

Angle-Resolved and Resonant Photoemission Study of the Valence Bands of α -La(0001) on W(110)

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Abstract

We report a photoelectron spectroscopic study of the valence bands of double hexagonal-close-packed (dhcp) α -La(0001) films epitaxially grown on W(110) at room temperature. The La 5*d* photoemission cross section in the photon energy region from 20 eV to 130 eV was obtained and the valence-band structure of α -La was determined. Except for 4*f*-related structures, the valence-band structures of dhcp α -La and dhcp β -Ce were found to resemble each other. From the band structure, the crystal structure of the La film was confirmed. No evidence for the existence of a 5*d*-like surface state near the Fermi energy at the $\bar{\Gamma}$ point of the surface Brillouin zone was obtained and a 6*s* band bottom was identified.

Keywords

Lanthanum Thin Film, Electronic Structure, Valence Band, Tungsten Surface, Photoemission, Angle-Resolved Spectroscopy

1. Introduction

The unusual physical properties of the various phases of metallic La have been the subject of many experimental and theoretical investigations [1]-[8]. La exists in three different phases at atmospheric pressure: double hexagonal-close-packed (dhcp) α phase below 609 K, fcc β phase between 609 K and 1138 K, and bcc γ phase between 1138 K and its melting point (1191 K). Although the stable form of La below 609 K is the dhcp α phase, the free energy difference between the α and β phases is so small that the fcc phase can coexist below 609 K, in

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a metastable form, with the α phase.

The ground-state electronic configuration is $[\text{Xe}]4f^05d^16s^2$ for La atom and can be denoted as $[\text{Xe}]4f^0(5d6s)^3$ for La metal considering possible hybridization. Compared with the isoelectronic elements Y and Sc, which are not superconducting under atmospheric pressure, La has a high superconducting transition temperature T_C of 4.88 K (dhcp α) and 6.05 K (fcc β). It is worth noting that the physical properties of La are quite unusual in the β phase which is metastable below ~ 609 K at atmospheric pressure, rather than in the stable α phase. Under pressure, T_C for the β phase rises sharply from ~ 6 K at ambient pressure and saturates at a value of ~ 13 K around 2×10^7 Pa. This rapid rise of T_C with pressure for fcc β -La is the most dramatic among all the elements. In addition, the temperature dependence of the thermal-expansion coefficient is quite anomalous for fcc β -La: it becomes negative at ~ 40 K, and reaches its largest negative value at ~ 18 K.

These remarkable properties have led to speculations and suggestions about their electronic origin that a mechanism involving $4f$ electrons is responsible and the electronic wave functions at the Fermi level contain a significant admixture of $4f$ character. However, bremsstrahlung-isochromat [9] and inverse photoemission [10] experiments indicate that the $4f$ state in La lies ~ 5 eV above the Fermi level (E_F) and is ~ 1 eV wide. Therefore, it is generally considered that the $4f$ hybridization with the occupied states is negligibly small and does not affect the ground-state properties such as lattice dynamics or superconductivity of fcc β -La.

Understanding the properties of La in relation to the electronic structure is, of course, important. Furthermore, in addition to its intrinsic interest, La is of importance as a reference system for the chemically similar $4f$ metals. As the first member of the rare-earth series of elements, the properties of La are frequently compared with those of other rare earths to help interpret properties be associated with the occupied $4f$ levels. The valence-band photoemission of Ce exhibits two peaks attributed to $4f$ emission. To understand better the electronic properties of Ce, it is useful to study the neighboring element La with no $4f$ electrons. Several photoemission experiments have been done to determine the valence-band electronic structure of La [11]-[14], suggesting the existence of a surface state near E_F at the $\bar{\Gamma}$ point of the surface Brillouin zone (SBZ). The interest in surface states arises from the strong influence of the surface electronic structure on physical and chemical properties of metal surfaces. However, the valence-band photoemission investigation for La is still not satisfactory and incomplete, since only part of the bands was studied. It cannot be denied that there has been relatively little theoretical effort given to determine the electronic structure of La, although there are a large number of papers [15]-[22]. This is no doubt related to the scarcity of experimental data available for comparison.

