

Numerical Solution of MHD Boundary Layer Flow of Non-Newtonian Casson Fluid on a Moving Wedge with Heat and Mass Transfer and Induced Magnetic Field

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Abstract

The paper investigates the numerical solution of the magnetohydrodynamics (MHD) boundary layer flow of non-Newtonian Casson fluid on a moving wedge with heat and mass transfer. The effects of thermal diffusion and diffusion thermo with induced magnetic field are taken in consideration. The governing partial differential equations are transformed into nonlinear ordinary differential equations by applying the similarity transformation and solved numerically by using finite difference method (FDM). The effects of various governing parameters, on the velocity, temperature and concentration are displayed through graphs and discussed numerically. In order to verify the accuracy of the present results, we have compared these results with the analytical solutions by using the differential transform method (DTM). It is observed that this approximate numerical solution is in good agreement with the analytical solution. Furthermore, comparisons of the present results with previously published work show that the present results have high accuracy.

Keywords

Casson Fluid, Induced Magnetic Field, Boundary Layer, Moving Wedge, Radiation, Soret and Dufour Effects

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1. Introduction

The induced magnetic field is arisen due to a strong magnetic field. The problems with induced magnetic field play an important role in a number of industrial applications such as fiber or granular insulation, liquid-metals, electrolytes, ionized gases as well as the geothermal systems. The study of the boundary layer flow under the influence of a magnetic field with the induced magnetic field was considered by a few authors. For example, Raptis and Perdakis [1] studied the MHD free convection boundary layer flow past an infinite vertical porous plate. Very little attention has been paid to the boundary layer flow and heat transfer taking the effect of the induced magnetic field and thermal radiation into consideration. For example, Kumari *et al.* [2] considered the MHD flow and heat transfer over a stretching surface by considering the effect of the induced magnetic field. On the other hand, Takhar *et al.* [3] investigated the unsteady free convection flow at the stagnation point in the presence of a magnetic field. Recently Ali *et al.* [4] studied MHD stagnation-point flow and heat transfer towards stretching sheet with induced magnetic field. Khamisah Jafar *et al.* [5] studied viscous dissipation and radiation effects in MHD stagnation point flow towards a stretching sheet with induced magnetic field. The thermal reversal in MHD stagnation point flow towards a stretching sheet with induced magnetic field and viscous dissipation effects have been investigated by Kashif Ali and Muhammad Ashraf [6]. The MHD mixed convection boundary layer flow toward a stagnation point on a vertical surface with induced magnetic field was studied by Ali *et al.* [7]. Khamisah Jafar *et al.* [8] discussed the MHD boundary layer flow due to a moving wedge in a parallel stream with the induced magnetic field.

The analysis of non-Newtonian fluids has many applications in industrials such as ground water hydrology, petroleum reservoir, nuclear waste disposal, geothermal energy production, transpiration cooling, design of solid matrix heat exchange and paced bed chemical catalytic reactors. Also, there are some models of non-Newtonian fluids such as Casson model which concerned one of the most important applications of theoretical fluid mechanics to problems arising in physiology, mainly in describing the flow of the blood. Casson [9] proposed a model to describe the flow curves of suspensions of pigments in lithographic varnishes used for the preparation of printing inks. Eldabe *et al.* studied some problems of non-Newtonian fluids such as ([10]-[14]).

In the present work, we extend and generalize the work of [8] to include the boundary layer motion of non-Newtonian fluid obeying Casson model with heat and mass transfer in the presence of radiation and mixed diffusion. The main objective of the present work is to obtain the numerical solutions of the MHD boundary layer flow of Casson fluid on a moving wedge with heat and mass transfer. The effects of thermal diffusion and diffusion thermo with induced magnetic field are taken into consideration. The governing equations of motion with the boundary conditions are solved numerically by using (FDM) and analytically by using (DTM). The solutions are obtained as a function of the physical parameters of the problem; then the effects of these various parameters of the problem on these solutions have been computed and discussed in detail and illustrated through some figures. The special cases of our problem are compared with the previous published works. This showed that there is a good agreement between them.

2. Mathematical Formulation

Let us consider a steady, laminar, hydromagnetic coupled heat and mass transfer by mixed convection flow on a wedge plate in a parallel free stream with a variable induced magnetic field applied parallel to the wedge walls outside the boundary layer. The equations governing the steady motion of an incompressible non-Newtonian electrically conducting fluid in presence of a magnetic field are:

$$\nabla \cdot V = 0, \quad \nabla \cdot H = 0 \quad (1)$$

$$(V \cdot \nabla)V - \frac{\mu_0}{4\pi\rho}(H \cdot \nabla)H = -\frac{1}{\rho}\nabla P + \frac{1}{\rho}\nabla \cdot \tau, \quad (2)$$

$$\nabla \times (V \times H) + \mu_e \nabla^2 H = 0, \quad (3)$$

$$\rho c_p (V \cdot \nabla)T = k \nabla^2 T + \tau \cdot (\nabla V) - \frac{\partial q_r}{\partial y} + \frac{\rho D_m k_T}{c_s} \nabla^2 C \quad (4)$$

$$(V \cdot \nabla)C = D_m \nabla^2 C + \frac{D_m k_T}{T_m} \nabla^2 T \quad (5)$$

where V is the fluid velocity vector, H is the induced magnetic field vector, $P = p + \frac{\mu |H|^2}{8\pi}$ is the magnetohydrodynamic pressure, p is the fluid pressure, μ_0 is the magnetic permeability, τ is the stress tensor, σ is the electrical conductivity, T is the temperature, k is the thermal conductivity, c_p is the specific heat at constant pressure, ρ is the density of the fluid, $\mu_e = \frac{1}{4\pi\sigma}$ is the magnetic diffusivity, q_r is the radiative heat flux, C is the species concentration, c_s is the concentration susceptibility, k_T is thermal diffusion ratio, T_m is mean fluid temperature and D_m is the coefficient of chemical molecular diffusivity.

