# Periodic Solutions of a Class of Second-Order Differential Equation 

Zeyneb Bouderbala ${ }^{1}$, Jaume Llibre ${ }^{2}$, Amar Makhlouf ${ }^{1}$<br>${ }^{1}$ Department of Mathematics, University of Annaba, Elhadjar, Annaba, Algeria<br>${ }^{2}$ Departament de Matemàtiques, Universitat Autònoma de Barcelona, Barcelona, Spain<br>Email: zeynebbouderbala@yahoo.fr, jllibre@mat.uab.cat, makhloufamar@yahoo.fr

Received 25 December 2015; accepted 26 February 2016; published 29 February 2016

Copyright © 2016 by authors and Scientific Research Publishing Inc.
This work is licensed under the Creative Commons Attribution International License (CC BY).
http://creativecommons.org/licenses/by/4.0/


Open Access

## Abstract

We study the periodic solutions of the second-order differential equations of the form

$$
\ddot{x}+3 x \dot{x}+x^{3}+F(t)\left(\dot{x}+x^{2}\right)+G(t) x+H(t)=0,
$$

where the functions $F(t), G(t)$ and $H(t)$ are periodic of period $2 \pi$ in the variable $t$.

## Keywords

Periodic Solution, Differential Equation, Averaging Theory

## 1. Introduction and Statement of the Main Results

In this paper we shall study the existence of periodic solutions of the second-order differential equation of the form

$$
\begin{equation*}
\ddot{x}+3 x \dot{x}+x^{3}+F(t)\left(\dot{x}+x^{2}\right)+G(t) x+H(t)=0 \tag{1}
\end{equation*}
$$

where the dot denotes derivative with respect to the time $t$, and the functions $F(t), G(t)$ and $H(t)$ are periodic of period $2 \pi$ in the variable $t$.

We note that the second-order differential Equation (1), when $F=G=H=0$, appears in the Ince's catalog of equations possessing the Painlevé property (see [1]). Moreover, the differential equation $\ddot{x}+3 x \dot{x}+x^{3}=0$ is well known in many areas of mathematics and physics, and it possesses the algebra $\operatorname{sl}(3, \mathbb{R})$ of Lie point symmetries (see for more details in the paper [2] and the references quoted there).

In a recent paper [3] (see also [4] [5]), the second-order differential Equation (1) has been studied when $F=H=0$. A study of coupled quadratic unharmonic oscillators in terms of the Painlevé analysis and inte-
grability can be seen in [6], and studies on the second-order differential equations can be seen in [7]. Other approach to the periodic solutions of second-order differential equations can be found in [8].

Here we study the periodic solutions of the second-order differential Equation (1) when $F(t)=\varepsilon f(t)$, $G(t)=1+\varepsilon g(t)$, and $H(t)=\varepsilon^{k} h(t)$ with $k=1,2$. Our main results are the following ones.
Theorem 1. We define the functions

$$
\begin{align*}
& \mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=-\int_{0}^{2 \pi} F\left(t, X_{0}, Y_{0}\right) \sin t \mathrm{~d} t  \tag{2}\\
& \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=\int_{0}^{2 \pi} F\left(t, X_{0}, Y_{0}\right) \cos t \mathrm{~d} t
\end{align*}
$$

where

$$
\begin{gathered}
F\left(t, X_{0}, Y_{0}\right)=-h(t)-g(t) A(t)-f(t) B(t)-3 A(t) B(t), \\
A(t)=X_{0} \cos t+Y_{0} \sin t \\
B(t)=-X_{0} \sin t+Y_{0} \cos t
\end{gathered}
$$

Assume that the functions $F(t)=\varepsilon f(t), G(t)=1+\varepsilon g(t)$ and $H(t)=\varepsilon^{2} h(t)$ are $2 \pi$-periodic. Then for $\varepsilon \neq 0$ sufficiently small and for every $\left(X_{0}^{*}, Y_{0}^{*}\right)$ solution of the system $\mathcal{F}_{j}\left(X_{0}, Y_{0}\right)=0$ for $j=1,2$, satisfy

$$
\begin{equation*}
\left.\operatorname{det}\left(\frac{\partial\left(\mathcal{F}_{1}, \mathcal{F}_{2}\right)}{\partial\left(X_{0}, Y_{0}\right)}\right)\right|_{\left(x_{0}, Y_{0}\right)=\left(x_{0}^{*}, Y_{0}^{*}\right)} \neq 0 \tag{3}
\end{equation*}
$$

the differential Equation (1) has a $2 \pi$-periodic solution $\quad x(t, \varepsilon)=\varepsilon\left(X_{0}^{*} \cos t+Y_{0}^{*} \sin t\right)+O\left(\varepsilon^{2}\right)$.
Theorem 1 is proved in section 3 using the averaging theory described in section 2 . Two applications of Theorem 1 are the following.

