction is then slightly extended. This recovery is slight, and a permanent strain of approximately 0.31 is retained (**Figure 3(b)**). Comparing the snapshots in **Figure 2(k)** and **Figure 2(1)**, the microstructure does not change in this process; hence, the unloading behavior reveals to be elastic. Finally, a refined grain distribution is obtained. As a conclusion, the changes in microstructure resulting from severe compressive deformation are properly simulated.

## 4. Simulation Results for Other Models

## 4.1. Model A2

Simulations for other models are conducted under the same conditions. Figure 5(a) shows the variation in the microstructure of Model A2. As for Model A1, the initial configuration is maintained except for slight relaxation at the grain boundaries, and the model deforms elastically before the 20,000th timestep. Inelastic deformation is then initiated, and at the 24,000th timestep, as shown in Figure 5(a)(iii), a large deformation region depicted in dense blue is generated in Grain III (marked by red circle) along with a small one in Grain II (yellow circle). The larger deformation region is originated at the triple junction, and it seems that this region was dominative for yielding before the 24,000th timestep. The deformation regions including belatedly generated regions spread rapidly, and spreading is complete by the 26,000th timestep (Figure 5(a)(iv)). New grain boundaries are also generated, and Grains II and IV are clearly divided into multiple grains. In contrast, the grain boundary between Grains I and III disappear, and the two grains are connected because of the change in crystal orientation. Continuing the compression, the shapes of grains are varied and divided into fine grains. Finally, a refined microstructure is generated, which retains even when unloaded, as shown in Figure 5(a)(xii). Compared to the results for Model A1, the variation in microstructure is more apparent in Model A2, and the original grain shape is not discriminable. This type of behavior, overall similarity and slight differences, are also observed in the stress-strain curve, as shown in Figure 6. The second peak for Model A2 (green curve) is rather dulled than for A1 (orange curve). Comparing the snapshots before the second peaks for Model A1 and A2 (Figure 2(e) and Figure 5(a)(v), respectively), the grain shape for Model A2 is more distorted, and the original vertical grain boundaries are declined. This is considered to cause of the blunting of the second peak in stress.



**Figure 5.** Variation in microstructure under compression for Models A2 and A3: (a) Model A2; (b) Model A3.

## 4.2. Model A3

**Figure 5(b)** shows the variation in atomic configuration for Model A3. The characteristics in the initial grain boundaries are different from above presented Models A1 and A2; the horizontal grain boundaries between Grains I/III, I/IV, II/III, and II/IV are continuous and form straight lines, while the vertical boundaries are discrete. The elastic deformation behavior is, nevertheless, similar to that observed in the other models, both in atomic arrangement and stress evolution. However, a different feature is found around the first peak in stress (blue curve in **Figure 6(b)**). In Model A3, a short step on the increasing curve is seen



**Figure 6.** Stress-strain curves (a) and the magnified view around the first and second peaks (b) for Models A1-A3.

at the 18,000th timestep, whereas no specific change in atomic configuration is observed at this moment, as shown in the snapshot at the 20,000th timestep (**Figure 5(b)(ii)**). A similar short step appears at the 20,000th timestep, and the stress then reaches the peak at the 23,000th timestep. Unlike in the other models, the stress does not begin to decrease instantly after the peak; instead, it remains high for approximately 3000 timesteps.

The variation in atomic configuration during this stress plateau is shown in **Figure 7**. No defects appear before the 23,000th timestep. A small source of linear defect is generated at the 24,000th timestep (red circle in **Figure 7(d)**). This defect spreads straightly toward the right-bottom direction in Grain II and reaches the opposite grain boundary by the 25,500th timestep. The generation of defects at grain boundaries occurs successively, and the defects spread in the grains to form straight linear defects. Although the lines are all parallel in the grains, the direction is different for every grain; thus, they are assumed to be dislocations. This behavior continues until the 25,500th timestep, which corresponds to the end of the plateau in the stress curve.

Subsequently, the stress decreases drastically at the 26,000th timestep. At this moment, a planar deformation region is generated (**Figure 7(h)**) and then spreads in the grain. Nearly the entire region is occupied by the deformed region by the 28,000th timestep, corresponding to the end of decrease in the stress curve. This behavior is similar to the behavior observed in Models A1 and A2 during the stress decrease after the first peak.

The second peak in stress for Model A3 is not remarkable compared to those in Models A1 and A2. Instead, obvious peaks and decreases in stress are repeated. The horizontal grain boundary traversing the model is quite stable and immobile. The vertical grain boundaries also become clear and remain straight after the first drop in stress, as shown in **Figure 5(b)(v)**. Consequently, deformation region is not generated continuously but as isolated events. In addition, the newly generated grain boundaries are rather ambiguous, and most of them disappear. As a result, the grains are not clearly distinguished at the final stage



Figure 7. Detailed variation in microstructure around the first stress peak for Model A3.

with the exception of the horizontal clear boundary. The crystallographic relationship among grains and grain boundaries relative to the loading direction might have caused the specific behavior observed in this model.