In this paper, we report the results of angle-resolved (ARPES) and resonant valence-band photoemission of epitaxial thin films of La grown on W(110) using synchrotron radiation. Photon-energy ($h\nu$) dependence of the intensities of the photoemission peaks is measured to characterize the energy bands of La and the $5d6s$ valence-band structures of La(0001) film are determined.

2. Experimental

The experiments were performed at the beamline BL-3B of the Photon Factory, Institute of Materials Structure Science, High Energy Accelerator Research Organization (KEK). ARPES spectra were measured at room temperature using a hemispherical electron analyzer with an acceptance of $\pm 1^\circ$. Total instrumental energy resolution was 50 - 100 meV, depending on the photon energy ($h\nu$) in the range of 20 - 130 eV. Each series of spectra was normalized to the relative flux of incident photons. The light incidence angle from the surface normal was fixed to be 45° .

The clean W(110) surface, having sharp (1×1) low-energy electron diffraction (LEED) patterns with low backgrounds, was prepared by repeated heating it to $\sim 2300^\circ\text{C}$ in ultrahigh vacuum. The amounts of impurities were below the detection limit of Auger electron spectroscopy (AES).

La (purity 99.99%) was deposited *in situ* on the W(110) surface by electron-beam evaporation, with the substrate held at room temperature, after a long outgassing of the La source. The base pressure in the experimental chamber was 1.3×10^{-8} Pa, rising to 4.2×10^{-8} Pa during deposition. No traces of surface impurities including oxygen and carbon contaminations were found by AES. The films (up to 15 monolayers) grown under these conditions showed sharp hexagonal LEED patterns corresponding to the formation of well-ordered dhcp α -La(0001) or fcc β -La(111) surfaces. The LEED patterns could not be further improved by sample annealing at 370 - 570 K. The interplanar spacing of dhcp La(0001) (3.04 Å) is almost equal to that of fcc La(111) (3.06 Å), and therefore the distinguishing between dhcp La and fcc La phases is difficult in the LEED experiment. As

stated above, the possibility of coexisting of fcc β -La in a metastable form cannot be denied, but our ARPES band dispersion data are consistent with the dhcp structure rather than the fcc structure. Therefore, throughout this paper, only the stable dhcp α -phase below 609 K is referred. The amounts of deposits were determined from the Auger-peak-intensity ratio $I_{\text{La}(78\text{eV})}/I_{\text{W}(169\text{eV})}$, as described previously [23]. The thicknesses of La films are quoted below in units of equivalent monolayers (ML), representing the number of close-packed La layers of the dhcp La(0001) orientation by assuming uniform thickness.

3. Results and Discussion

Figure 1 shows normal-emission ARPES spectra of α -La(0001)/W(110) in the valence-band region measured at $h\nu = 34$ eV, as a function of the coverage of La (from 0 ML to 5 ML). Each spectrum was measured for a newly deposited La film. The signal from the underlying W(110) substrate will be attenuated due to inelastic scattering of some of the photoelectrons as they traverse through the layer of La metal, depending on the thickness of the layer. The 0-ML spectrum is reduced by 0.06 times considering attenuation of electrons due to inelastic scattering in 1.8-ML La. In **Figure 1**, two La-induced emissions are seen: one labeled *A* is near E_F and the other labeled *B* at ~ 1.3 -eV binding energy (E_B). These emissions grow as increasing La deposition. Besides *A* and *B*, a weak structure is also seen at $E_B \sim 2$ eV. This 2-eV structure was found to decrease to zero with La deposition, and therefore we consider it a trace of underlying W emission.

To obtain information about the character of peaks *A* and *B*, we examine the $h\nu$ dependence of ARPES spectrum of La. **Figure 2** shows a series of normal-emission spectra of ~ 4 -ML α -La(0001)/W(110) in the $h\nu$ range from 20 eV to 130 eV. In addition to peaks *A* and *B*, a weak emission labeled *C* is indistinctly observed at $E_B \sim 3.5$ eV. It is noted that peaks *A* and *B* are positioned at almost constant binding energies of about 0.2 eV and 1.3 eV for all $h\nu$, that is, both peaks hardly show dispersion along the axis of k_{\perp} (wave vector perpendicular to the surface) for the 4-ML thickness film. This suggests that the k_{\perp} is not a good quantum number for La films with thickness less than 4 ML. In the case of Ni films on W(110), it has been reported that the wave vector k_{\perp} becomes a good quantum number and the k_{\perp} dispersion of Ni starts to develop only when film thickness exceeds 5 ML [24]. Although the energy positions of peaks *A* and *B* are independent of $h\nu$, the intensities depend strongly on $h\nu$. In order to gain some insight into character of the peaks, we examine $h\nu$ dependence of the intensities of the peaks *A* and *B*.

