

On Universal Mechanics and Superluminal Velocities*

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Abstract

In this work we continue to set up the theory of universal space and time and derive the Euclidean form of the scaling transformations. Two types of velocities emerge, inertial and universal, with the former bound by the light velocity c whereas the latter is unbound, and may accommodate consistently particles' velocities possibly exceeding c . The inertial velocity is the ratio of the simultaneous source's displacement and the corresponding length of the light trip to the observer, whereas the universal velocity has its familiar meaning for motion in a synchronous inertial frame. Defining the momentum as the product of universal velocity and mass, and utilizing the already established mass-energy equivalence, the mechanics constructed on the bases of the new concepts, named universal mechanics, admits superluminal velocities but yet coincides with the relativistic mechanics in its basic dynamical components and their inter-relations. The possibility of superluminal velocities provides a straight forward explanation of the presence of the μ -meson particles abundantly at the sea level despite their generation at high altitude and their short life-time.

Keywords

Inertial Velocity, Superluminal Speeds, Universal Mechanics

1. Introduction

The basic concepts and main structure of the theory of the universal space and time (UST) were discussed in [1]. It was argued that time intervals are essentially measured by spatial intervals associated with light trips. A light trip follows the same path in the universal space which can be identified with any inertial frame, one at a time.

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The duration and direction of a light trip are the same, but its geometric lengths differ from a frame to another. The scaling transformations (STs) determine the relation between the geometric lengths of the same light trip in two inertial frames. The STs were shown to map intrinsic units in a frame on intrinsic units in another, and were utilized to explain the presence of the μ -meson particles which are generated at a high latitude at the sea level despite their short lifetime. It was also shown that in UST the longitudinal Doppler's effect coincides with the relativistic one, but the transverse effect at right angle is absolutely absent in UST.

In the current work which continues constructing UST we furnish the following points:

- We derive the scaling transformations in a concise and transparent way illuminating its Euclidean form.
- Two types of velocity are distinguished, inertial and universal, with the former bound by light velocity, while the latter is unbound and can exceed c .
- The inertial velocity is defined by the quotient of the distance travelled by a body to the corresponding distance travelled by a light pulse. The universal velocity is the quotient of the distance travelled by the body to the corresponding time interval in a universal space.
- Defining the particle's momentum by the product of its mass and universal velocity, we obtain all basic relativistic mechanical expressions and relations [2]-[6].
- The paradox of the presence of the meta-stable particles at sea level finds its natural explanation in superluminal speeds of these particles.

In future works, the change of direction of a light trip between two frames and its application to explain the stellar aberration will be presented. Also, a second type of scaling transformation (STII) will be derived and applied to explain the drag effect, Sagnac effect, Michelson and Morley experiment, Michelson and Gale experiment.

2. The Euclidean Form of the Scaling Transformations

Assume that a body b is moving in the universal frame S at a velocity $u = ui$ and $B \in S$ is a body at rest in S at a geometric time distance T seconds from the observer $O \in S$. Suppose that at an instant $t = 0$ corresponding to b passing by $B \in S$, a spherical light wave emanates from the observer $O \in S$ (Figure 1). The wave arrives at B at an instant T and intercepts the body b at a point $b' \in S$ at an instant t . Because the geometry of the universal frame is Euclidean by hypothesis, the sides' lengths of the triangle OBb' satisfy all triangle relations in Euclidean trigonometry. In particular, the relation, $Ob' = OB + Bb'$, has its familiar meaning in Euclidean geometry. Since O and B are at rest in S , we have $OB = cT(-e)$, but since the body b is moving in S we set provisionally $Bb' = \gamma uti$ and $Ob' = \gamma ct(-e_L)$, where γ is a factor that depends on the magnitude of the velocity of the body b and must go to 1 as u goes to zero. The vectors

$$OB = cT(-e), \quad Bb' = \gamma tu, \quad Ob' = \gamma ct(-e_L), \tag{2.1}$$

satisfy the relation

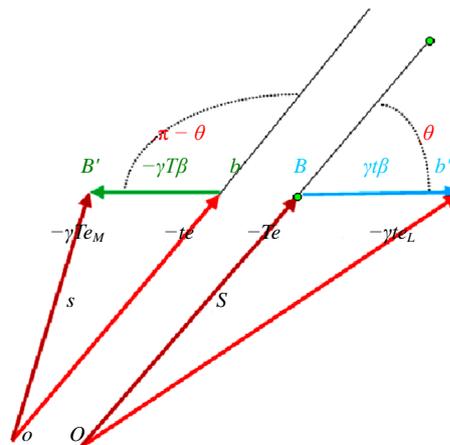


Figure 1. The light's trips to stationary and moving targets as seen from two frames when each is considered universal.

