THz-Wave Generation from GaP THz Photonic Crystal Waveguides under Difference-Frequency Mixing

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ABSTRACT

GaP terahertz (THz) two-dimensional (2D) photonic crystal (PC) waveguides with line defects were fabricated using inductively-coupled plasma reactive-ion etching (ICP-RIE) in Ar/Cl_2 gas chemistries. THz-wave generation from the fabricated PC waveguides was demonstrated under collinear phase-matched difference-frequency generation (DFG), using Cr:Forsterite (Cr:F) lasers as the incident source. We compared the THz-wave output characteristics of the PC waveguides with that of GaP planar waveguides. The collinear phase-matching conditions in the DFG process were satisfied for 300- and 500-µm-wide PC waveguide structures at 0.7 and 0.6 THz, respectively. The additional output peaks that appeared near the edge of the photonic band gap, around 0.5 THz, were attributed to the guiding modes in the PC waveguide; no such peaks appeared in the non-patterned ridge waveguides. These experimental results suggest that the phonon-polariton confinement in THz-PC waveguides based on the GaP crystal could be used to enhance the nonlinear optical effect for THz-wave generation.

Keywords: THz-Wave Generation; Photonic Crystal Waveguide; Nonlinear Optical Effect

1. Introduction

Terahertz-wave (THz-wave) sources have attracted a great deal of attention because the THz frequency region has a number of important potential applications in the medical and imaging fields [1,2]. Many coherent THz sources, based on photoconductive antenna, quantum cascade lasers, and non-linear optical effects, such as difference-frequency generation (DFG), optical parametric oscillation, and optical rectification, have been developed for these applications [3-17]. In particular, DFG can be used to produce tunable, narrow linewidth, high-power THz-waves, with ns-pulsed and continuous wave (CW) operation at room temperature [18,19]. This DFG process is suitable for wave sources used in THz telecommunications and *in-situ* security screening.

A GaP crystal is an attractive nonlinear optical material for coherent THz-wave generation via the DFG process, due to its transparent properties in the infrared and THz regions and relatively high conversion efficiency due to THz phonon-polariton excitation [7-10,12-14]. Our previous works have demonstrated THz-wave generation from bulk GaP crystals under small-angle, noncollinear $\chi(2)$ phase-matched DFG conditions [7-9]. Using a THz-wave source, a frequency-tunable THz spectrometer was constructed to study the THz spectra of biomolecules (e.g. sugars, nucleosides, nucleotides, and amino acids) [20-22]. When the cross-section size of the GaP rod-waveguide structure coincides with the wavelength of the THz wave, the THz wave is confined in the waveguide, thus, inducing a change in the refractive index. As a result, THz waves are produced under a collinear $\chi(2)$ phase-matching configuration. We have reported THz-wave generation from GaP-THz rod and rib-waveguide structures and confirmed the enhancement of the conversion efficiency in the DFG process, using the wave-guiding effect for THz waves [23,24].

Recently, photonic crystal (PC) waveguides have been investigated because they allow the optical properties of the material to be modified. PC components have been developed in the optical region [25-28]. PCs in the THz region are also of interest because they could contribute to the development of THz technology. In contrast to passive devices, such as filters and PC fibers, few active devices have been investigated [29-33]. PC-embedded waveguides in the THz region make it possible to enhance the nonlinear optical effect, such as DFG in GaP, due to the strong phonon-polariton confinement in the PC waveguides.

In this paper, we fabricated two-dimensional (2D) photonic structures in GaP waveguides using inductivelycoupled plasma reactive-ion etching (ICP-RIE). The THz-wave output characteristics of the PC waveguides



were compared with those of GaP planar waveguides.

2. Experiment

GaP-PC waveguides were fabricated using the ICP-RIE technique on an undoped semi-insulating (001) GaP wafer (Sumitomo Metal Mining, Co. Ltd.; resistivity: >10⁵ W cm; thickness: 300 mm). The fabrication process involved spin-coating an XP KMPR-1025 negative photoresist onto a 60-µm-thick GaP wafer, patterned using UV (i-line) exposure. Using a Cl₂/Ar plasma chemistry in the ICP-IRE process, etch rates as high as 2 mm min⁻¹ were achieved to fabricate the PC waveguide structures. The waveguide samples were cleaved into 10×10 mm rectangles, along the <110> crystal direction. No antireflection coating was applied to the input or output surfaces of the waveguide. Figure 1 shows a scanning electron microscopy (SEM) image of one of the fabricated PC waveguides, with line defects in the G-K direction. Holes were arranged in a triangular lattice, with a lattice constant, a, of 200 mm; the radius of the air holes, r, was 82 mm. By removing rows of holes, we created 300- and 500-mm-wide line-defect waveguides along the G and K directions, respectively.

We calculated the photonic band diagram of the PC slab without line defects using the plane-wave expansion method (**Figure 2(a)**). The lowest stop band-gap region was estimated to exist in the frequency range of 0.4 to 0.55 THz for transverse electric (TE) polarization. We also calculated the THz transmission spectra for the THz-PC slab and waveguide using three-dimensional (3-D) finite-difference time-domain (3D-FDTD) simulations, shown in **Figure 2(b**). The modeled structures were the same as the PC without line defects and the 300-mm-wide PC waveguide. The THz-wave transmission spectra indicated high transmittance in the stop band-gap region for the PC waveguide. Thus, this suggests the existence of defect modes originating from the line defects in this frequency region.