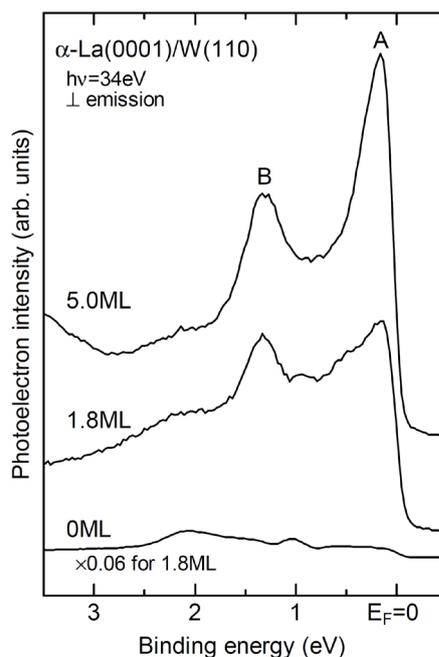


Figure 1. Normal-emission valence-band spectra of α -La(0001)/W(110), as a function of the coverage of La, measured at $h\nu = 34$ eV. The 0-ML spectrum is a spectrum of W(110) substrate.

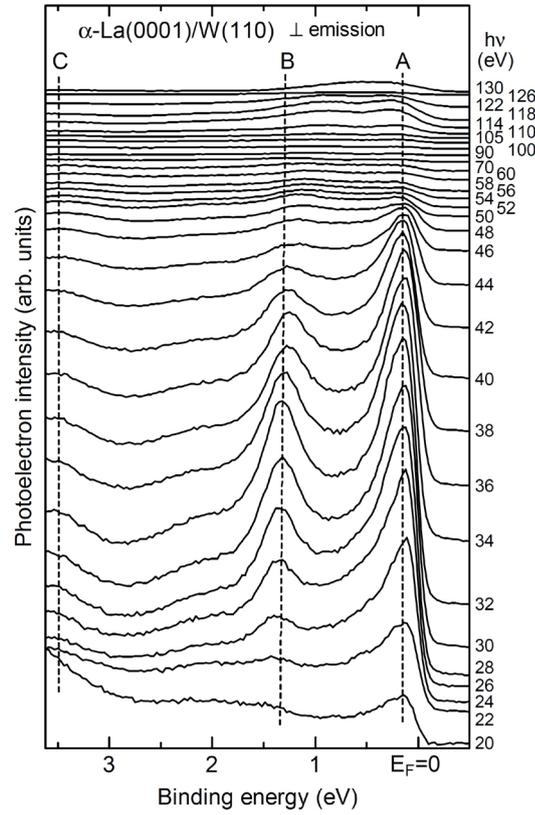


Figure 2. Normal-emission valence-band spectra of ~ 4 -ML α -La(0001)/W(110) as a function of $h\nu$ from 20 eV to 130 eV.

Figure 3 shows $h\nu$ dependence of the heights (excitation spectra) of peaks A and B in the $h\nu$ region from 20 eV to 130 eV, obtained by averaging the data of **Figure 2** and other two sets of similar series. The peak height was measured after properly subtracting a linear background from each spectrum. The heights of both peaks A and B rise up with increasing $h\nu$ to reach a maximum at $h\nu \sim 32$ eV, thereafter falls monotonically to approach zero until ~ 100 eV, and then, when $h\nu$ exceeds ~ 100 eV, again increases to reach a second maximum centered at $h\nu \sim 116$ eV. Now, we compare the excitation spectrum of La with that of Ce. The excitation spectrum of Ce has its own peculiar structure centered at $h\nu \sim 42$ eV [23]. This 42-eV structure is due to the Ce 4*f* electron photoemission. It should be noted that the present excitation spectrum of La does not have such a structure originated in 4*f* electrons, supporting the previous conclusion that the 4*f* hybridization with the occupied states is negligibly small in La. However, in the lower $h\nu$ region up to ~ 100 eV, the excitation spectra for peaks A and B resemble that for Ce 5*d* peak [23]. For $20 \text{ eV} < h\nu < 40 \text{ eV}$, the Hartree-Fock-Slater atomic photoionization cross section for La 6*s* is 10 - 30 times as small as that for La 5*d* [25]. Thus, the profiles of the excitation spectra for peaks A and B indicate that their initial states are of La 5*d* character. The 116-eV structure is due to a two-step process: the La 4*d* resonating giant photoexcitation $h\nu + 4d^{10}(5d6s)^3 \rightarrow 4d^9(5d6s)^3 4f^1(^1P_1)$ and the subsequent autoionization decay $4d^9(5d6s)^3 4f^1 \rightarrow 4d^{10}(5d6s)^2 + e^-$ (e^- denotes the ejected photoelectron), *i.e.*, so-called 4*d* resonant photoemission of La [26]-[29].

