

Cluster Model of Formation of Subnuclear and Subatomic Objects

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ABSTRACT

The paper describes the development results on one-dimensional (1D) asymptotic model of the formation kinetics for the objects (clusters) of subnuclear (quark) and subatomic (nuclear) matters. A concept of the objects distribution density wave $\varphi(a, t)$ in space of sizes a lies in the basis for analytical description of the processes under consideration. The proposed formalism makes it possible to describe in an adequate way the final outcomes of the well-known catastrophic phenomena in the world of elementary particles. Mass characteristics of different processes of approach to equilibrium in nuclear reactions are calculated.

KEYWORDS

Cluster Model; Kinetics of Formation; Objects; Subnuclear/Nuclear Matter

1. Introduction

Different analytical and numerical methods are used to study the mechanism of reaching the equilibrium in nuclear processes (e.g., see [1-4]). Such investigations are based on the model concepts of the nature of nuclear matter and the dynamics of the processes under consideration. Work [5] describes an attempt to look at intranuclear phenomena from the unified point of view, treating them as processes of formation and growth of compact objects (clusters) with distinct collective quantum properties (strong interaction) in a closed stochastic system of nucleons oscillating around local positions of equilibrium. The closed system in question can be either a parent nucleus or two nuclei in contact after their colliding (a double system). At the initial moment of time, this system contains seeds of nuclear matter: alpha-particles and (or) their fragments — tritons and deuterons. The random interaction of these seeds results in their coalescence, and compact nuclear matter clusters are formed in the system. It is assumed that, as in the case of the growth processes of nanocrystals with strong interatomic bonds [6], oscillations of nucleons on the surface of contacting nuclear clusters with temporarily (virtually)

broken external bonds can lead to the mutual compensation of these unused bonds. As a result, consolidation of adjacent clusters takes place and larger objects of nuclear medium with different masses are formed. Due to Coulomb repulsion, these objects can escape from the area of interaction. Values of most propable mass numbers of clusters-nuclides calculated in [5] correspond to the well-known data for such processes of approach to equilibrium as cluster radioactivity, deep inelastic interaction of heavy ions, spontaneous nuclear fission and synthesis of superheavy elements in stars.

The aims of this work are as follows: 1) verification of the proposed cluster model using the example of hard collision processes of elementary particles and formation of stable hadron jets [7]; 2) consideration of the nuclide “stability islands” problem [8] on the basis of cluster model [5]; 3) determining average mass numbers of nuclei formed in the processes of spontaneous fission and nucleosynthesis in stars.

2. Analytical Approach

A concept of the objects distribution density wave $\varphi(a, t)$ in space of sizes a lies in the basis of analytical method

for describing the processes under consideration [5,6]. This 1D approach makes it possible to neglect deviations of the object geometric shape from the spherical one. An evolution of the wave $\varphi(a, t)$ during the stochastic process of the objects aggregation can be described in the diffusion approximation with the help of Fokker-Planck kinetic equation written for the space of clusters sizes a (where t is time):

$$\frac{\partial \varphi(a, t)}{\partial t} + \frac{\partial}{\partial a} [v\varphi(a, t)] - \frac{1}{2} \frac{\partial^2}{\partial a^2} [\eta\varphi(a, t)] = 0 \quad (1)$$

Here, $v = \langle da \rangle / dt$ and $\eta = \langle (da)^2 \rangle / dt = \hbar / 2m$ are the average rate of kinematic transfer of φ and the diffusion coefficient in the space a , respectively; m is the cluster mass; and \hbar is the reduced Planck constant. The carried out study of the asymptotic properties of function $\varphi(a, t)$ showed that in the end of the irreversible aggregation of objects as a result of interaction of large clusters with small seeds in a stochastic conservative system the distribution density of large clusters with $a \gg a_0$ (a_0 is a size of a seed) is inversely proportional to their masses: $\varphi \propto m^{-1}$ (an ‘‘inverse-mass’’ law) [6].