The rheological equation of state for an isotropic flow of a Casson fluid can be expressed as [9]:

$$\tau_{ij} = \begin{cases} 2 \left(\mu_B + \frac{p_y}{\sqrt{2\pi}} \right) e_{ij}, & \pi > \pi_c \\ 2 \left(\mu_B + \frac{p_y}{\sqrt{2\pi}} \right) e_{ij}, & \pi < \pi_c \end{cases} \quad (6)$$

In the above equation $\pi = e_{ij}e_{ij}$ and e_{ij} denotes the $(i, j)^{th}$ component of the deformation rate, π is the product of the component of deformation rate with itself, π_c is the a critical value of this product based on the non-Newtonian model, μ_B is the plastic dynamic viscosity of the non-Newtonian fluid and p_y is the yield stress of the fluid.

We take the Cartesian coordinates x measured along the surface of the wedge and y normal to it, respectively. If (u, v) and (H_1, H_2) are the velocity and magnetic components in (x, y) directions, respectively, subject to the boundary layer approximations, Equations (5) and (6) for the problem under consideration can be reduced to:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (7)$$

$$\frac{\partial H_1}{\partial x} + \frac{\partial H_2}{\partial y} = 0, \quad (8)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - \frac{\mu_0}{4\pi\rho} \left(H_1 \frac{\partial H_1}{\partial x} + H_2 \frac{\partial H_1}{\partial y} \right) = U_e \frac{dU_e}{dx} - \frac{\mu_0 H_e}{4\pi\rho} \frac{dH_e}{dx} + v \left(1 + \frac{1}{\beta} \right) \frac{\partial^2 u}{\partial y^2}, \quad (9)$$

$$u \frac{\partial H_1}{\partial x} + v \frac{\partial H_1}{\partial y} - H_1 \frac{\partial u}{\partial x} - H_2 \frac{\partial u}{\partial y} = \mu_e \frac{\partial^2 H_1}{\partial y^2} \quad (10)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2} + \frac{v}{c_p} \left(1 + \frac{1}{\beta} \right) \left(\frac{\partial u}{\partial y} \right)^2 - \frac{1}{\rho c_p} \left(\frac{\partial q_r}{\partial y} \right) + \frac{D_m k_T}{c_s c_p} \left(\frac{\partial^2 C}{\partial y^2} \right) \quad (11)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_m \frac{\partial^2 C}{\partial y^2} + \frac{D_m k_T}{T_m} \left(\frac{\partial^2 T}{\partial y^2} \right) \quad (12)$$

The boundary conditions are as follows:

$$\begin{aligned} v = 0, \quad u = u_w(x) = U_w x^m, \quad H_1 = H_2 = 0, \quad T = T_w, \quad C = C_w, \quad \text{at } y = 0 \\ u = U_e(x) = U_\infty x^m, \quad H_1 = H_e(x) = H_0 x^m, \quad T = T_\infty, \quad C = C_\infty \quad \text{as } y \rightarrow \infty \end{aligned} \quad (13)$$

where $\beta = \mu_B \sqrt{2\pi c_p / p_y}$ is the non-Newtonian Casson parameter, ν is the kinematic viscosity, $U_e(x)$ and $H_e(x)$ are the x-velocity and magnetic field at the edge of the boundary layer, respectively and $\alpha = k / \rho c_p$ is the thermal diffusivity.

We assume here that $U_e(x) = U_\infty x^m$ and $H_e(x) = H_0 x^m$, U_∞ is the constant velocity at the outer edge of the boundary layer and H_0 is the value of $H_e(x)$ at $x = 0$. Further, m is also a constant, which varies in the

range $0 \leq m \leq 1$, U_w is a positive or a negative constant [8].

By using the Rosseland approximation, the radiative heat flux q_r is given by

$$q_r = -\frac{4\sigma^*}{3k^*} \frac{\partial T^4}{\partial y} \quad (14)$$

where σ^* is the Stefan-Boltzmann constant and k^* is the mean absorption coefficient. It should be noted that by using the Rosseland approximation, the present analysis is limited to optically thick fluids. If temperature differences within the flow are sufficiently small, then Equation (14) can be linearized by expanding T^4 into the Taylor series about T_∞ , which after neglecting higher order terms takes the form

$$T^4 \approx 4T_\infty^3 T - 3T_\infty^4 \quad (15)$$

In view of Equations (14) and (15), Equation (11) reduces to

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \left(1 + \frac{16\sigma^* T_\infty^3}{3kk^*} \right) \frac{\partial^2 T}{\partial y^2} + \frac{\nu}{c_p} \left(1 + \frac{1}{\beta} \right) \left(\frac{\partial u}{\partial y} \right)^2 + \frac{D_m k_T}{c_s c_p} \left(\frac{\partial^2 C}{\partial y^2} \right) \quad (16)$$

3. Method of Solutions

To solve system of Equations (7)-(13), we will consider the following similarity transformations:

$$\begin{aligned} \eta &= \sqrt{\frac{(m+1)U_e(x)}{2\nu x}} y, \quad \psi = \sqrt{\frac{2\nu x U_e(x)}{(m+1)}} f(\eta), \quad \phi = H_0 \sqrt{\frac{2\nu x x^m}{(m+1)U_\infty}} g(\eta), \\ \theta &= \left(\frac{T - T_\infty}{T_w - T_\infty} \right), \quad \varnothing = \left(\frac{C - C_\infty}{C_w - C_\infty} \right) \end{aligned} \quad (17)$$

where η is the independent similarity variable, $f(\eta)$ the dimensionless stream function, $g(\eta)$ the dimensionless Induced magnetic field, $\theta(\eta)$ the dimensionless temperature and $\phi(\eta)$ the dimensionless concentration. Further, ψ is the stream function which is defined in the usual way as $u = \partial\psi/\partial x$ and $v = \partial\psi/\partial y$ and ϕ is the magnetic stream function which is defined in the usual way as $H_1 = \partial\phi/\partial x$ and $H_2 = \partial\phi/\partial y$.