$$\gamma t(-\mathbf{e}_L) = T(-\mathbf{e}) + \gamma t\boldsymbol{\beta}, \tag{2.2}$$

where $\boldsymbol{\beta} = \mathbf{u}/c$ Solving for t we obtain

$$t = \gamma^{-1} E(\boldsymbol{\beta}, \pi - \theta) T, \tag{2.3}$$

where

$$E(\pi - \theta, \boldsymbol{\beta}) = \frac{\boldsymbol{\beta} \cos \theta + \sqrt{1 - \boldsymbol{\beta}^2 \sin^2 \theta}}{1 - \boldsymbol{\beta}^2} \tag{2.4}$$

We shall call $E(\pi - \theta, \boldsymbol{\beta})$ the *Euclidean factor*.

Since time in the UST theory is a frame-independent entity we therefore seek the transformation from a universal frame S to a frame s co-moving with b such that the transformation and its inverse yield time durations as frame-independent. The latter demand is essential to enable one to start with either S or a frame co-moving with b (but not both) as universal. Consider thus an inertial frame s moving in the universal frame S at velocity u ; the body b will be at rest in s . Let $o \in s$ be the s -observer that is contiguous to $O \in S$ at the instant $t = 0$, *i.e.*, when the spherical wave emanates from $O \in S$. While b remains at rest in the moving frame s , the body $B \in S$ moves at velocity $(-\mathbf{u})$. In the frame s , the wave hits b at the instant t and intercepts B at the instant T at a position $B' \in s$. We seek thus the transformation that maps the triangle OBb' in S to the triangle $oB'b$ in s , with the time lengths of OB and Ob' are equal to the time lengths of oB' and ob respectively, and vice versa. Since the role of the ordered pair (t, T) in s must be identical to the role of the pair (T, t) in S , the relation between t and T in s must result from that in S by interchanging t and T in (2.3) and replacing \mathbf{u} by $(-\mathbf{u})$ or equivalently $\pi - \theta$ by θ . This yields

$$T = \gamma^{-1} E(\boldsymbol{\beta}, \theta) t. \tag{2.5}$$

On the other hand, the quantities t and T already satisfy (2.3). Substituting one of the Equations (2.3) or (2.5) in the other yields

$$1 = \gamma^{-2} E(\pi - \theta, \boldsymbol{\beta}) E(\theta, \boldsymbol{\beta}) = \frac{1}{\gamma^2 (1 - \boldsymbol{\beta}^2)} \tag{2.6}$$

or

$$\gamma \equiv 1 / \sqrt{1 - \boldsymbol{\beta}^2} \tag{2.7}$$

The required transformation is therefore

$$t = \Gamma(\theta, \boldsymbol{\beta}) T \equiv \sqrt{1 - \boldsymbol{\beta}^2} E(\pi - \theta, \boldsymbol{\beta}) T \tag{2.8a}$$

or

$$\frac{t}{\sqrt{1 - \boldsymbol{\beta}^2}} = E(\pi - \theta, \boldsymbol{\beta}) T. \tag{2.8b}$$

We call the function $\Gamma(\theta, \boldsymbol{\beta})$ the *scaling factor* and the relation $t = \Gamma(\theta, \boldsymbol{\beta}) T$ the *scaling transformation* (ST). The relation (2.8b) is called the *Euclidean form of the scaling transformation*. The inverse of the latter relation is obtained through interchanging t and T and replacing θ by $\pi - \theta$ (or $\boldsymbol{\beta}$ by $-\boldsymbol{\beta}$) which amounts to (2.8) itself. The ST allows for each frame to be taken stationary, and it holds within the same frame whether it was the stationary or the moving frame, as well as between the universal and moving frame. Note that the geometric lengths of the corresponding trips in S and s are different.

It was shown in [1] that if a light wave emanates from b when at B then the period it takes to arrive at O is also given by (2.8), *i.e.* if a light wave emanates from b simultaneously with a wave emanating from O then when the former arrives at O , the latter arrives at b' .