We define two mechanisms of growth of nuclear matter clusters: 1) the small flux of seeds, when each of them has time to occupy its place on the surface of the cluster before starting interaction with the next seed; 2) the high flux of seeds, when they affect the cluster almost simultaneously (equivalent to collision of clusters). Linearization of Equation (1) makes it possible to produce analytical expressions for increase in the mean size $\langle a \rangle$ of large clusters ($a \gg a_0$) with time [6]. In case of the first growth mechanism the approximate law of increasing the average size of clusters can takes the following form

$$\langle a \rangle_1 = a_0 (t/t_1)^{2/5}, \quad t_1 = (2\sqrt{2}/5) a_0 (m_0 t_0 / \hbar)^{1/2}. \quad (2)$$

Here, t_1 is the unit of time in the processes with small flux of seeds, t_0 is a typical time scale of objects interaction, m_0 is the seed mass. The second mechanism has a typical time scale equal to $t_0 = 2\langle a \rangle / c_0$, where $c_0 = 5 \times 10^7 \text{ m}\cdot\text{s}^{-1}$ is the ‘‘sound’’ velocity or average nucleon thermal velocity determined by the average thermal energy of the degenerate Fermi gas (22 MeV per nucleon [8,9]). Then, the law of increasing the average size can be written as

$$\langle a \rangle_2 = (3/2)^{1/3} a_0 (t/t_2)^{1/3}, \quad t_2 = a_0 (a_0 m_0 / \hbar c_0)^{1/2}. \quad (3)$$

Here, t_2 is a unit of time in the processes with a large flow of seeds.

The size of a spherical nucleus is related with mass number A in the following way [8,9]:

$$a \approx 2r_0 A^{1/3}. \quad (4)$$

Here, $r_0 = 1.3 \text{ fm}$ is a typical space scale of strong in-

teraction. The total time of the process is determined by the Heisenberg rule basing on the level of shell energy: [10]. It is correct for isolated excitation of level in the quantum system.

The following formula can be written for the most probable mass numbers of cluster-nuclides [5]:

$$A_n \approx A_0 \left(\frac{5}{2\sqrt{2}} \frac{\sqrt{\beta}}{\lambda} \pi n + 1 \right)^{6/5}, \quad n = 1, 2, 3, \dots \quad (5)$$

Here, A_0 is the mass number of the seed; $\lambda > 1$ is an arbitrary real number; and small parameter $\beta \approx \hbar t_0 / 2m_0 a_0^2$ is defined by seed mass m_0 and size a_0 , and by the characteristic time scale t_0 of the objects interaction. A reasonable choice for this scale is the period of high-frequency nucleon oscillations in the nucleus. This parameter can be defined as $t_0 = 2r_0/c_0 \approx 5 \times 10^{-23} \text{ s}$. The approximate quantity λ is found by matching the solutions for the first maximum of φ , which corresponds to the seeds, and the maximum of φ for small clusters [6]. The calculations yield $\lambda \approx \sqrt{15/2}$.

3. Subnuclear Processes

Verification of the suggested formalism can be performed using the example of catastrophic (deeply inelastic) processes with hadrons [7,11]. One can consider these processes as the following chain of events: head-on collisions of ‘‘lepton-nucleon’’ or ‘‘nucleon-nucleon’’ \rightarrow the quark is split from the gluon cloud \rightarrow random interactions in the continuous intrahadron medium \rightarrow formation of new particles (hadron jets). New particles are considered as clusters containing the elements of a continuous intrahadron medium (partons): quarks, gluons, quark-antiquark pairs, and so on. Then the law of evolution of the mean cluster size (3) gives the following relation between the units of length a_{unit} , time t_{unit} and mass m_{unit} at deep inelastic interaction of fundamental particles:

$$m_{\text{unit}} a_{\text{unit}}^3 / t_{\text{unit}}^2 = \hbar c \quad (6)$$

Here, c is maximum velocity of interaction propagation in hadron medium, namely, the velocity of light in vacuum. On the basis of this formula and other known data [7,11,12] for ‘‘conventional’’ quark masses (u- and d-quarks) deep inside hadrons and for typical space scales, it is possible to get the following evaluation of a time unit $t_{\text{unit}}^{\text{af}}$ for current quarks (cq) in the state of asymptotic freedom (af): $a_{\text{unit}}^{\text{af}} \sim 10^{-16} \text{ m}$, $m_{\text{unit}} \sim m_u^{\text{cq}} \approx 5 \text{ MeV}$, $m_{\text{unit}} \sim m_d^{\text{cq}} \approx 7 \text{ MeV} \Rightarrow t_{\text{unit}}^{\text{af}} \sim 10^{-26} \text{ s}$. Transition time t_{trans} for quarks from asymptotic freedom to constrained state (confinement) inside the hadron is estimated by formula (3) as $t_{\text{trans}} \sim t_{\text{unit}}^{\text{af}} (a_{\text{unit}}^{\text{conf}} / a_{\text{unit}}^{\text{af}})^3 \sim 10^{-23} \text{ s}$ ($a_{\text{unit}}^{\text{conf}} \sim 10^{-15} \text{ m}$). The produced value corresponds to the

time scale of strong interaction.

In the processes considered in [7], the minimum value of the length unit is $a_{\text{unit}}^{\text{min}} \sim 10^{-18}$ m. It can be taken as a top estimation of the current quark size. Then, from the relation (6) we get that $t_{\text{unit}}^{\text{min}} \sim 10^{-29}$ s, and the time $t_{\text{hadr}}^{\text{form}}$ of hadron formation with mean size $\langle a \rangle \sim 1$ fm, in compliance with formula (3), makes

$$t_{\text{hadr}}^{\text{form}} \approx t_{\text{unit}}^{\text{min}} \left(\langle a \rangle / a_{\text{unit}}^{\text{min}} \right)^3 \sim 10^{-20} \text{ s}.$$

The produced value is much higher than the life time of unstable hadrons (resonances) $t_{\text{life}}^{\text{res}} \leq 10^{-21}$ s : $t_{\text{hadr}}^{\text{form}} \gg t_{\text{life}}^{\text{res}}$. This reflects the formation of stable hadron jets in the processes of catastrophic collisions being described. If we take, in compliance with the results of [13], that the velocity of excitation propagation in quark-gluon medium is equal to the sound velocities $c_s = 0.3c$ and $0.57c$, then multipliers 1.8 and 1.3 will appear in the evaluations given above for time units, the time of quark transition into the bound state, and the time of hadron formation. These corrections will not change the estimation of time scale for strong interaction and will make stronger inequality $t_{\text{hadr}}^{\text{form}} \gg t_{\text{life}}^{\text{res}}$ that reflects the formation of stable hadron jets.

Asymptotic distribution $\varphi \propto m^{-1}$ means that: 1) the probability of nucleon fragmentation under deep inelastic scattering is higher than the chance of its preservation; 2) the number of pions in hadron jet is 1.5 times higher than the number of nucleons. These conclusions correspond to the notions on the nature of hard processes given in [7]. In particular, the following reaction takes place at frontal collision of protons: $pp \rightarrow \pi^0 \pi^+ \pi^+ nn$. One can see that there are 3 pions and 2 neutrons in the formed hadron jet.

Thus, the proposed formalism provides adequate determination of asymptotic states reached at well-known catastrophic phenomena in the world of subnuclear objects, namely, fundamental and elementary particles. This gives us a reason for making an attempt to apply this formalism for considering the intranuclear processes mentioned in the Section 1.

4. Processes in Nuclear Scales

The values A calculated in [5] correspond to the mass numbers of the nuclides and their isotopes over the currently known range, including the transfermium elements. In this paper we regard the problem of "stability islands" of nuclides [8,14]. As an example the element roentgenium with $Z = 111$ has the mass number of $A = 273$ [8]. In the proposed model, this value of A obtained from formula (5) corresponds to tritons as the seeds. At $Z = 114$, in compliance with 2β -stability, the number of neutrons is equal to 184 and the mass number is equal to $A = 298$ [14]. In our model, the mass numbers closed to it are 299 and 302. These values obtained from the formula (5)