By using Equation (17) in the Equations (9)-(12) and Equation (16), we get the following ordinary differential equations in dimensionless form

$$\left(1 + \frac{1}{\beta} \right) f''' + ff'' + \gamma(1 - f'^2) - M[gg'' + \gamma(1 - g'^2)] = 0 \quad (18)$$

$$\lambda g''' + fg'' - f''g = 0 \quad (19)$$

$$\left(1 + \frac{4}{3} R_a \right) \theta'' + P_r f \theta' + P_r E \left(1 + \frac{1}{\beta} \right) f''^2 + P_r D_u \varnothing'' = 0 \quad (20)$$

$$\varnothing'' + S_c f \varnothing' + S_c S_0 \theta'' = 0 \quad (21)$$

where $\gamma = 2m/(m+1)$ is the wedge parameter, $M = \mu H_0^2 / 4\pi \rho U_\infty^2$ is the magnetic parameter, $\lambda = 1/4\pi\nu\sigma$ is the reciprocal magnetic Prandtl number, $E = U_e^2(x)/c_p(T_w - T_\infty)$ is the Eckert number, $P_r = \nu/\alpha$ is the Prandtl number, $R_a = 4\sigma T_\infty^3 / kk^*$ is the radiation parameter, $S_c = \nu/D_m$ is the Schmidt number, $D_u = D_m k_T (C_w - C_\infty) / \nu c_s c_p (T_w - T_\infty)$ is the Dufour number and $S_0 = D_m k_T (T_w - T_\infty) / \nu T_m (C_w - C_\infty)$ is the Soret number.

Also, the subjected boundary conditions Equation (13) will now take the form:

$$\begin{aligned} f(0) &= 0, \quad f'(0) = \delta, \quad g(0) = 0, \quad g'(0) = 0, \quad \theta(0) = 1, \quad \varnothing(0) = 1, \\ f'(\infty) &\rightarrow 1, \quad g'(\infty) \rightarrow 1, \quad \theta(\infty) \rightarrow 0, \quad \varnothing(\infty) \rightarrow 0 \end{aligned} \quad (22)$$

where $\delta = U_w/U_\infty$ is the moving parameter.

3.1. The Differential Transform Method (DTM)

The differential transformation of an analytical function $f(t)$ for one variable is defined as [15].

$$F(k) = \frac{1}{k!} \left[\frac{d^k f(t)}{dt^k} \right]_{t=t_0} \quad (23)$$

where $f(t)$ is the original function and $F(k)$ is the transformed function. The differential inverse transformation of $F(k)$ is defined as:

$$f(t) = \sum_{k=0}^{\infty} F(k)(t-t_0)^k \quad (24)$$

Combining Equations (23) and (24), we obtain

$$f(t) = \sum_{k=0}^{\infty} \frac{(t-t_0)^k}{k!} \left[\frac{d^k f(t)}{dt^k} \right]_{t=t_0} \quad (25)$$

From Equations (23)-(25), it can be seen that the differential transformation method is derived from Taylor's series expansion, but the method does not calculate the derivatives representatively. However, the relative derivatives are calculated by an iterative way which is described by the transformed equations of the original function. For implementation purposes, the function $f(t)$ is expressed by a finite series and Equation (24) can be written as

$$f(t) \approx \sum_{k=0}^N F(k)(t-t_0)^k \quad (26)$$

By Equation (23), the following theorems can be deduced:

Theorem 3.1.1 If $f(t) = u(t) \pm v(t)$ then $F(k) = U(k) \pm V(k)$.

Theorem 3.1.2 If $f(t) = \alpha u(t)$ then $F(k) = \alpha U(k)$.

Theorem 3.1.3 If $f(t) = t^m$ then $F(k) = \delta(k-m) = \begin{cases} 1 & k=m \\ 0 & \text{otherwise} \end{cases}$.

Theorem 3.1.4 If $f(t) = \frac{du(t)}{dt}$ then $F(k) = (k+1)U(k+1)$.

Theorem 3.1.5 If $f(t) = \frac{d^n u(t)}{dt^n}$ then $F(k) = \frac{(k+n)!}{k!} U(k+n)$.

Theorem 3.1.6 If $f(t) = u(t) \cdot v(t)$ then $F(k) = \sum_{r=0}^k U(r)V(k-r)$.

Theorem 3.1.7 If $f(t) = \frac{du(t)}{dt} \frac{dv(t)}{dt}$ then $F(k) = \sum_{r=0}^k (r+1)(k-r+1)U(r+1)V(k-r+1)$.

3.2. Basic Concepts of the Multi-Step Differential Transform Method (MDTM)

When the DTM is used for solving differential equations with the boundary condition at infinity or problems that have highly non-linear behavior, the obtained results were found to be incorrect (when the boundary-layer variable go to infinity, the obtained series solutions are divergent). Besides that, power series are not useful for large values of the independent variable.

To overcome this shortcoming, the MDTM that has been developed for the analytical solution of the differential equations is presented in this section. For this purpose, the following non-linear initial-value problem is considered:

$$u(t, f, f', \dots, f^{(p)}) = 0 \quad (27)$$

subject to the initial conditions $f^{(k)}(0) = c_k$, for $k = 0, 1, 2, \dots, p-1$.

Let $[0, T]$ be the interval over which we want to find the solution of the initial-value problem (27). In actual applications of the DTM, the approximate solution of the initial value problem (27) can be expressed by the fol-

lowing finite series:

$$f(t) = \sum_{n=0}^N a_n (t-t_0)^n, \quad t \in [0, T] \quad (28)$$

The multi-step approach introduces a new idea for constructing the approximate solution. Assume that the interval $[0, T]$ is divided into M subintervals $[t_{m-1}, t_m]$, $m = 0, 1, 2, \dots, M$ of equal step size $h = (T/M)$ by using the nodes $t_m = mh$. The main ideas of the MDTM are as follows. First, we apply the DTM to Equation (24) over the interval $[0, t_1]$, we will obtain the following approximate solution:

$$f_1(t) = \sum_{n=0}^N a_{1n} (t-t_0)^n, \quad t \in [0, t_1] \quad (29)$$

using the initial conditions $f_1^{(k)}(0) = c_k$. For $m \geq 2$ and at each subinterval $[t_{m-1}, t_m]$ we will use the initial conditions $f_m^{(k)}(t_{m-1}) = f_{m-1}^{(k)}(t_{m-1})$ and apply the DTM to Equation (27) over the interval $[t_{m-1}, t_m]$, where t_0 in Equation (23) is replaced by t_{m-1} . The process is repeated and generates a sequence of approximate solutions $f_m(t)$, $m = 1, 2, \dots, M$, for the solution $f(t)$:

$$f_m(t) = \sum_{n=0}^N a_{mn} (t-t_{m-1})^n, \quad t \in [t_m, t_{m-1}] \quad (30)$$

where $N = K \cdot M$. In fact, the MDTM assumes the following solution:

$$f(t) = \begin{cases} f_1(t), & t \in [0, t_1] \\ f_2(t), & t \in [t_1, t_2] \\ \vdots \\ f_M(t), & t \in [t_{M-1}, t_M] \end{cases} \quad (31)$$

The new algorithm, MDTM, is simple for computational performance for all values of h . It is easily observed that if the step size $h = T$, then the MDTM reduces to the classical DTM. As we will see in the next section, the main advantage of the new algorithm is that the obtained series solution converges for wide time regions and can approximate non-chaotic or chaotic solutions.