Corollary 1. We consider the differential Equation (1) with $F(t)=\varepsilon\left(1-\cos ^{2} t\right), G(t)=1+\varepsilon \sin ^{2} t$ and $H(t)=\varepsilon^{2} \sin t$. Then for $\varepsilon \neq 0$ sufficiently small, this differential equation has a $2 \pi$-periodic solution $x(t, \varepsilon)=\varepsilon 2(\sin t-\cos t) / 3+O\left(\varepsilon^{2}\right)$.

Corollary 2. We consider the differential Equation (1) with $F(t)=\varepsilon\left(1-\cos ^{2} t+2 \cos ^{4} t\right)$,
$G(t)=1+\varepsilon\left(\sin ^{2} t+2 \sin ^{4} t\right)$ and $H(t)=\varepsilon^{2}\left(\sin t+\sin ^{3} t\right)$. Then for $\varepsilon \neq 0$ sufficiently small, this differential equation has a $2 \pi$-periodic solution $x(t, \varepsilon)=\varepsilon(21 \cos t-7 \sin t) / 20+O\left(\varepsilon^{2}\right)$.

Corollaries 1 and 2 are also proved in section 3.
Theorem 2. Assuming that

$$
\int_{0}^{2 \pi} h(t) \sin t \mathrm{~d} t=0, \quad \int_{0}^{2 \pi} h(t) \cos t \mathrm{~d} t=0
$$

and setting

$$
\begin{align*}
& \mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=-\int_{0}^{2 \pi} f\left(t, X_{0}, Y_{0}\right) \sin t \mathrm{~d} t \\
& \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=\int_{0}^{2 \pi} f\left(t, X_{0}, Y_{0}\right) \cos t \mathrm{~d} t \tag{4}
\end{align*}
$$

with

$$
\begin{gathered}
f\left(t, X_{0}, Y_{0}\right)=-g(t) A(t)-f(t) B(t)-3 A(t) B(t), \\
A(t)=X_{0} \cos t+Y_{0} \sin t-\int_{0}^{t} h(\tau) \sin (t-\tau) \mathrm{d} \tau \\
B(t)=-X_{0} \sin t+Y_{0} \cos t-\int_{0}^{t} h(\tau) \cos (t-\tau) \mathrm{d} \tau
\end{gathered}
$$

Assume that $F(t)=\varepsilon f(t), G(t)=1+\varepsilon g(t)$ and $H(t)=\varepsilon h(t)$ are $2 \pi$-periodic functions. Then for $\varepsilon \neq 0$ sufficiently small and for every $\left(X_{0}^{*}, Y_{0}^{*}\right)$ solution of the system $\mathcal{F}_{j}\left(X_{0}, Y_{0}\right)=0$ for $j=1,2$ satisfy (3), the differential Equation (1) has a periodic solution

$$
x(t, \varepsilon)=\varepsilon\left(X_{0}^{*} \cos t+Y_{0}^{*} \sin t-\int_{0}^{t} h(\tau) \sin (t-\tau) \mathrm{d} \tau\right)+O\left(\varepsilon^{2}\right)
$$

Theorem 2 is proved in section 4. Two applications of Theorem 2 are the following.
Corollary 3. We consider the differential Equation (1) with $F(t)=\varepsilon(\sin (2 t)+\cos (2 t)), G(t)=1+\varepsilon \sin t$ and $H(t)=\varepsilon 2 \cos ^{2} t$. Then for $\varepsilon \neq 0$ sufficiently small, this differential equation has a $2 \pi$-periodic solution

$$
x(t, \varepsilon)=\varepsilon\left((-2 \cos t+15 \sin t) / 31+2 \cos ^{2} t(\cos t-1)\right)+O\left(\varepsilon^{2}\right)
$$

Corollary 4. We consider the differential Equation (1) with $F(t)=\varepsilon \sin t, G(t)=1+\varepsilon \sin ^{2} t$ and $H(t)=\varepsilon 2 \cos (2 t)$. Then for $\varepsilon \neq 0$ sufficiently small, this differential equation has a periodic solution

$$
x(t, \varepsilon)=\varepsilon\left(2(\cos t-1) \cos (2 t)-\frac{8}{5} \sin t\right)+O\left(\varepsilon^{2}\right)
$$

Corollaries 3 and 4 are also proved in section 4.