3. The Euclidean Body-Observer Triangle

The initial and final positions of the moving body in the universal frame S together with the observer's position form a Euclidean triangle (Figure 2) with sides length's

$$T, \quad T'' = \frac{\beta t}{\sqrt{1-\beta^2}}, \quad T' = \frac{t}{\sqrt{1-\beta^2}}. \quad (3.1)$$

The given lengths satisfy all triangle relations in Euclidean geometry, and yield a value t for the time duration as it is prescribed by the STI. An elementary fact asserts that when three numbers are legitimate to form a triangle, this triangle is unique (up of course to arbitrary rotations, translations, or reflections). Thus the latter values determine a unique Euclidean triangle; it is the *body-observer triangle*.

By (3.1) we have

$$\frac{T''}{\beta} = \frac{T'}{1} = \frac{t}{\sqrt{1-\beta^2}}, \quad (3.2)$$

which are equivalent to the relations

$$T'' = \beta T', \quad T'^2 - T''^2 = t^2. \quad (3.3)$$

In terms of the initial geometric distance T we have

$$\frac{T''}{\beta} = \frac{T'}{1} = \frac{\Gamma(\beta, \theta) T}{\sqrt{1-\beta^2}} \equiv E(\beta, \pi - \theta) T. \quad (3.4)$$

By the sinuses law in trigonometry,

$$\frac{\gamma \beta t}{\sin \delta} = \frac{\gamma t}{\sin \theta} = \frac{T}{\sin(\theta - \delta)}, \quad (3.5)$$

we have

$$\sin \delta = \beta \sin \theta, \quad \sin(\theta - \delta) = \sin \theta / E(\beta, \pi - \theta) \quad (3.6)$$

By (3.2) the pair of sides (T'', T') are in 1-1 correspondence with the pair (β, t) , and the body-observer triangle is thus determined by (β, t, T) . This expresses the obvious fact that the direction of the vector β relative to the observer, *i.e.* θ , is determinable by the quantities (β, t, T) through the scaling transformations $\Gamma(\beta, \theta) = t/T$, by which we can determine one out of the quantities $\{T, t, \beta, \theta\}$ in terms of the remaining three. In other words, the body-observer triangle is fully determined by three out of the four variables T, t, β, θ .

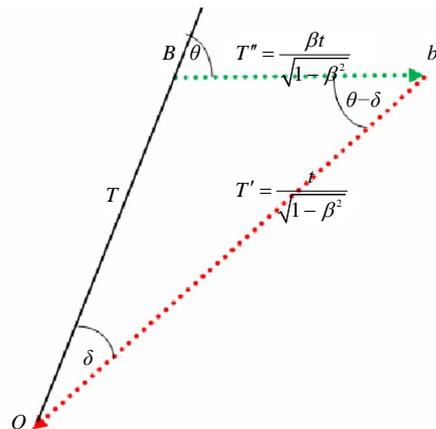


Figure 2. The body-observer triangle.

This implies that, in correspondence with each body-observer triangle there is one value of β , and hence the same value of β is obtained whether calculated from the expressions of T' or T'' or from the STI.

The Inertial Velocity: By (3.2), the displacement of the source, the distance travelled by the signal, and their duration t are in 1-1 correspondence. The quantity

$$\beta = T''/T' \quad (3.7)$$

which is obtained from (3.2), will be called the *inertial velocity* of the body b . The definition (3.7) expresses the inertial speed in terms of geometric distances; *it is the quotient of the distance T'' travelled by the body b to the distance T' travelled by the light emanating from the observer O when they intercept each other.* The initial time for both motions is the instant (b at B). By [1], T' is also the distance travelled by the pulse emitted from b when at B till arriving at O .

4. The Universal Mechanics

The *universal velocity* of a body b refers to its velocity in a universal space; it is defined exactly as the familiar velocity in classical mechanics:

$$U = |\mathbf{Ob}' - \mathbf{OB}|/t = |\mathbf{Bb}'|/t \quad (4.1)$$

The universal velocity of the body b is

$$U = c\beta_U \quad (4.2)$$

where by (3.2)

$$\beta_U \equiv \frac{T''}{t} = \frac{\beta}{\sqrt{1-\beta^2}}. \quad (4.3)$$

It is clear that β_U tends to zero with β and it can assume any non-negative value with no upper bound. The *momentum* of the particle b is defined by the product of its mass m and universal velocity:

$$\mathbf{p} = m\mathbf{U} = \frac{mc\boldsymbol{\beta}}{\sqrt{1-\beta^2}} = \frac{m\mathbf{u}}{\sqrt{1-(u/c)^2}} \quad (4.4)$$