3.3. Analytical Solution Be the MDTM

By applying the MDTM to Equations (18)-(21), gives the following recursive relations in each sub-domain (t_i, t_{i+1}) , $i = 0, 1, \dots, N-1$.

$$\begin{aligned} & \left(1 + \frac{1}{\beta}\right) (k+1)(k+2)(k+3) F[k+3] + \sum_{r=0}^k (k-r+1)(k-r+2) F[r] F[k-r+2] \\ & + \gamma \left(\delta[k] - \sum_{r=0}^k (r+1)(k-r+1) F[r+1] F[k-r+1] \right) + M \left[\sum_{r=0}^k (k-r+1)(k-r+2) G[r] G[k-r+2] \right] \\ & + M \left[\gamma \left(\delta[k] - \sum_{r=0}^k (r+1)(k-r+1) G[r+1] G[k-r+1] \right) \right] = 0 \end{aligned} \quad (32)$$

$$\begin{aligned} & \lambda (k+1)(k+2)(k+3) G[k+3] + \sum_{r=0}^k (k-r+1)(k-r+2) F[r] G[k-r+2] \\ & - \sum_{r=0}^k (k-r+1)(k-r+2) G[r] F[k-r+2] = 0 \end{aligned} \quad (33)$$

$$\begin{aligned} & \left(1 + \frac{4}{3} R_a\right) (k+1)(k+2) \theta[k+2] + P_r \sum_{r=0}^k (k-r+1) F[r] \theta[k-r+1] \\ & + P_r E \left(1 + \frac{1}{\beta}\right) \sum_{r=0}^k (r+2)(r+1)(k-r+2)(k-r+1) F[r+2] F[k-r+2] \\ & + P_r D_u (k+1)(k+2) \varnothing[k+2] = 0 \end{aligned} \quad (34)$$

$$(k+1)(k+2)\mathcal{O}[k+2] + S_c \sum_{r=0}^k (k-r+1)F[r]\mathcal{O}[k-r+1] + S_c S_o (k+1)(k+2)\theta[k+2] = 0 \quad (35)$$

where $F[k], G[k], \theta[k]$ and $\mathcal{O}[k]$ are the differential transform of $f(\eta), g(\eta), \theta(\eta)$ and $\mathcal{O}(\eta)$.

The differential transformed boundary conditions in Equation (22) to:

$$\begin{aligned} F[0] &= 0, G[0] = 0, \theta[0] = 1, \mathcal{O}[0] = 1, \\ F[1] &= \lambda, F[2] = f_2, G[1] = 0, G[2] = g_2, \theta[1] = t_1, \mathcal{O}[1] = c_1. \end{aligned} \quad (36)$$

where f_2, g_2, t_1 and c_1 are constants. These constants are computed from the boundary condition. Moreover, substituting Equation (36) into Equation (32)-(35) and by using the recursive method, we can calculate other values of $F(k), G(k), \theta(k)$ and $\mathcal{O}(k)$. Hence, substituting all $F(k), G(k), \theta(k)$ and $\mathcal{O}(k)$, into Equation (26), we obtain series solutions.

The velocity, induced magnetic field, temperature and concentration distributes achieved with the aid of MATHAMATICA application software.

3.4. Numerical Solutions by Using the FDM

To solve the ordinary differential Equations (18)-(21), we use a finite difference based numerical algorithm. We reduce the order of Equations (18) and (19) by one with the help of the substitution:

$$q = \frac{df}{d\eta} = f', \quad p = \frac{dg}{d\eta} = g' \quad (37)$$

Equations (18)-(21) in view of Equation (37) can be written as:

$$\left(1 + \frac{1}{\beta}\right)q'' + fq' + \gamma(1 - q^2) - M \left[gp' + \gamma(1 - p'^2)\right] = 0 \quad (38)$$

$$\lambda p'' + fp' - pq' = 0 \quad (39)$$

$$\left(1 + \frac{4}{3}R_a\right)\theta'' + P_r f \theta' + P_r E \left(1 + \frac{1}{\beta}\right)q'^2 + P_r D_u \mathcal{O}'' = 0 \quad (40)$$

$$\mathcal{O}'' + S_c f \mathcal{O}' + S_c S_o \theta'' = 0 \quad (41)$$

The boundary conditions:

$$\begin{aligned} f(0) &= 0, q(0) = \delta, g(0) = 0, p(0) = 0, \theta(0) = 1, \mathcal{O}(0) = 1, \\ q(\infty) &= 1, p(\infty) = 1, \theta(\infty) = 0, \mathcal{O}(\infty) = 0. \end{aligned} \quad (42)$$

Now, we can write Equations (37)-(41) in the following finite difference form

$$\left(1 + \frac{1}{\beta}\right) \left(\frac{q_{i+1} - 2q_i + q_{i-1}}{(\Delta\eta)^2} \right) + f_i \left(\frac{q_{i+1} - q_{i-1}}{2\Delta\eta} \right) + \gamma(1 - q_i^2) - M \left[g_i \left(\frac{p_{i+1} - p_{i-1}}{2\Delta\eta} \right) + \gamma(1 - p_i^2) \right] = 0 \quad (43)$$

$$\lambda \left(\frac{p_{i+1} - 2p_i + p_{i-1}}{(\Delta\eta)^2} \right) + f_i \left(\frac{p_{i+1} - p_{i-1}}{2\Delta\eta} \right) - g_i \left(\frac{q_{i+1} - q_{i-1}}{2\Delta\eta} \right) = 0 \quad (44)$$

$$\begin{aligned} \left(1 + \frac{4}{3}R_a\right) \left(\frac{\theta_{i+1} - 2\theta_i + \theta_{i-1}}{(\Delta\eta)^2} \right) + P_r f_i \left(\frac{\theta_{i+1} - \theta_{i-1}}{2\Delta\eta} \right) + P_r E \left(1 + \frac{1}{\beta}\right) \left(\frac{q_{i+1} - q_{i-1}}{2\Delta\eta} \right)^2 \\ + P_r D_u \left(\frac{\mathcal{O}_{i+1} - 2\mathcal{O}_i + \mathcal{O}_{i-1}}{(\Delta\eta)^2} \right) = 0 \end{aligned} \quad (45)$$

$$\left(\frac{\mathcal{O}_{i+1} - 2\mathcal{O}_i + \mathcal{O}_{i-1}}{(\Delta\eta)^2} \right) + S_c f_i \left(\frac{\mathcal{O}_{i+1} - \mathcal{O}_{i-1}}{2\Delta\eta} \right) + S_c S_o \left(\frac{\theta_{i+1} - 2\theta_i + \theta_{i-1}}{(\Delta\eta)^2} \right) = 0 \quad (46)$$