## 4.3. Models B1-B3

**Figure 8** shows the variation in microstructure for Model B1. The stress-strain curves for Models B1, B2, and B3 along with the curve for A1 for comparison are shown in **Figure 9**. The change in microstructure is similar to that observed for Model A1, while the grain size at the 48,000th timestep is even smaller than that for Model A1. In Model B1, the stress level after the first peak in stress is lower on average than in Model A1, as shown in **Figure 9(a)**. Therefore, the resistance for the applied load is weaker in Model B1, resulting in finer grains. Note that the different color densities in **Figures 8(j)-(1)** indicate differences in crystal orientation in the depth (*z*) direction (as explained in Section 3.2), and are insignificant here. A comparison of Models B1–B3 indicates that the first peak-stress value is relatively low for Model B3, as shown in **Figure 9(b)**. Therefore, the variations in atomic configuration around the first stress peak for Models B1 and B3 are compared in **Figure 10**.

As shown in **Figure 10(a)(i)**, the horizontal grain boundaries in Model B1 are continuous, whereas the vertical boundaries between Grains I/III and II/IV are discrete. The gradient in the stress curve is somewhat dulled before the peak, although no corresponding change is observed in atomic configuration. A linear defect appears at the 23,000th timestep but soon disappears. The stress then reaches the peak at the 23,500th timestep. At this moment, multiple linear defects are generated in Grain III, as shown in **Figure 10(a)(iv)**. Several defects are continuously generated in other grains by the 24,500th timestep, while the stress value remains high approximately at the peak value in this duration. Planar



Figure 8. Variation in microstructure under compression for Model B1.



**Figure 9.** Stress-strain curves (a) and the magnified view around the first peak (b) for Models B1-B3.

defects are then generated at the 25,000th timestep (marked by circles), and the stress begins to decrease at this moment. Subsequently, the deformation regions spread to the entire domain by the 27,000th timestep, when the drop in stress is completed.

In Model B3, as shown in Figure 10(b)(i), the vertical straight grain boundaries are continuous, whereas the horizontal boundaries are all discrete. The grain boundary between Grains III/IV (marked with an oval) moves slightly by the 21,000th timestep, at which point the stress is still increasing monotonically before reaching the peak value. At the 21,500th timestep, a linear defect is generated from one of the defects on the grain boundary, as indicated by the red circle in Figure 10(b)(ii). The generation of this first defect occurs earlier than in Model B1. After this point, the deformation behavior is similar to those observed in the other models; multiple linear defects are generated and extended during



**Figure 10.** Detailed variation in microstructure around the first stress peak for Models B1 and B3: (a) Model B1; (b) Model B3.

the stress plateau followed by the spread of planar deformation regions with a drastic decrease in stress. Therefore, the difference in peak stress is attributed to the structures of the grain boundaries: when one of the grain boundaries is easy to become a source of the linear defect or dislocation, the model is easy to yield and the peak value in stress becomes lower.

## **5.** Conclusions

Microstructural changes of a polycrystalline material as a result of severe compressive deformation were simulated using molecular dynamics method. A simplified model consisting of four grains fulfilling the periodic boundary condition was prepared. Two types of grain arrangements and three combinations of crystal orientations were applied. The results obtained are summarized as follows. Following elastic deformation, yielding occurs in some of the grains. The source of the yielding is linear defects initiated at grain boundaries and planar deformation regions which are often generated at triple junctions. When the linear defects or dislocations are generated, the stress stops increasing and tends to remain constant, even as the number of dislocations increases. When a planar deformation defect or slip occurs, the stress decreases drastically until the deformation region extends to the entire model. New boundaries dividing the original grains into two or more grains then appear, and a refined microstructure containing various grain sizes is obtained. The first peak stress depends on the grain arrangement and crystal orientation relative to the loading direction, for which grain boundary structure is considered to be dominative.

The following points remain unclarified: the detailed relationship between the initial microstructure and yield stress; the effect of grain boundary structure on defect generation; the relationships between the initial microstructure and the characteristics of refined grains (size, shape, and distribution); and the change in material properties resulting from the grain refinement. These points will be investigated in future works using the present model, which was validated in this study. The model and method will be applicable to more practical problems for various materials including alloys and composites by using appropriate potential function. Comparison with the experimental results will also become available by accomplishing fully three-dimensional modeling.

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