Multiplying both sides of the identity

$$\frac{\beta^2}{1-\beta^2} = \frac{1}{1-\beta^2} - 1, \quad (4.5)$$

by m^2 we obtain

$$\frac{p^2}{c^2} = \frac{m^2}{1-\beta^2} - m^2 \quad (4.6)$$

In the reduced system of units (RSUI \equiv MKM) [1] $c = 1$, mass, energy and momentum are all measured by the same unit, *kilogram*. Thus, the *right hand-side* of (4.6) can be envisaged as a difference between the squares of two values of the mass or energy of the moving body corresponding to the states of motion and rest respectively. Denoting these values by E ($\equiv M$) and E_0 ($\equiv m$) respectively, *i.e.*,

$$E(\text{kg}) \equiv M(\text{kg}) = \frac{m}{\sqrt{1-\beta^2}}(\text{kg}), \quad E_0(\text{kg}) \equiv m(\text{kg}), \quad (4.7)$$

we write (4.6) in the form

$$p^2 = E^2 - E_0^2 = M^2 - m^2 \quad (4.8)$$

The latter relation reads: the state of motion of a body with rest mass m that is characterized by a momen-

tum of magnitude p is accompanied by a *total energy, or kinetic mass*,

$$E = \frac{m}{\sqrt{1-\beta^2}} = \sqrt{m^2 + p^2}. \quad (4.9)$$

When p goes to zero, the total energy (or kinetic mass) tends to the rest energy (or rest mass) $E_0 = m$. The Hamiltonian of the particle coincides with its total energy:

$$H = \sqrt{m^2 + p^2} \quad (4.10)$$

In MKM, $U = \beta_U$, $u = \beta$, $\mathbf{p} = m\beta_U$. The force acting on a particle is defined as in Newtonian mechanics by the time rate of the change in its momentum:

$$\mathbf{f} = \frac{d\mathbf{p}}{dt} = \frac{d}{dt}(m\beta_U) = \frac{d}{dt} \frac{m\beta}{\sqrt{1-\beta^2}} \quad (4.11)$$

If the magnitude of β_U remains constant while its direction changes,

$$\mathbf{f} = \frac{m}{\sqrt{1-\beta^2}} \frac{d\beta}{dt} = M \frac{d\beta}{dt} \quad (4.12a)$$

If only the magnitude of β_U changes while its direction remains constant,

$$\mathbf{f} = \frac{m d\beta/dt}{(1-\beta^2)^{3/2}} = m \frac{\beta_U^3}{\beta^3} \frac{d\beta}{dt} = \frac{M^3}{m^2} \frac{d\beta}{dt}. \quad (4.12b)$$

If both the magnitude and direction of β_U change then

$$\mathbf{f} = M \left(1 + \frac{M^2}{m^2} \right) \frac{d\beta}{dt} \quad (4.13)$$

Differentiating both sides of the equation $M^2 = m^2 + p^2$ with respect to time, we obtain

$$\frac{dM}{dt} = \frac{1}{M} \mathbf{p} \cdot \mathbf{f} = \mathbf{u} \cdot \mathbf{f} = \frac{\beta_U \mathbf{f}}{\sqrt{1+\beta_U^2}}. \quad (4.14)$$

The relation (4.14) determines the instantaneous rate at which the mass changes under the action of a force when moving at velocity U . Looking on M as the total energy of the particle, the Equation (4.14) also determines the power of the force, *i.e.* the rate at which it does work. The work done by the force during a displacement $d\mathbf{r} = \beta_U dt$ is given by

$$dW = \frac{\mathbf{f} \cdot d\mathbf{r}}{\sqrt{1+U^2/c^2}}. \quad (4.15)$$

It is customary to measure mass in kg, energy in (c^{-2} kg = Joule), and momentum in (c^{-1} kg = $m \cdot \text{sec}^{-1} \cdot \text{kg}$), which corresponds to using the reduced system of units $\text{RSUII} \equiv \text{MKcM}$ in which the expressions (4.7) become

$$E = \frac{mc^2}{\sqrt{1-\beta^2}} (c^{-2} \text{ kg}), \quad E_0 = mc^2 (c^{-2} \text{ kg}) \quad (4.16)$$