correspond to tritons and deuterons as the seeds, respectively. Work [15], which is dedicated to the production and the decay of element 114, mentions its isotopes with $A = 288$ and 289. In the model, the values $A = 290$ and 293 are the closest mass numbers corresponded to the tritons and deuterons as the seeds, respectively. Using various calculation options of the β -stability band the islands of stability are predicted also at $Z = 164$ and with neutron numbers of 272 or 318 [14]. Respective mass numbers $A = 436$ and 482 are closed to the approximate value of the mass number of the final nuclide $A_{\text{end}} \approx 470$ calculated in [5]. In the model, mass number $A = 436$ corresponds to deuterons as the seeds. The closest to the mass number of 482 are the following values calculated with the help of formula (5): $A = 483$ for the tritons as the seeds, and $A = 485$ for deuterons as the seeds. Thus, one can presume that the developed asymptotic model of clusters formation in the nuclear matter complements the method for predicting nuclear mass with the help of radial basis function [4], which makes it possible to find mass numbers in the range from 20 to 260.

One can rewrite expressions (2) and (3) in the following way more convenient for calculations:

$$\langle a \rangle_1 \equiv \left(5/2\sqrt{2} \right)^{2/5} (6\hbar/\pi\rho t_0)^{1/5} t^{2/5}. \quad (7)$$

$$\langle a \rangle_2 \equiv (3/2)^{1/3} (6\hbar c_0/\pi\rho)^{1/6} t^{1/3}. \quad (8)$$

Here, ρ is the density of nuclear matter taken as equal to $\rho = 2.5 \times 10^{17}$ kg·m⁻³ [5]. In order to eliminate parameter a_0 from the formulas it was taken that the mass of the seed is equal to $m_0 \approx (\pi/6)\rho a_0^3$.

One can try to describe spontaneous nuclear fission as a result of excitation of the first rotational level with energy $\Gamma_{\text{rot}} = 100$ keV using formulas (4) and (8). The total time of the intranuclear process that corresponds to the specified value of Γ_{rot} is equal to $\tau_{\text{rot}} = 6.287 \times 10^{-21}$ s, and the average mass number of light fragments is evaluated to be $\langle A \rangle_{\text{light}} \approx 100$. Average mass number of heavy fragments $\langle A \rangle_{\text{heavy}}$ is determined by the mass conservation law $A_{\text{mat}} = \langle A \rangle_{\text{light}} + \langle A \rangle_{\text{heavy}}$, A_{mat} is the mass number of parent nucleus. In case of light actinides (Th, U) we get $\langle A \rangle_{\text{heavy}} \approx 135 - 138$. The diagrams of mass distribution for the fragments of spontaneous fission of heavy nuclei ($A = 235, 238$) given in [16] show that the average mass number of lighter fragments is 90 - 100, and the average mass number of heavier fragments is about 140.

If we presume that high-speed nucleosynthesis in stars takes place as a result of transition of the system of nucleons from the high-frequency vibration level with the period of $t_0 = 2r_0/c_0 \approx 5 \times 10^{-23}$ s to the first rotational level with the energy of 100 keV during the life time $\tau_{\text{rot}} = 6.287 \times 10^{-21}$ s of this level, then formulas (4) and (7)

give the evaluation of the average mass number of superheavy elements as $\langle A \rangle_{\text{superheavy}} \approx 330$. As for the forced fission competed with the nucleosynthesis, in this case it is possible to take the travel time $t = a/c_0$ of the sound wave in the nucleus formed as a result of nucleosynthesis as a time scale t_0 . With regard to formula (4) when $\langle A \rangle = 330$ we get $t_0 = 3.6 \times 10^{-22}$ s. Then from the formulae (4) and (7) we get that, when the lifetime of the first rotational level is 6.287×10^{-21} s, the average mass number of light fragments of superheavy elements is about $\langle A \rangle_{\text{light}} \approx 100$. In compliance with the mass conservation law the average mass number of heavy fragments is equal to $\langle A \rangle_{\text{heavy}} \approx 230$. The produced values approximately agree with the first and the last peaks of the final element abundance in the Galaxy [2]. This agreement gives the credibility to the above-given evaluation of the average mass number of superheavy elements, which can be formed as a result of nucleosynthesis in supernovae and neutron stars.