Figure 1. SEM bird's-eye-view photograph of fabricated GaP photonic crystal (PC) waveguide with line defect widths of $300 \ \mu m$.



Figure 2. Calculated results of (a) photonic band diagram without line defects and (b) THz transmission spectra for the THz PC slab and waveguide (300-µm-wide linewidth).

Figure 3 illustrates a schematic of the experimental setup used for DFG of THz waves in GaP-PC waveguides, where Cr:Forsterite (Cr:F) lasers were used as the incident source. Both the pump and signal beams were delivered from two-channel Cr:F lasers (LOTIS TII) pumped with a double-pulse Q-switched Nd:YAG laser (repetition rate: 10 Hz, wavelength: 1064 nm). The wavelength of the pump source was fixed at 1216 nm; the signal source was variable from 1217 to 1221 nm, which corresponded to a frequency difference of 0.2 to 1.0 THz. The pulse energies of the pump and signal sources were adjusted to 0.5 mJ (pulse width: 22 ns; beam diameter: 1 mm). The pump and signal beam were collinearly combined using a beamsplitter and then directed to the input edge of the waveguide; the propagation direction of the two beams was parallel to the G-K direction. The generated THz wave was collected by a high-density polyethylene (HDPE) lens and detected by a liquid-helium-cooled Si-bolometer.

3. The Results and Discussions

THz-wave generation was carried out under collinear phase-matched DFG in GaP-PC waveguides. Figures 4(a) and (b) display the output characteristics of the generated THz wave from the 500- and 300-µm-GaP-PC waveguides, respectively. Figure 4(c) shows the THz-wave output from the GaP ridge waveguide, which had a thickness and width of 300 mm. The polarization of the generated THz wave is TE polarization. In the instances where the frequency of the signal source was varied at 1 GHz steps with respect to the fixed-frequency of the pump source, a 100 GHz bandwidth was observed at the peak frequency positions for phase matching, for both the 0.6 THz (corresponding to the 500-mm-wide waveguide) and 0.7 THz (300-um-wide waveguide) outputs. The peak frequency shift is due to the difference in the effective index for THz waves in the PC waveguides.



Figure 3. Schematic diagram of the experimental setup for THz-wave generation in GaP PC waveguides.





Figure 4. Frequency dependence of THz-wave output intensity from GaP-PC waveguides and non-patterned ridge waveguides with (a) line defect width of 500 μ m; (b) line defect width of 300 μ m; and (c) ridge width of 300 μ m. The insets in Figures 4(a) and (b) show enlarged view of the figures around 0.5 THz.

These bandwidths arise from the phase mismatch in the DFG process. The linewidth of THz wave is determined by that of incident infrared sources (500 MHz). Therefore, THz wave linewidth obtained by difference frequency process was about 1 GHz. The generated THz wave was tunable within the spectral bandwidth. The oscillation of the THz output power for both waveguides was observed due to the interference of the THz waves between the input and output faces of the waveguide. Therefore, we confirmed that the GaP-PC line-defect waveguides act as THz waveguides. For the 300-mmwide PC waveguide, a decrease in the THz output power appeared around 0.7 THz; this can be attributed to the transmission characteristics of the PC waveguide. The conversion efficiencies of generated THz waves from the waveguides were in the order of 10^{-4} %. This value is comparable as other nonlinear optical techniques such as optical parametric oscillation, optical rectification, and photoconductive antenna.

As mentioned above, the stop-band frequency was expected to be in the frequency range of 0.40 to 0.55 THz. For PC structures with line defects, the guiding modes should be shown in the stop-band frequency range. The insets of Figures 4(a) and (b) show the enlarged view around 0.5 THz. For the two PC waveguides, a THz output peak with narrow spectral bandwidth appeared around 0.5 THz, although phase matching was not satisfied in this frequency region. The intensity of the peaks for the 500-mm-PC waveguide was larger than that for the 300-mm-PC waveguide. Because no additional peak was observed in the phase-mismatched region for the non-patterned ridge waveguide (Figure 4(c)), the THz output characteristics were attributed to defect modes in the stop-band frequency region. The defect modes have a low group velocity, which corresponds to a large effective index of refraction, compared with non-patterned THz waveguides. This suggests that the modal dispersions of the defect modes contributed to the phasematching condition around 0.5 THz.

4. Conclusion

We have fabricated THz-GaP waveguides based on 2D PCs produced by reactive-ion etching in Ar/Cl₂ plasma chemistries. THz waves were generated from the fabricated PC waveguides with line defects. The collinear phase-matching conditions in the DFG process were satisfied for 300- and 500-mm-wide PC waveguide structures at 0.7 and 0.6 THz, respectively. Furthermore, the THz output peak with narrow spectral width appeared around 0.5 THz for both PC waveguides. This frequency region is within the stop-band (non phase-matching) region, as suggested by the numerical calculations. These THz output peaks were unique for these THz-PC waveguide structures. Due to the phonon-polariton confinement in GaP-PC line-defect waveguides, our experimental results indicated the possibility of enhancing the nonlinear optical effect for THz-wave generation using THz-PC waveguides based on the GaP crystal.

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