Figure 4 shows off-normal emission spectra of 4.0-6.0-ML α -La(0001)/W(110) as a function of the emission angle θ_e along the $[10\bar{1}0]$ ($\bar{\Gamma}$ - \bar{M}) azimuth. Three major peaks in the spectra are marked by tick marks and are labeled A-C as above. From the spectra in **Figure 4**, we determined the E -versus- k_{\parallel} dispersion [$E(k_{\parallel})$] along the $\bar{\Gamma}$ - \bar{M} direction of SBZ. The results are summarized in **Figure 5**. The four sets of band structures derived from **Figures 4(a)-(d)** are well in agreement.

As seen in **Figure 5**, the obtained band structure $E(k_{\parallel})$ is rather complicated. Band A is split into two bands at $k_{\parallel} \sim 0.45 \text{ \AA}^{-1}$: the upper and lower bands are labeled A1 and A2, respectively. Peak A1 is weak and shows substantially no k_{\parallel} dispersion. As k_{\parallel} increases up to 0.94 \AA^{-1} , band A2 disperses downward by ~ 0.7 eV, to reach the

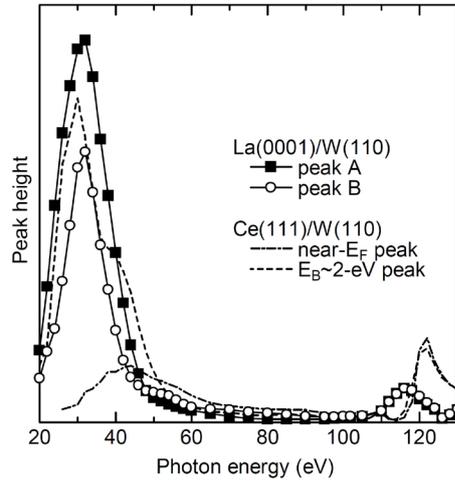


Figure 3. $h\nu$ dependence of the heights (excitation spectra) of the peaks A and B, obtained by averaging the data of Figure 2 and other two sets of similar series. The excitation spectra of near- E_F $4f^1$ -final-state peak and $E_B \sim 2$ -eV $4f^0$ -final-state peak for γ -Ce(111) (from [23]) are shown to be compared with those of La. The Ce excitation spectra are normalized to the average of the 32-eV maxima of peaks A and B.