The velocity, induced magnetic field, temperature and concentration distributes at all interior nodal points computed by successive applications of the above finite difference equations and these are achieved with the aid of MATLAB application software.

4. Results and Discussion

The system of the equations which describe the motion of Casson fluid in the boundary layer with heat and mass transfer and in the presence of induced magnetic field with radiation are solved analytically by using (DTM) and numerically by using (FDM). The values of the velocity, temperature, concentration and induced magnetic field are obtained as a functions of the physical parameters of the problem such as Casson parameter β , wedge parameter γ , magnetic parameter M , magnetic Prandtl number λ , radiation parameter R_a , Prandtl number P_r , Eckert number E , Dufour parameter D_u , Schmidt number S_c , and Soret number S_0 . These effects are illustrated graphically through a set of (1 - 18).

Figures 1-5 represent the velocity profiles for the flow parameters. The effect of β on the velocity profiles is presented in **Figure 1**. It is seen that the velocity increases when β increases. **Figure 2** clears the effect of γ on the velocity profiles. It is shown that the velocity increases with increasing γ . **Figure 3** clears the effect of M on the velocity profiles. It is shown that the velocity decreases with increasing M , also it decrease with increasing λ as shown in **Figure 4**. The effect of δ on the velocity profiles is presented in **Figure 5**. It is shown that the velocity increases with increasing δ . **Figures 6-10** represent the induced magnetic field profiles for the flow parameters. It is observed that the induced magnetic field increases with increasing β, γ and δ , while decreases with increasing M and λ . **Figures 11-15** represent the temperature profiles for the flow parameters. It is observed that the temperature increases with increasing R_a, E and D_u , while decreases with increasing β and P_r . **Figures 16-18** represent the concentration profiles for the flow parameters. It is noted that concentration increases with an increases S_0 , while it decrease with the increases β and S_c .

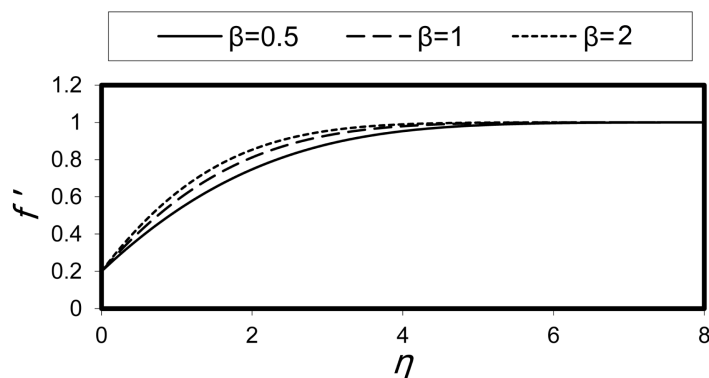


Figure 1. Velocity profiles $f'(\eta)$ for some values of β .

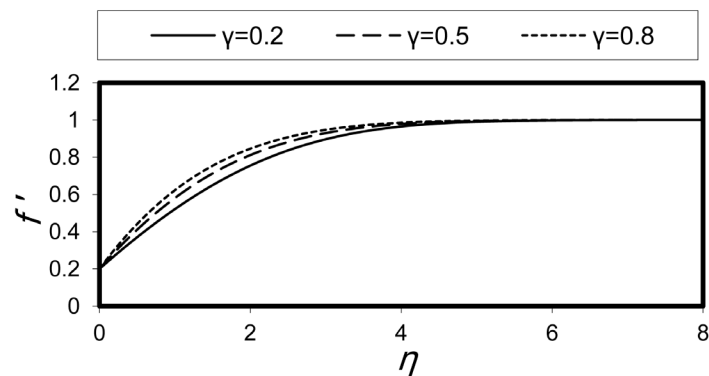


Figure 2. Velocity profiles $f'(\eta)$ for some values of γ .

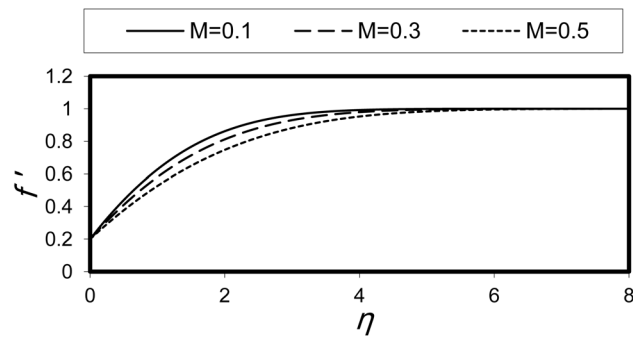


Figure 3. Velocity profiles $f'(\eta)$ for some values of M .

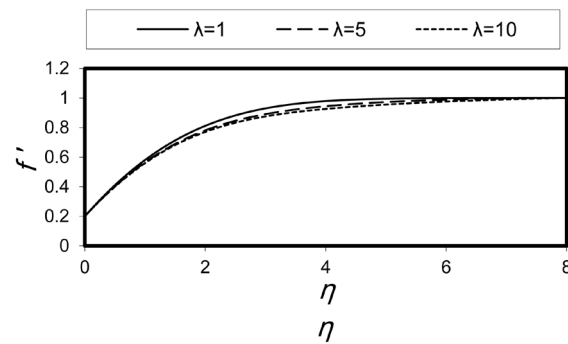


Figure 4. Velocity profiles $f'(\eta)$ for some values of λ .