## 2. Basic Results on Averaging Theory

We state the results from the averaging method that we shall use for proving the results of this work.
We consider differential systems of the form

$$
\begin{equation*}
\boldsymbol{x}^{\prime}=F_{0}(t, \boldsymbol{x})+\varepsilon F_{1}(t, x)+\varepsilon^{2} F_{2}(t, x, \varepsilon) \tag{5}
\end{equation*}
$$

where $\varepsilon$ is a small parameter, and the functions $F_{0}, F_{1}: \mathbb{R} \times \Omega \rightarrow \mathbb{R}^{n}$ and $F_{2}: \mathbb{R} \times \Omega \times\left(-\varepsilon_{0}, \varepsilon_{0}\right) \rightarrow \mathbb{R}^{n}$ are $\mathcal{C}^{2}$ functions, $T$-periodic in the variable $t$, and $\Omega$ is an open subset of $\mathbb{R}^{n}$. Suppose that the unperturbed system

$$
\begin{equation*}
x^{\prime}=F_{0}(t, x), \tag{6}
\end{equation*}
$$

has a submanifold of dimension $n$ of $T$-periodic solutions, i.e. of periodic solutions of period $T$.
We denote by $\boldsymbol{x}(t, \mathbf{z}, 0)$ the solution of system (6) such that $\boldsymbol{x}(0, \mathbf{z}, 0)=\mathbf{z}$. We consider the first variational equation of system (6) on the periodic solution $\boldsymbol{x}(t, \mathbf{z}, 0)$, i.e.

$$
\begin{equation*}
\boldsymbol{y}^{\prime}=D_{x} F_{0}(t, \boldsymbol{x}(t, \mathbf{z}, 0)) \boldsymbol{y} \tag{7}
\end{equation*}
$$

where $y$ is an $n \times n$ matrix. Let $M_{z}(t)$ the fundamental matrix of system (7) such that $M_{z}(0)$ is the identity matrix of $\mathbb{R}^{n}$.

By assumption there exists an open set $V$ such that $\mathrm{Cl}(V) \subset \Omega$ and for each $\mathbf{z} \in \mathrm{Cl}(V), \boldsymbol{x}(t, \mathbf{z}, 0)$ is $T$-periodic. Therefore we have the following result.

Theorem 3. We suppose that there is an open and bounded set $V$ with $\operatorname{Cl}(V) \subset \Omega$ such that for each $\mathbf{z} \in \mathrm{Cl}(V)$, the solution $\mathbf{x}(t, \mathbf{z}, 0)$ is T-periodic, and let $\mathcal{F}: \mathrm{Cl}(V) \rightarrow \mathbb{R}^{n}$ be the function defined by

$$
\begin{equation*}
\mathcal{F}(\mathbf{z})=\int_{0}^{T} M_{z}^{-1}(t) F_{1}(t, x(t, \mathbf{z}, 0)) \mathrm{d} t \tag{8}
\end{equation*}
$$

If there is $\alpha \in V$ with $\mathcal{F}(\alpha)=0$ and $\operatorname{det}((\mathrm{d} \mathcal{F} / \mathrm{d} \mathbf{z})(\alpha)) \neq 0$, then there is a $T$-periodic solution $\boldsymbol{x}(t, \varepsilon)$ of system (5) satisfying $\boldsymbol{x}(t, \varepsilon)=\boldsymbol{x}(t, \mathbf{z}, 0)+O(\varepsilon)$.

Theorem 3 is due to Malkin [9] and Roseau [10], for a new and shorter proof (see [11]).