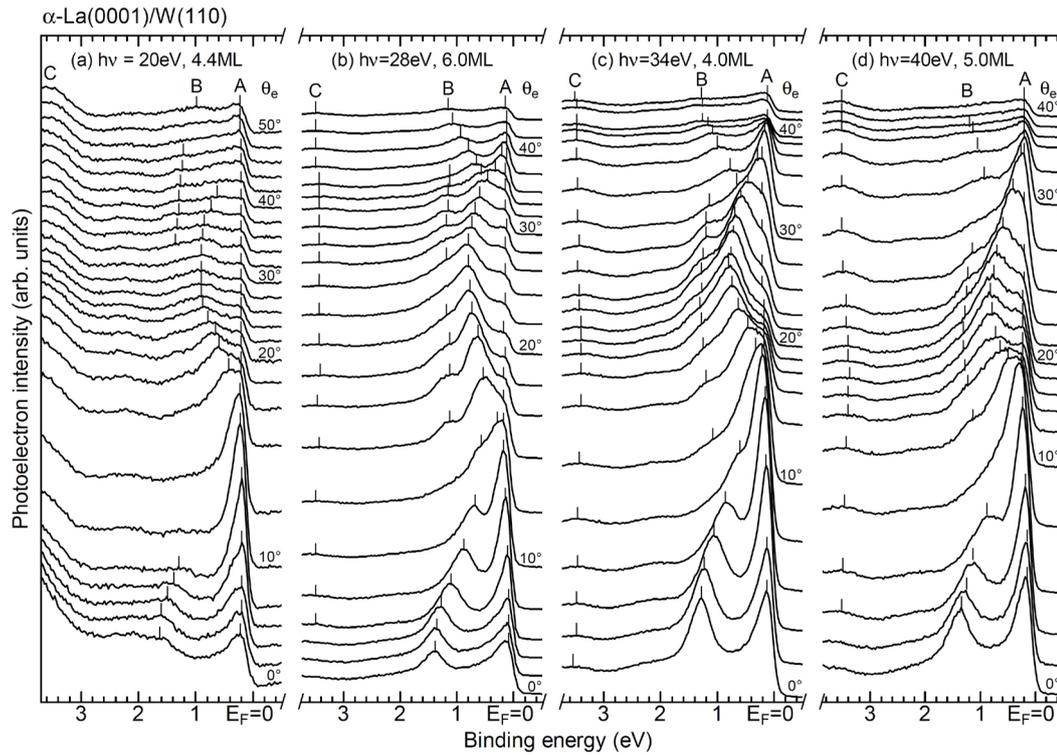


Figure 4. Off-normal emission valence-band spectra of 4.0 - 6.0-ML α -La(0001)/W(110) measured at (a) $h\nu = 20$ eV; (b) 28 eV; (c) 34 eV; and (d) 40 eV, with the emission angle θ_e changed by 2° along the $[10\bar{1}0]$ ($\bar{\Gamma}-\bar{M}$) azimuth. The emission angle θ_e is measured from the surface normal in the photoelectron collection plane.

bottoms at $E_B \sim 0.85$ eV. Thereafter, band A2 disperses upward and bands A1 and A2 are merged into one at $k_{//} \sim 1.5 \text{ \AA}^{-1}$. Band B disperses upward as $k_{//}$ increases up to $\sim 0.35 \text{ \AA}^{-1}$ and then is split into two bands: upper band B1 and lower band B2. Thereafter, band B1 still disperses upward until $\sim 0.5 \text{ \AA}^{-1}$ and then downward: bands B1 and A2 unite to one labeled A2 + B1 between ~ 0.5 and $\sim 1.35 \text{ \AA}^{-1}$. Band B2 disperses downward to reach the rather flat bottom at $E_B \sim 1.27$ eV between ~ 0.6 and $\sim 1.3 \text{ \AA}^{-1}$. Thus, the bands B1 and B2 show very complicated dis-

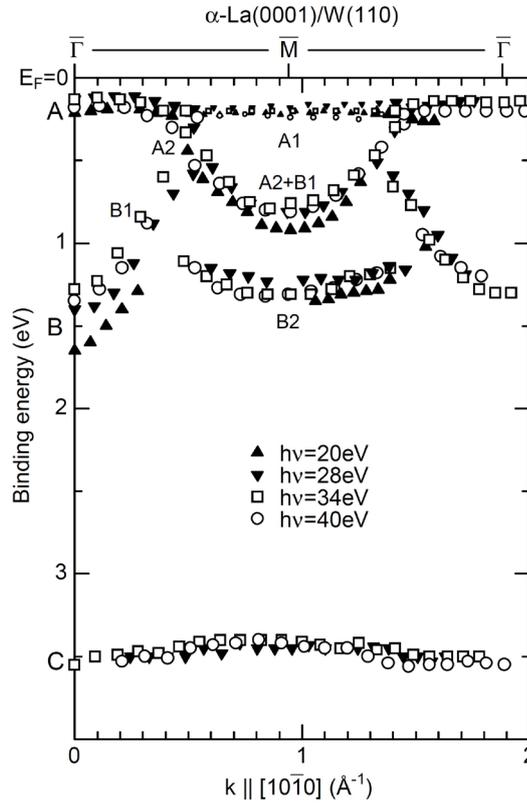


Figure 5. Measured valence-band dispersions $E(k_{||})$ of peaks A (A1, A2), B (B1, B2), and C for 4.0 - 6.0-ML α -La(0001)/W(110) in the $[10\bar{1}0]$ ($\bar{\Gamma} - \bar{M}$) azimuth. Solid upward triangles, solid downward triangles, open squares, and open circles are data points for $h\nu = 20$ eV, 28 eV, 34 eV, and 40 eV derived from **Figures 4(a)-(d)**, respectively.