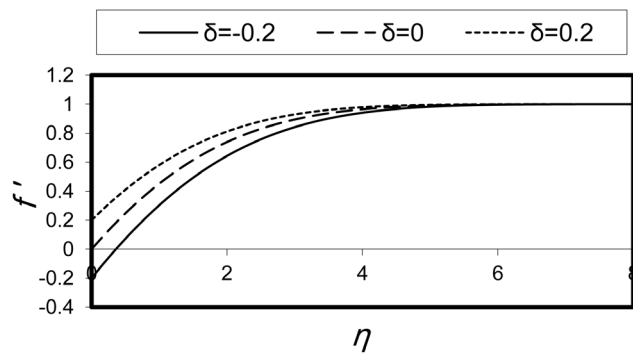


Figure 5. Velocity profiles $f'(\eta)$ for some values of δ .

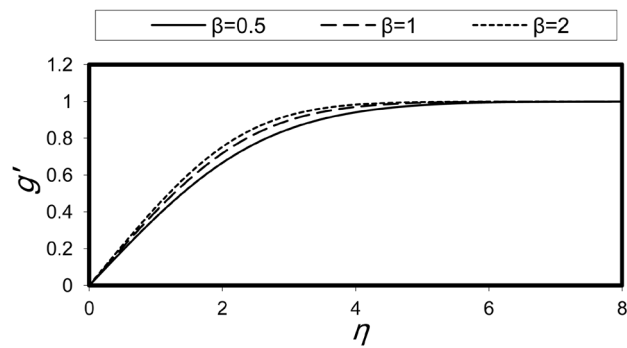


Figure 6. Induced magnetic profiles $g'(\eta)$ for some values of β .

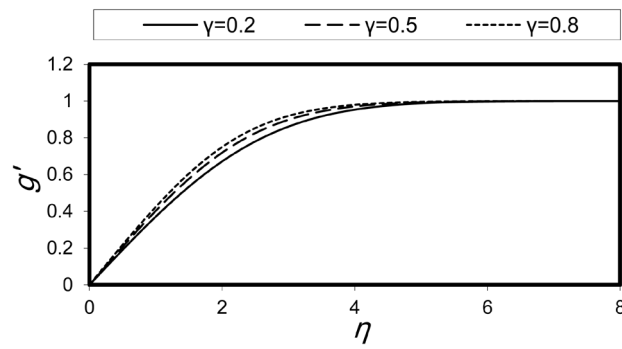


Figure 7. Induced magnetic profiles $g'(\eta)$ for some values of γ .

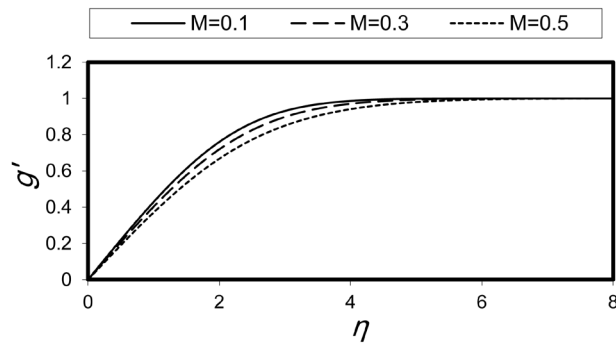


Figure 8. Induced magnetic profiles $g'(\eta)$ for some values of M .

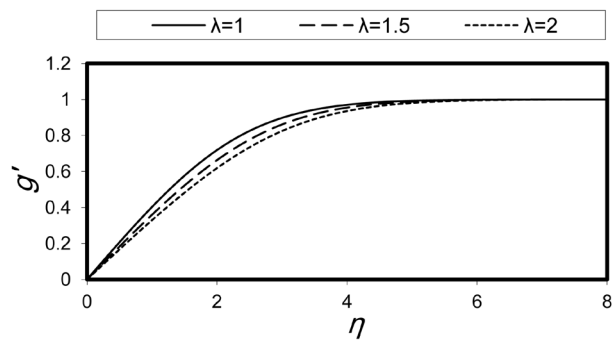


Figure 9. Induced magnetic profiles $g'(\eta)$ for some values of λ .

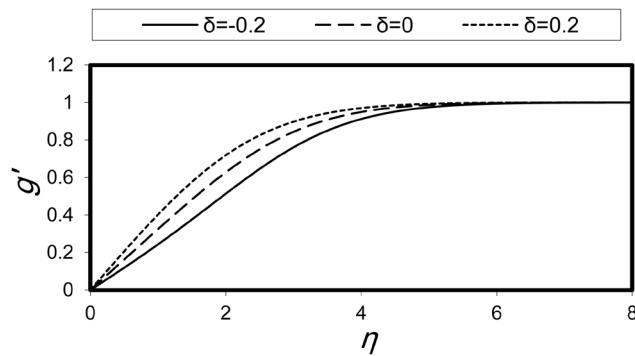


Figure 10. Induced magnetic profiles $g'(\eta)$ for some values of δ .

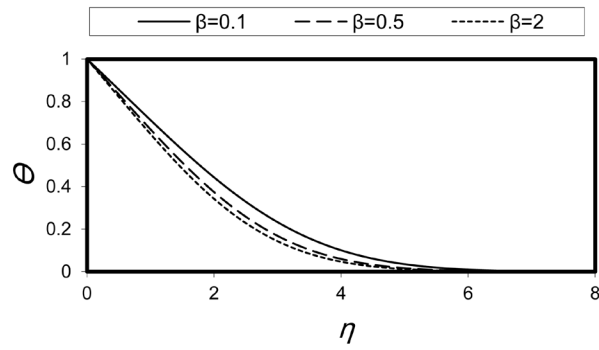


Figure 11. Temperature profiles $\theta(\eta)$ for some values of β .

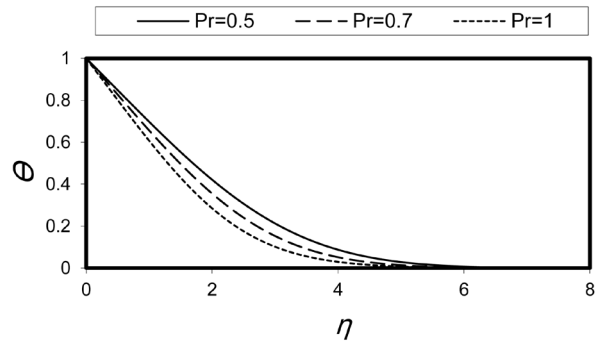


Figure 12. Temperature profiles $\theta(\eta)$ for some values of Pr .