## 3. Proof of Theorem 1 and Its Two Corollaries

Proof of Theorem 1. Introducing the variable $y=\dot{x}$, we can write the second-order differential Equation (1) as the following first-order differential system

$$
\begin{align*}
& \dot{x}=y \\
& \dot{y}=-3 x y-x^{3}-F(t)\left(y+x^{2}\right)-G(t) x-H(t) \tag{9}
\end{align*}
$$

Doing the rescaling $(x, y)=(\varepsilon X, \varepsilon Y)$, we obtain the system

$$
\begin{align*}
& \dot{X}=Y \\
& \dot{Y}=-X+\varepsilon(-h(t)-g(t) X-f(t) Y-3 X Y)+\varepsilon^{2}\left(-f(t) X^{2}-X^{3}\right) . \tag{10}
\end{align*}
$$

System (10) with $\varepsilon=0$ is the unperturbed system, otherwise system (10) is the perturbed system. The unperturbed system has a unique singular point, the origin of coordinates. The solution $(X(t), Y(t))$ of the unperturbed system such that $(X(0), Y(0))=\left(X_{0}, Y_{0}\right)$ is

$$
X(t)=X_{0} \cos t+Y_{0} \sin t, \quad Y(t)=-X_{0} \sin t+Y_{0} \cos t .
$$

Note that all these periodic orbits have period $2 \pi$. Using the notation introduced in section 2 . We have that $\boldsymbol{x}=(X, Y), \quad \mathbf{z}=\left(X_{0}, Y_{0}\right), \quad F_{0}(x, t)=(Y,-X), \quad F_{1}(x, t)=(0,-h(t)-g(t) X-f(t) Y-3 X Y)$ and $F_{2}(x, t)=\left(0,-f(t) X^{2}-X^{3}\right)$.
The fundamental matrix solution $M_{z}(t)$ is independent of the initial condition $\mathbf{z}$, and denoting it by $M(t)$ we obtain

$$
M(t)=\left(\begin{array}{cc}
\cos t & \sin t \\
-\sin t & \cos t
\end{array}\right) .
$$

Now we compute the function $\mathcal{F}(\mathbf{z})=\left(\mathcal{F}_{1}\left(X_{0}, Y_{0}\right), \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)\right)$ given in (8), and we get the functions (2) of the statement of Theorem 1.

By Theorem 3 each zero $\left(X_{0}^{*}, Y_{0}^{*}\right)$ of system $\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=0$ satisfying (3), provides a $2 \pi$ periodic solution $(X(t, \varepsilon), Y(t, \varepsilon))$ of system (10) with $\varepsilon \neq 0$ sufficiently small such that

$$
(X(t, \varepsilon), Y(t, \varepsilon))=\left(X_{0}^{*} \cos t+Y_{0}^{*} \sin t,-X_{0}^{*} \sin t+Y_{0}^{*} \cos t\right)+O(\varepsilon) .
$$

Going back through the change of variables for every periodic solution $(X(t, \varepsilon), Y(t, \varepsilon))$ of system (10) with $\varepsilon \neq 0$ sufficiently small, we obtain a $2 \pi$-periodic solution $x(t, \varepsilon)=\varepsilon\left(X_{0}^{*} \cos t+Y_{0}^{*} \sin t\right)+O\left(\varepsilon^{2}\right)$ of the differential Equation (1) with $\varepsilon \neq 0$ sufficiently small. This completes the proof of Theorem 1 .

Proof of Corollary 1. We must apply Theorem 1 with

$$
f(t)=1-\cos ^{2} t, \quad g(t)=\sin ^{2} t, \quad h(t)=\sin t .
$$

We compute the functions $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$ of the statement of Theorem 1, and we obtain

$$
\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\frac{\pi}{4}\left(4-3 X_{0}+3 Y_{0}\right), \quad \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=\frac{\pi}{4}\left(-X_{0}-Y_{0}\right)
$$

System $\mathcal{F}_{1}=\mathcal{F}_{2}=0$ has the zero $\left(X_{0}^{*}, Y_{0}^{*}\right)=(2 / 3,-2 / 3)$. Since the Jacobian (3) at this zero is $3 \pi^{2} / 8$, we obtain using Theorem 1 the periodic solution given in the statement of the corollary.

Proof of Corollary 2. We apply Theorem 1 with

$$
f(t)=1-\cos ^{2} t+2 \cos ^{4} t, \quad g(t)=\sin ^{2} t+2 \sin ^{4} t, \quad h(t)=\sin t+\sin ^{3} t .
$$

Computing the functions $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$ of Theorem 1 we get

$$
\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\frac{\pi}{4}\left(7-4 X_{0}+8 Y_{0}\right), \quad \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=-\frac{\pi}{2}\left(X_{0}+3 Y_{0}\right) .
$$

System $\mathcal{F}_{1}=\mathcal{F}_{2}=0$ has the zero $\left(X_{0}^{*}, Y_{0}^{*}\right)=(21 / 20,-7 / 20)$. Since the Jacobian (3) at this zero is $5 \pi^{2} / 2$ the corollary follows.