persions, up and down. Peak C is weak in intensity (see **Figure 4**) and shows no dispersion or, if any, a very small upward dispersion by ~ 0.1 eV with the increase of $k_{||}$ from 0 to 0.94 \AA^{-1} (peak C in **Figure 4(d)**).

Here, we note that the band structure is symmetrical with respect to the point of $k_{||} = 0.94 \text{ \AA}^{-1}$. If this turning $k_{||}$ -point is \bar{M} in the fcc(111) SBZ, the corresponding lattice constant is estimated to be 5.46 \AA since $k_{\bar{\Gamma}\bar{M}} = 2(2/3)^{1/2}\pi/a$. This a value is larger than that for fcc β -La ($a_{\text{fcc}} = 5.3 \text{ \AA}$) by 3.0%. If this turning $k_{||}$ -point is \bar{M} in the dhcp(0001) SBZ, the corresponding lattice constant is estimated to be 3.86 \AA since $k_{\bar{\Gamma}\bar{M}} = 2(1/3)^{1/2}\pi/a$. This a value is larger than that for dhcp α -La ($a_{\text{dhcp}} = 3.77 \text{ \AA}$) by 2.4%. This comparison may indicate that the dhcp phase is slightly favorable, but the difference between 2.4 and 3.0 is too small to distinguish definitely between fcc La and dhcp La phases. Looking on fcc(111) as dhcp(0001), $a_{\text{fcc}}/a_{\text{dhcp}} = \sqrt{2}$. Note that $5.3 \text{ \AA}/3.77 \text{ \AA} = 1.41$. Therefore, the lattice-constant consideration alone cannot determine the crystal structure of the present La film.

It is the electronic energy dispersion relation $E(k)$ which determines the phase of La film. A great number of band-structure calculations of La have been reported so far, but most of them were for fcc β -phase, and for dhcp α -phase one and only one has been reported by Jarlborg *et al.* [22]. All the calculated results for fcc β -phase are similar. We compare the experimental band structure (**Figure 5**) with calculated one for fcc La (**Figure 5** in [22]) or dhcp La (**Figure 6** in [22]), assuming that most of the observed bands are of bulk nature.

In the present case, as stated above, k_{\perp} is not a good quantum number, and it would be expected that $k_{||}$ -resolved structures in the one-dimensional density of bulk states ($k_{||}$ -resolved DOS) are observed by tuning θ_e . Singularities in the density of states (*i.e.*, high occupied density of states) are prominent in the $k_{||}$ -resolved DOS spectrum and therefore the $k_{||}$ -resolved DOS spectrum may reflect, not rigorously but approximately, the $E(k_{||})$ dispersion [30] [31]. The $[10\bar{1}0]$ ($\bar{\Gamma} - \bar{M}$) azimuth of dhcp(0001) corresponds to the $[11\bar{2}]$ ($\bar{\Gamma} - \bar{M}$) azimuth of fcc(111). We first examine the band structure of fcc β -La. A plane defined by the fcc(111) surface normal and the $[11\bar{2}]$ direction is the $\Gamma K L U X$ ($\bar{1}10$) mirror plane. The $\Gamma - K$, $\Gamma - L$, $\Gamma - X$, and $L - U$ axes are in the $(\bar{1}10)$ plane, but, except for $L - U$, not parallel to the $[11\bar{2}]$ direction. It is found that the calculated band structures along $\Gamma - K$, $\Gamma - L$,