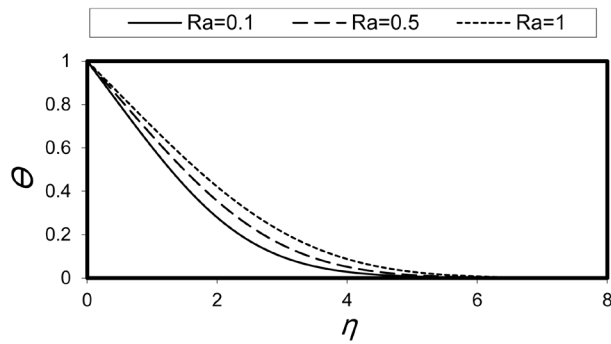


Figure 13. Temperature profiles $\theta(\eta)$ for some values of Ra .

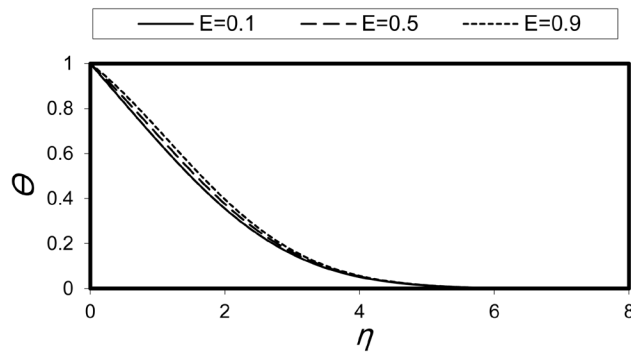


Figure 14. Temperature profiles $\theta(\eta)$ for some values of E .

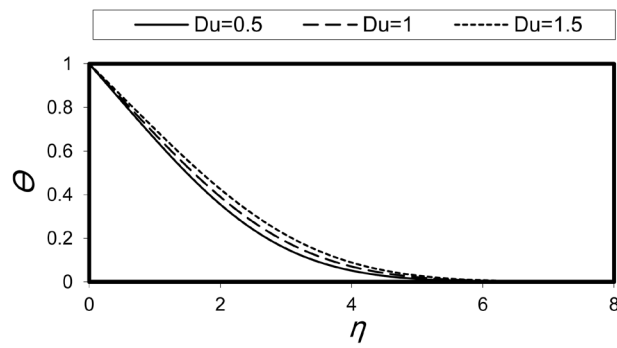


Figure 15. Temperature profiles $\theta(\eta)$ for some values of D_u .

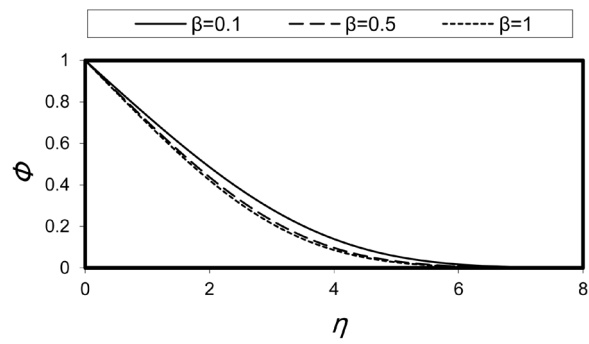


Figure 16. Concentration profiles $\phi(\eta)$ for some values of β .

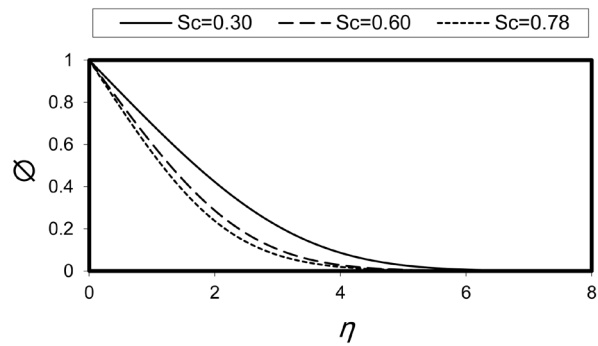


Figure 17. Concentration profiles $\phi(\eta)$ for some values of Sc .

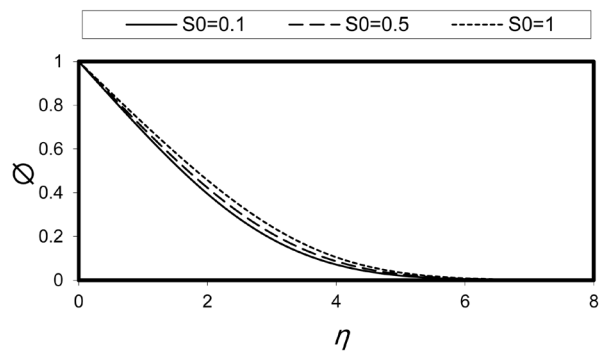


Figure 18. Concentration profiles $\phi(\eta)$ for some values of S_0 .

In order to verify the accuracy of the numerical solution of present work by using (FDM) we have compared these results for $f''(0)$ when $\delta = 0$ and $M = 0$ with Rajagopal *et al.* [16], Kuo [17], Ishak *et al.* [18], K. Jafar [8]. It is observed that this numerical solution in excellent agreement with the results previous obtained values as shown in Table 1. Finally, Tables 2-4 shows comparison between numerical solution by using (FDM) and the analytical solution using (DTM). It is observed that this approximate numerical solution is in excellent agreement with the results of the analytical solution of by using (DTM).