## 4. Proof of Theorem 2 and Its Corollaries

Proof of Theorem 2. As in the proof of Theorem 1, the second-order differential Equation (1) can be written as the first order differential system (9). Doing the rescaling $(x, y)=(\varepsilon X, \varepsilon Y)$, we obtain the system

$$
\begin{align*}
& \dot{X}=Y \\
& \dot{y}=-X-h(t)+\varepsilon(-g(t) X-f(t) Y-3 X Y)+\varepsilon^{2}\left(-f(t) X^{2}-X^{3}\right) . \tag{11}
\end{align*}
$$

System (11) with $\varepsilon=0$ is the unperturbed system, otherwise it is the perturbed system.
The solution $(X(t), Y(t))$ of the unperturbed system such that $(X(0), Y(0))=\left(X_{0}, Y_{0}\right)$ is

$$
\begin{aligned}
& X(t)=X_{0} \cos t+Y_{0} \sin t-\int_{0}^{t} h(\tau) \sin (t-\tau) \mathrm{d} \tau \\
& Y(t)=-X_{0} \sin t+Y_{0} \cos t-\int_{0}^{t} h(\tau) \cos (t-\tau) \mathrm{d} \tau
\end{aligned}
$$

Note that these periodic orbits have period $2 \pi$. Using the notation introduced in section 2 . We have that $\boldsymbol{x}=(X, Y), \quad \mathbf{z}=\left(X_{0}, Y_{0}\right), \quad F_{0}(\boldsymbol{x}, t)=(Y,-X-h), \quad F_{1}(\boldsymbol{x}, t)=(0,-g(t) X-f(t) Y-3 X Y)$ and $F_{2}(x, t)=\left(0,-f(t) X^{2}-X^{3}\right)$.

The fundamental matrix solution $M_{z}(t)$ is independent of the initial condition $\mathbf{z}$ and it is

$$
M(t)=\left(\begin{array}{cc}
\cos t & \sin t \\
-\sin t & \cos t
\end{array}\right)
$$

We compute the function $\mathcal{F}(\mathbf{z})=\left(\mathcal{F}_{1}\left(X_{0}, Y_{0}\right), \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)\right)$ given in (8), and we get the functions (4) of the statement of Theorem 2.

By Theorem 3, each zero $\left(X_{0}^{*}, Y_{0}^{*}\right)$ of system $\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=0$ satisfying (3), provides a $2 \pi$ periodic solution $(X(t, \varepsilon), Y(t, \varepsilon))$ of system (11) with $\varepsilon \neq 0$ sufficiently small such that

$$
\binom{X(t, \varepsilon)}{Y(t, \varepsilon)}=\binom{X_{0}^{*} \cos t+Y_{0}^{*} \sin t-\int_{0}^{t} h(\tau) \sin (t-\tau) \mathrm{d} \tau}{-X_{0}^{*} \sin t+Y_{0}^{*} \cos t-\int_{0}^{t} h(\tau) \cos (t-\tau) \mathrm{d} \tau}+O(\varepsilon)
$$

Going back through the change of variables for every periodic solution $(X(t, \varepsilon), Y(t, \varepsilon))$ of system (11) with $\varepsilon \neq 0$ sufficiently small, we obtain a $2 \pi$-periodic solution

$$
x(t, \varepsilon)=\varepsilon\left(X_{0}^{*} \cos t+Y_{0}^{*} \sin t-\int_{0}^{t} h(\tau) \sin (t-\tau) \mathrm{d} \tau\right)+O\left(\varepsilon^{2}\right)
$$

of the differential Equation (1) for $\varepsilon \neq 0$ sufficiently small. This completes the proof of Theorem 2.
Proof of Corollary 3. We apply Theorem 2 with

$$
f(t)=\sin (2 t)+\cos (2 t), \quad g(t)=\sin t, \quad h(t)=2 \cos ^{2} t
$$

We compute the functions $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$ of the statement of Theorem 2, and we obtain

$$
\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\frac{\pi}{2}\left(2+X_{0}-4 Y_{0}\right), \quad \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=\frac{\pi}{2}\left(1+8 X_{0}-Y_{0}\right)
$$

System $\mathcal{F}_{1}=\mathcal{F}_{2}=0$ has the solution $\left(X_{0}^{*}, Y_{0}^{*}\right)=(-2 / 31,15 / 31)$. Since the Jacobian (3) is $31 \pi^{2} / 4$, by Theorem 2 we obtain the periodic solution of the statement of the corollary.