And Equations (4.9) and (4.10) become

$$E = \frac{mc^2}{\sqrt{1-\beta^2}} = \sqrt{m^2 c^4 + p^2 c^2}, \quad (4.17)$$

$$H = c\sqrt{m^2c^2 + p^2}. \quad (4.18)$$

The Equation (4.14) which gives the rate of kinetic energy change is written as follows

$$\frac{dE}{dt} = \frac{c^2}{E} \mathbf{p} \cdot \mathbf{f} = \mathbf{u} \cdot \mathbf{f} = \frac{c\mathbf{U} \cdot \mathbf{f}}{\sqrt{c^2 + U^2}}. \quad (4.19)$$

And the work done by the force during a displacement $d\mathbf{r} = \mathbf{U}dt$ is written as

$$dW = dE = \frac{c\mathbf{f} \cdot d\mathbf{r}}{\sqrt{c^2 + U^2}}. \quad (4.20)$$

From its definition, the *momentum of a particle* is related to its universal velocity and energy by

$$\mathbf{p} = \frac{E_0}{c^2} \mathbf{U} = \frac{E}{c^2} \mathbf{u}. \quad (4.21)$$

Or equivalently

$$E_0 = \frac{c^2}{U} p, \quad E = \frac{c^2}{u} p. \quad (4.22)$$

For *particles* travelling at a universal velocity $\mathbf{U} = \mathbf{c}$, we have

$$E_0 = cp \quad (4.23)$$

We recall that \mathbf{U} is unbound in magnitude, and hence, the value c is attainable. For such particles

$$E = \sqrt{2}E_0 \quad (4.24)$$

The energy needed to give a stationary particle a universal velocity c is

$$E - E_0 = (\sqrt{2} - 1)mc^2 = 0.414mc^2$$

In the relations (4.21) and (4.23) the rest energy (or rest mass) is directly measurable, and the momentum of the particle is known if the body's velocity is known. For particles that can exist only in a state of motion, like photons in vacuum, it is meaningless to talk in S , which is the observer frame, about the photon's rest mass or rest energy. However, the relation (4.21), written in terms of inertial velocity

$$\mathbf{p} = \frac{E}{c^2} \mathbf{u} \quad (4.25)$$

and applicable to material particle, may be extended to comprise photons whose velocity is c in S . For material particles, both E and \mathbf{p} are determined on knowing m and \mathbf{u} , which are immediately measurable. Setting $u = c$ in the latter relation we obtain for photons

$$E = \mathbf{c} \cdot \mathbf{p} = cp \quad (4.26)$$

Unlike the relation (4.23) which connects directly measurable quantities, the latter relation which extends (4.25) to photons, requires in order to be fully meaningful a way by which either the energy or the momentum of a photon is prescribed. It is not satisfactory to merely deduce the energy and the momentum of the photon through its earlier or subsequent interaction with an external system. The contrary is required; the theory should be capable to quantify the photon's energy and momentum in terms of inherent characteristics, and thus capable of predicting the magnitudes of its interaction with an external system. Moreover, and whereas (4.23) assigns to a particle with a rest energy E_0 a definite momentum, the relation (4.25) does not distinguish between photons; it only states the relation between their energy and momentum. The additional information about the energy and momentum of a photon comes from Plank hypothesis which relate the energy (or momentum) of a photon to a macroscopic measurable quantity, the frequency ν of photon, by $E = h\nu$ (or $p = h\nu/c = h/\lambda$).

5. Point-Wise Measurement of the Inertial Velocity

Suppose that the free source of light b is detected at a point $B \in S$ at $t=0$ and at $b' \in S$ at an instant of time t . The pulse which was emitted from b when was at B arrives at the instant t at $O \in S$. *i.e.*, the arrivals of b at b' and the pulse at O are simultaneous; they both occur at t . Let $T'' = |\mathbf{Bb}'|$ be the geometric length of the displacement \mathbf{Bb}' . The source b , at the time t , is at geometric distance

$$\frac{\beta t}{\sqrt{1-\beta^2}} \equiv T'' \quad (5.1)$$

from $B \in S$. This yields

$$\beta = \frac{1}{\sqrt{1+(t/T'')^2}} \quad (5.2)$$

The latter formula determines the inertial velocity of a moving body b in S in terms of the time t read at its final position $b' \in S$ and the geometric distance T'' of b' from the body's initial position $B \in S$. The Formula (5.2) digresses from the Newtonian (or universal) definition of the velocity as the ratio of distance travelled to the corresponding time interval, which takes in our units the form $\beta_U = T''/t$. In terms of the universal velocity, the inertial velocity takes the form

$$\beta = \frac{\beta_U}{\sqrt{\beta_U^2 + 1}} = \frac{U}{\sqrt{U^2 + c^2}} \quad (5.3)$$

The following comments illustrate some facts concerning the inertial velocity which applies, of course, only to material bodies, but not to light signals.