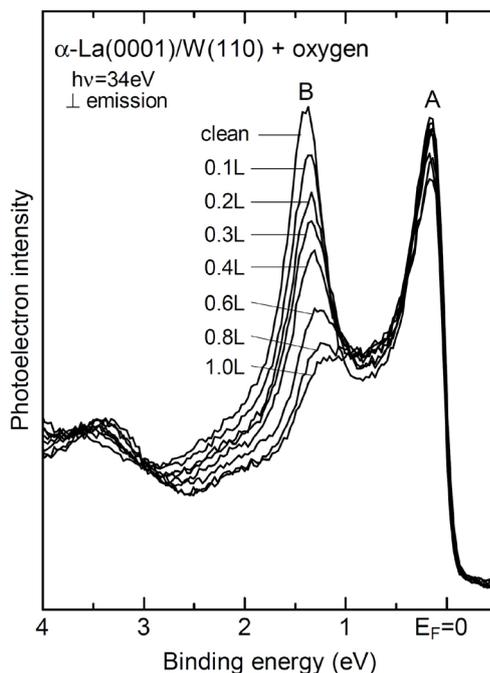


Figure 6. Comparison of the $h\nu = 34$ -eV 4.0-ML α -La(0001)/W(110) spectra before and after O_2 exposure. The exposure increases from 0.1 L to 1.0 L by 0.1 L or 0.2 L.

Γ -X, and L -U can explain experimental band A2 + B1, but are inconsistent with other bands. More important, in sharp contrast to observations, the bulk bands of fcc β -La are not symmetrical about \bar{M} .

In the case of dhcp α phase, the $(1\bar{2}10)$ mirror plane, defined by the (0001) surface normal and the $[10\bar{1}0]$ ($\bar{\Gamma}$ - \bar{M}) direction, contains the Γ - \bar{M} and A - L axes, which are parallel to $[10\bar{1}0]$ and furthermore have the same one-dimensional periodicity as the $\bar{\Gamma}$ - \bar{M} axis. That is, the bulk bands along Γ - \bar{M} and A - L are symmetrical about \bar{M} in agreement with the observations. Comparing the experimental band structure (Figure 5) with the calculated one for dhcp La (Figure 6 in [22]), we realize fairly good correspondence between the two, with regard to position, dispersion, and periodicity. The intensity of peak A, which is rather weak at and around $\bar{\Gamma}$, increases with increasing k_{\parallel} up to $\sim 0.45 \text{ \AA}^{-1}$ where peak A is split into two, A1 and A2. The reverse of this is realized from the A1-A2 merging point ($k_{\parallel} \sim 1.5 \text{ \AA}^{-1}$) to $\bar{\Gamma}_{2nd}$ point in second zone ($k_{\parallel} \sim 1.9 \text{ \AA}^{-1}$). Therefore, a band (or bands) should disperse down across E_F near $k_{\parallel} \sim 0.45 \text{ \AA}^{-1}$ and back up near $k_{\parallel} \sim 1.5 \text{ \AA}^{-1}$, in qualitative agreement with the calculated band structure of dhcp La [22]. Like this, a band corresponding to observed band A2 + B1, B1, or B2 (except for A1) is found in the calculated bands along Γ - \bar{M} or A - L . A band corresponding to band A1 is missing: this low-intense Fermi-edge transition A1 can be ascribed to a kind of DOS transition visible due to the life-time broadening of the final states [32]. Thus, as for the bands from E_F to $E_B \sim 2$ eV, which are $5d$ -like according to [22], and except for band A1, our ARPES band dispersion data are consistent with the dhcp structure rather than the fcc structure.

The band structure calculation for dhcp La predicts a $6s$ -like band along Γ - \bar{M} (from $E_B \sim 2$ eV at M to ~ 3.4 eV at Γ) [22], but is not observed in ARPES owing to its low photoionization cross section as compared with that of $5d$. For $20 \text{ eV} < h\nu < 40 \text{ eV}$, the Hartree-Fock-Slater atomic photoionization cross section for La $6s$ is 10 - 30 times as small as that for La $5d$ [25]. Note that a weak feature C is observed at $E_B \sim 3.5$ eV. Band structure calculations of dhcp La predict that the bottom of $6s$ band is at $E_B \sim 3.4$ eV [22], in agreement with our observations. Singularities in the density of states (*i.e.*, high occupied density of states) may lead to low-intensity structures in the spectra.