Proof of Corollaryc 4. We apply Theorem 2 with

$$
f(t)=\sin t, \quad g(t)=\sin ^{2} t, \quad h(t)=2 \cos (2 t)
$$

We compute the functions $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$ of the statement of Theorem 2, and we obtain

$$
\mathcal{F}_{1}\left(X_{0}, Y_{0}\right)=\frac{3 \pi}{4}\left(8+5 Y_{0}\right), \quad \mathcal{F}_{2}\left(X_{0}, Y_{0}\right)=\frac{11 \pi}{4} X_{0}
$$

System $\mathcal{F}_{1}=\mathcal{F}_{2}=0$ has the solution $\left(X_{0}^{*}, Y_{0}^{*}\right)=(0,-8 / 5)$. Since the Jacobian (3) is $-165 \pi^{2} / 16$, by Theorem 2 we obtain the periodic solution of the statement of the corollary.

## Acknowledgements

The second author is partially supported by a MINECO grant MTM2013-40998-P, an AGAUR grant number 2014SGR568, and the grants FP7-PEOPLE-2012-IRSES 318999 and 316338.

## References

[1] Ince, E.L. (1927) Ordinary Differential Equations. Longmans, London, 1927.
[2] Karasu, A. and Leach, P.G.L. (2009) Nonlocal Symmetries and Integrable Ordinary Differential Equations: $\ddot{x}+3 x \dot{x}+x^{3}=0$ and Its Genralizations. Journal of Mathematical Physics, 50, Article ID: 073509, 17 p.
[3] Chandrasekar, V.K., Senthilvelan, M. and Lakshmanan, M. (2005) Lienard-Type Nonlinear Oscillator. Physical Review E, 72, Article ID: 066203, 8 p.
[4] Chandrasekar, V.K., Senthilvelan, M. and Lakshmanan, M. (2012) A Systematic Method of Finding Linearizing Transformations for Nonlinear Ordinary Differential Equations: I. Scalar Case. Journal of Nonlinear Mathematical Physics, 19, Article ID: 1250012, 21 p.
[5] Chandrasekar, V.K., Senthilvelan, V.K. and Lakshmanan, M. (2012) A Systematic Method of Finding Linearizing Transformations for Nonlinear Ordinary Differential Equations: II. Extension to Coupled ODEs. Journal of Nonlinear Mathematical Physics, 19, Article ID: 1250013, 23 p.
[6] Lakshmanan, M. and Sahadevan, R. (1985) Coupled Quadratic Anharmonic Oscillators, Painlevé Analysis and Integrability. Physical Review A, 31, 861-876. http://dx.doi.org/10.1103/PhysRevA.31.861
[7] Ferreira, C., López, J.L. and Pérez, S. (2014) Ester Convergent and Asymptotic Expansions of Solutions of SecondOrder Differential Equations with a Large Parameter. Analysis and Applications, 12, 523-536. http://dx.doi.org/10.1142/S0219530514500328
[8] Li, J., Luo, J. and Wang, Z. (2014) Periodic Solutions of Second Order Impulsive Differential Equations at Resonance via Variational Approach. Mathematical Modelling and Analysis, 19, 664-675. http://dx.doi.org/10.3846/13926292.2014.980864
[9] Malkin, I.G. (1956) Some Problems of the Theory of Nonlinear Oscillations. Gosudarstv. Izdat. Tehn-Teor. Lit., Moscow. (In Russian).
[10] Roseau, M. (1985) Vibrations non linéaires et théorie de la stabilité. Springer Tracts in Natural Philosophy, Vol. 8, Springer, New York, 1985.
[11] Buica, A., Françoise, J.P. and Llibre, J. (2007) Periodic Solutions of Nonlinear Periodic Differential Systems with a Small Parameter. Communications on Pure and Applied Analysis, 6, 103-111.