- (i) By (5.2), $0 \leq \beta < 1$, and hence the inertial velocity β cannot reach velocity of light.
- (ii) If a body b is at rest at $B \in S$ which is distinct from $b' \in S$, then b will never be found at b' . Setting $t = \infty$ in (5.2) yields $\beta = 0$.
- (iii) In spite of the fact that the inertial velocity of any object cannot reach the velocity of light, the object itself can overtake the pulse emanating from its starting position $B \in S$. If the moving body b and the light signal emitted from $B \in S$ when b passed by, arrive simultaneously at b' , then $t = T''$, and the inertial velocity of the body is $\beta = \frac{1}{\sqrt{2}} \approx 0.707$. However, the universal velocity of the same body is $\beta_U = T''/t = 1$.

While inertial velocity cannot reach the velocity of light, universal velocity is unbound. These facts demonstrate that there is nothing odd about the result of an experiment yielding a superluminal speed for an elementary particle.

(iv) For a fixed value of T'' , β is a decreasing function of t ; it tends to zero for t tending to infinity and to 1 for t tending to zero. This expresses the obvious fact, the faster the particle is the shorter time it takes to arrive at b' .

(v) For small velocities, $t/T'' \gg 1$ and

$$\beta = \frac{T''}{t} \frac{1}{\sqrt{1+(T''/t)^2}} \approx \frac{T''}{t} - \frac{1}{2} \left(\frac{T''}{t} \right)^3 = \beta_U - \frac{1}{2} \beta_U^3 \quad (5.4)$$

If the inertial velocity is sufficiently small we can neglect the third order term in comparison with the first order term and write

$$\beta \approx \frac{T''}{t} = \beta_U \quad (5.5)$$

We also obtain the same result simply by neglecting 1 in the dominator on the right hand-side of (5.2) in comparison with the much larger term $(t/T'')^2$. Therefore, for small values, the classical expression is an approximation of the inertial velocity Formula (5.2).

For high velocities, $t/T'' \ll 1$ and the Formula (5.2) can be approximated by

$$\beta \approx 1 - \frac{1}{2} \left(\frac{t}{T''} \right)^2 + \frac{3}{8} \left(\frac{t}{T''} \right)^4 = 1 - \frac{1}{2} \beta_U^{-2} + \frac{3}{8} \beta_U^{-4} \quad (5.6)$$

The inertial velocity can also be deduced in terms of T' . Indeed, on solving the expression of T' for β we obtain the following equivalent expression of the inertial velocity,

$$\beta = \sqrt{1 - (t/T')^2} \quad (5.7)$$

The μ -meson particles lifetime once more

The μ -meson particles whose mean lifetime is $t \approx 2 \times 10^{-6}$ sec, are generated at about 60 km above the earth surface, that is, $T'' \approx 60 \text{ km} / (300,000 \text{ km/sec}) = 2 \times 10^{-4}$ sec. Inserting $t/T'' = 10^{-2}$ in Formula (5.2) yields the inertial velocity of the particles that just arrive at the earth surface as

$$\beta = \frac{1}{\sqrt{1 + 10^{-4}}} \approx 1 - \frac{1}{2} 10^{-4} = 0.99995 \quad (5.8)$$

The particles with inertial velocities not less than the latter value can cover 60 km in the earth's frame S in spite of the fact that a pulse of light emitted from an S -observer can travel during the period of the μ -meson lifetime only

$$t \times c = 2 \times 10^{-6} \times 3 \times 10^5 \text{ km} = 0.6 \text{ km} \quad (5.9)$$

Relative to the earth frame the universal velocity of the mesons particles that reach the earth surface is

$$\beta_U = \frac{T''}{t} = \frac{2 \times 10^{-4}}{2 \times 10^{-6}} = 100 \quad (5.10)$$

one hundred times of the velocity of light!