Rare earth elements are characterized by their partially filled $4f$ -shell. In the case of Ce with the ground-state electronic configuration of $[\text{Xe}]4f^1(5d6s)^3$, its valence-band photoelectron spectrum reveals a characteristic double-peaked structure: one near E_F and one at $E_B \sim 2$ eV [23]. It is by now generally accepted that the peak near E_F is assigned to a well-screened $4f^1$ final state, with the $4f$ hole filled by a valence electron, and the other at $E_B \sim 2$ eV is assigned to a poorly screened $4f^0$ final state, with valence electrons ($5d$) partially screening the

hole. Contrary to Ce, in the present case of La, the observed spectra do not show such a double-peaked structure and, e.g., a poorly-screened $4f^0$ final-state feature is obviously missing. This result reconfirms that the $4f$ hybridization with the occupied states is negligibly small in La. Furthermore, it is found that, except for $4f$ -related structures, the valence-band structure of dhcp α -La resembles closely that of dhcp β -Ce (or fcc γ -Ce) [23].

Finally, we want to examine band A and the possibility of its relation to a surface state. Recent development in valence-band photoemission studies of rare-earth metals revealed that a d -like surface state exists just below E_F around the center ($\bar{\Gamma}$) of the SBZ for the close-packed surfaces of almost all rare-earth metals including La [12] [14]. In fact, according to [22], there is a large gap between $E_B = -0.5$ eV and $+1.7$ eV at and around $\bar{\Gamma}$ for dhcp La(0001) and, for example, according to the band calculation of an 11 layer slab for dhcp Ce, a surface state exists very close to the top of a band gap along $\bar{\Gamma}-\bar{M}$ in the SBZ (see Figure 2 in [33]). Thus, there is a possibility that band $A - A2 - A2 + B1$ is a surface state. In order to get some insight into the character of peak A and also peak B , we examine the influence of surface contamination on the peaks. The sensitivity to contamination is generally used as one test for surface states, as is the case of La [12]. Note that a surface state is not necessarily sensitive to the adsorption of every kind of foreign atom and some bulk state also show high sensitivity to foreign atoms [34]. Therefore, we should take notice that sensitivity to the presence of foreign atoms on the surface is not a universally correct rule for the identification of a surface state, but only a guideline. Figure 6 displays comparison of the $h\nu = 34$ -eV 4.0-ML α -La(0001)/W(110) spectra before and after O_2 exposure. The exposure increases from 0.1 L to 1.0 L (1 L = 1 langmuir = 10^{-6} Torr-sec). It is found that peak B is rapidly quenched by exposing the La(0001) surface to O_2 up to 1 L, whereas peak A is almost unchanged. These findings may suggest that the peak A is ascribed to the bulk electronic structure of La, while the peak B apparently to the surface electronic structure. However, the energy position of peak B is strongly $h\nu$ -dependent (that of peak A is weakly $h\nu$ -dependent), as clearly seen in Figure 5. Although the possibility of the existence of a surface state on dhcp La(0001) cannot be denied completely, both peaks A and B are considered to be originated from the bulk-like states.

4. Conclusion

The electronic structure of well-ordered and atomically clean dhcp α -La(0001) epitaxially grown on a W(110) surface has been studied by means of photoelectron spectroscopy. The photoemission cross sections of the La $5d$ states were clarified in the $h\nu$ region from 20 eV to 130 eV, confirming character of the photoemission peaks. The energy-band structure of the dhcp α -La(0001) film was determined along the $[10\bar{1}0]$ ($\bar{\Gamma}-\bar{M}$) azimuth direction and was found to be consistent with the theoretical calculation [22]. In addition, by comparison between La and Ce, except for $4f$ -related structures, the valence-band structures of dhcp α -La and dhcp β -Ce were found to resemble closely each other. Structural information of the film was obtained from the experimentally-determined dispersions: the experimental dispersions reveal a lattice constant of 3.86 Å, which is larger than the lattice constant of bulk dhcp α -La only by 2.4%. Contrary to expectations, no undeniable evidence for the existence of a $5d$ -like surface state near E_F at the $\bar{\Gamma}$ point of SBZ was obtained. The $6s$ band bottom was identified. Our data demonstrated importance of using well-ordered and atomically clean La(0001) epitaxial films grown on a well-ordered and atomically clean W(110) surface for carrying out La photoemission experiments.

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