Table 1. Comparison between the previous obtained values of $f''(0)$ by Rajagopal *et al.* [16], Kuo [17], Ishak *et al.* [18], K. Jafar [8] and our results when $\delta = 0$ and $M = 0$.

| γ | Rajagopal <i>et al.</i> [16] | Kuo [17] | Ishak <i>et al.</i> [18] | K. Jafar [8] | Present result |
|----------|------------------------------|----------|--------------------------|--------------|----------------|
| 0.0 | - | 0.469600 | 0.4696 | 0.4696 | 0.46960075 |
| 0.1 | 0.587035 | 0.587880 | 0.5870 | 0.5871 | 0.5870358 |
| 0.3 | 0.774755 | 0.775524 | 0.7748 | 0.7748 | 0.7747554 |
| 0.5 | 0.927680 | 0.927905 | 0.9277 | 0.9277 | 0.9276813 |
| 1.0 | 1.232585 | 1.231289 | 1.2326 | 1.2326 | 1.232590 |

Table 2. Comparison between the previous obtained values of $f(\eta)$ and $f'(\eta)$ where $\beta = 1, \gamma = 0.5, M = 0.3, \lambda = 1, \delta = 0.2, P_r = 0.7, R_a = 0.5, E = 0.1, D_u = 0.5, S_c = 0.3$ and $S_o = 0.5$.

| η | $f(\eta)$ | | | $f'(\eta)$ | | |
|--------|----------------------|----------|-----------|-----------------------|----------|------------|
| | DTM | MDTM | FDM | DTM | MDTM | FDM |
| 0.0 | 0.0 | 0.0 | 0.0 | 0.2 | 0.2 | 0.2 |
| 1.0 | 0.403766 | 0.403766 | 0.4037649 | 0.580616 | 0.580616 | 0.58061605 |
| 2.0 | 1.11098 | 1.11098 | 1.1109816 | 0.811421 | 0.811421 | 0.81142107 |
| 3.0 | 1.98901 | 1.98901 | 1.9890103 | 0.929608 | 0.929609 | 0.92960949 |
| 4.0 | 2.95048 | 2.94735 | 2.9473479 | 1.0034 | 0.97895 | 0.97895070 |
| 5.0 | 6.32369 | 3.93601 | 3.9360082 | 14.7103 | 0.995059 | 0.99505878 |
| 6.0 | 415.231 | 4.93359 | 4.9335872 | 1902.95 | 0.999104 | 0.99910402 |
| 7.0 | 28673.5 | 5.93319 | 5.9331922 | 111857 | 0.999884 | 0.99988388 |
| 8.0 | 1.0757×10^6 | 6.93315 | 6.9331527 | 3.62989×10^6 | 1.0 | 1.0 |

Table 3. Comparison between the previous obtained values of $g(\eta)$ and $g'(\eta)$ where $\beta = 1, \gamma = 0.5, M = 0.3, \lambda = 1, \delta = 0.2, P_r = 0.7, R_a = 0.5, E = 0.1, D_u = 0.5, S_c = 0.3$ and $S_o = 0.5$.

| η | $g(\eta)$ | | | $g'(\eta)$ | | |
|--------|------------------------|----------|-----------|------------------------|----------|-----------|
| | DTM | MDTM | FDM | DTM | MDTM | FDM |
| 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 |
| 1.0 | 0.20504 | 0.20504 | 0.2050395 | 0.402848 | 0.402848 | 0.4028479 |
| 2.0 | 0.776432 | 0.776432 | 0.7764307 | 0.718855 | 0.718855 | 0.7188556 |
| 3.0 | 1.59576 | 1.59579 | 1.5957888 | 0.897193 | 0.897483 | 0.974845 |
| 4.0 | 2.45379 | 2.53591 | 2.5359114 | 0.411927 | 0.970245 | 0.9702454 |
| 5.0 | -31.0684 | 3.51998 | 3.5199796 | -185.501 | 0.993106 | 0.9931059 |
| 6.0 | -4620.89 | 4.51661 | 4.5166091 | -20617.8 | 0.998756 | 0.9987563 |
| 7.0 | -281829 | 5.51606 | 5.5160617 | -1.07009×10^6 | 0.99984 | 0.9998396 |
| 8.0 | -9.71748×10^6 | 6.51601 | 6.5160072 | -3.21308×10^7 | 1.0 | 1.0 |

Table 4. Comparison between the previous obtained values of $\theta(\eta)$ and $\varnothing(\eta)$ where $\beta = 1, \gamma = 0.5, M = 0.3, \lambda = 1, \delta = 0.2, P_r = 0.7, R_a = 0.5, E = 0.1, D_u = 0.5, S_c = 0.3$ and $S_o = 0.5$.

| η | $\theta(\eta)$ | | | $\varnothing(\eta)$ | | |
|--------|----------------|-------------|-------------|---------------------|------------|------------|
| | DTM | MDTM | FDM | DTM | MDTM | FDM |
| 0.0 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 |
| 1.0 | 0.65649 | 0.65649 | 0.6564895 | 0.696176 | 0.696176 | 0.6961763 |
| 2.0 | 0.355842 | 0.355842 | 0.3558407 | 0.422562 | 0.422562 | 0.4225611 |
| 3.0 | 0.153485 | 0.15349 | 0.1534889 | 0.213823 | 0.213822 | 0.2138215 |
| 4.0 | 0.0417624 | 0.052241 | 0.052240 | 0.0885228 | 0.0866998 | 0.0866989 |
| 5.0 | -3.56762 | 0.0141772 | 0.01417667 | 0.649388 | 0.0273532 | 0.0273525 |
| 6.0 | -393.027 | 0.00306269 | 0.00306253 | 67.9412 | 0.00653717 | 0.00653686 |
| 7.0 | -20071.9 | 0.000478198 | 0.000478176 | 3451.29 | 0.00107547 | 0.00107538 |
| 8.0 | -592012 | 0.0 | 0.0 | 101282 | 0.0 | 0.0 |

5. Conclusions

In this work we have obtained the numerical solution of MHD boundary layer flow of non-Newtonian Casson fluid on a moving wedge with heat and mass transfer and induced magnetic field. The effects of thermal diffusion and diffusion thermo with induced magnetic field are taken into consideration. The resulting partial differential equations which describe the problem are transformed into ordinary differential equations by using a similarity transformation and then solved numerically by using the finite difference method (FDM) and analytically by using (DTM). A representative set of numerical results for velocity, induced magnetic field, temperature and concentration profiles is presented graphically and discussed. The figures and tables clearly show that the results by using (FDM) are in excellent agreement with previously published works and with the results of the analytical solution by using (DTM). The important results for this study summarized as follows:

- The velocity distribution decreases with the increase of M and λ , while it increases with an increase of β, γ and δ .
- The temperature distribution increases with the increase of R_a, E and D_u , while it decreases with the increases of P_r and β .
- The induced magnetic field distribution increases with the increase of β, γ and δ , while it decreases with the increases of M and λ .
- The concentration distribution increases with increase of S_0 , while it decreases with the increases of S_c and β .

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