6. The Simultaneous Positions of a Particle and Its Emitted Signal

Suppose that the body b heads towards O with inertial velocity β . A pulse emitted from b when at $B \in S$, say at $t = 0$, arrives at $O \in S$ simultaneously with the body arriving at $b' \in S$. The Galilean picture in which light is envisaged to emanate from the body's current position, corresponds to body-observer triangle (3.1), which reduces in the current head-on motion to a 3 straight segments (Figure 3), with

$$\frac{T''}{\beta} = \frac{T'}{1} = \frac{T}{1 + \beta} = \gamma t \quad (6.1)$$

Equivalently,

$$T'' = \gamma \beta t, \quad T' = \gamma t T = T'' + T' = \gamma (\beta + 1) t \quad (6.2)$$

The first and third ratios show that when the body is at a distance T'' from B , the pulse it emits is at a distance

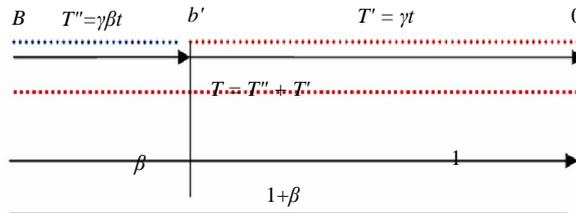


Figure 3. The simultaneous positions of the moving body and its emitted signal.

$$T = \frac{1+\beta}{\beta} T'' \quad (6.3)$$

from B . Thus, a particle always lags behind the pulse it emits in its direction. As an example, a *hypothetical* light signal emitted from the meson particles when generated, travels by the time during which a meson particle arrives at the earth surface the distance

$$T = \frac{1+0.99995}{0.99995} \times 2 \times 10^{-4} \text{ sec} \approx 4 \times 10^{10} \text{ sec} = 120 \text{ km}$$

which (assuming not absorbed) is almost twice as much the distance travelled by the particle itself.

By (6.1), the particle and the light, arrive at b' , at t_p and t_l respectively, where

$$t_p = \sqrt{1-\beta^2} \frac{T''}{\beta} \quad (6.4)$$

$$t_l = \sqrt{\frac{1-\beta}{1+\beta}} T'' = \sqrt{1-\beta^2} \frac{T''}{1+\beta} = \frac{\beta}{1+\beta} t_p \quad (6.5)$$

Thus light arrives first at b' advancing b by

$$t_p - t_l = t_p \left(1 - \frac{\beta}{1+\beta} \right) = \frac{t_p}{1+\beta} = \frac{t_l}{\beta} \quad (6.6)$$

Employing the STI to determine the position of the light front when b arrives at b' ; *i.e.* at the instant of time t_p , we get

$$T = \sqrt{\frac{1+\beta}{1-\beta}} t_p = \sqrt{\frac{1+\beta}{1-\beta}} \sqrt{1-\beta^2} \frac{T''}{\beta} = \frac{1+\beta}{\beta} T'' \quad (6.7)$$

which coincides with (6.3). The same relation has been obtained by the Galilean picture, which on scaling T'' and $T' = T - T''$ yields the particle travelling the distance $\sqrt{1-\beta^2} T''$ with velocity β and light travelling the distance $\sqrt{1-\beta^2} T'$ with velocity 1. *i.e.*

$$\sqrt{1-\beta^2} (T - T'') = t_p$$

Substituting for t_p from (6.4) we obtain T as given by (6.3). On changing to distances instead of geometric time distances we write (7.3) in the form

$$\frac{X}{c+u} = \frac{X''}{u}$$

which is the same as the classical picture, apart from the fact that the quotients are γt , but not t , and hence u refers to the inertial velocity. For small velocities, $X''/u \approx t$, and $X \approx (c+u)t$.

7. Conclusion

The UST which was shown to produce the mass-energy equivalence relation in a natural way proved also capable of producing a universal mechanics that is almost identical to the relativistic mechanics. Indeed, and through defining momentum as the product of mass and universal velocity, which has no upper bound, the mechanics constructed admits superluminal velocities but yet coincides with the relativistic mechanics in its basic dynamical components and their inter-relations. The possibility of superluminal velocities provides a straight forward explanation of the abundant presence of the μ -meson particles at the sea level in spite of their generation at high latitude and their short lifetime.

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