5. Discussion of Results

The proposed model corresponds to the problem of motion of 1D wave packet $\phi(a, t)$ in the space a [6]. The group velocity $\Psi = d\langle a \rangle/dt$ goes down with time:

$$d\langle a \rangle/dt \propto t^{Z-1}, \quad 0 < Z \leq 1.$$

This circumstance reflects the expansion of wave packet and the slower-down of its propagation. At the final stage of irreversible aggregation of clusters in the closed system this wave packet with non-Gaussian shape has an attenuating discontinuity in the wave “front” related to the maximum possible size of a_{max} :

$$\phi \propto a^{-3}, \quad \langle a \rangle \ll a \leq a_{\text{max}}; \quad \phi = 0, \quad a > a_{\text{max}}.$$

The presented results and notions allow us to make a conclusion that the developed in [6] asymptotic method for describing the formation of clusters under consideration has sufficiently general character to be implemented in respective problems of high-energy physics.

In Section 3 it was found that the least value of time unit in quark-gluon medium is 10^{-29} s, and the unit of time for current quarks in the state of asymptotic freedom is 10^{-26} s. Using relation (6) one can get the following evaluation of the unit of time in the interaction processes with constituent quarks in confinement state described in [11,12]: $a_{\text{unit}}^{\text{conf}} \sim 10^{-15}$ m, $m_{\text{unit}} = m_u \cong m_d \cong 300 \text{ MeV} \Rightarrow t_{\text{unit}}^{\text{conf}} \sim 4 \times 10^{-24}$ s. The transfer time from asymptotic freedom to confinement is evaluated in Section 3 as $t_{\text{trans}} \sim 10^{-23}$ s that agrees with the time scale of strong interaction, and the formation time of stable hadrons is estimated as $t_{\text{form}}^{\text{hadr}} \sim 10^{-20}$ s. Relation (6) shows that the processes in the world of fundamental and elementary particles are characterized with the spectrum of time

units. Close to the upper boundary this spectrum overlaps with the “lower” area of times $t = 10^{-23} - 10^{-22}$ s which are typical for direct nuclear reactions. Thus, formula (6), which shows that “regular” space-and-time relations are valid up to the distances of about $10^{-18} - 10^{-15}$ m and time of $10^{-29} - 10^{-26}$ s, complies with the generally accepted notions about space-and-time scales in microphysics [17].

We can try to determine the value of phenomenological “fundamental mass” of m_{fund} , if we take that the least space unit (a fundamental length) is the value of $a_{\text{fund}} \sim 10^{-18}$ m and that this value is connected with the fundamental time scale t_{fund} by the reasonable relation $t_{\text{fund}} = a_{\text{fund}}/c$. Then, from relation (6) we get the following expression for the fundamental mass: $m_{\text{fund}} = \hbar/ca_{\text{fund}}$. Then we get that $m_{\text{fund}}c^2 = \hbar c/a_{\text{fund}} = 196 \text{ GeV}$ [6]. The produced value lies within the supposed experimental range specified in [8], where mass m_{H^0} of the neutral Higgs boson H^0 should be: $115 < m_{H^0} < 250 \text{ GeV}$. Besides, this value is closed to the “critical” mass of 180 - 200 GeV, above which H^0 -boson can decay into the pairs of W- and Z-bosons [18]. The reason of the obtained agreement is in the fact that the proposed model of formation of “quantum” clusters is scalar and does not contain such parameters of spins and of electric charges, and, therefore, it can be applied to the scalar neutral boson H^0 with the zero spin. It is also important to mark approximate correspondence of obtained evaluation for m_{fund} to the value of the upper limit of the Higgs mass identified in [19] as 170 GeV.

It could be interesting to carry out a qualitative analysis of various processes with fundamental and elementary particles with the help of relation (6). For example, we could try to estimate the state of the heaviest of all quarks — t -quark (tq), which has a conventional mass of 176 GeV [12]. One can get from relation (6) that the minimum time unit for t -quark, when having the size $a_{\text{unit}}^{\text{min}} \sim 10^{-18}$ m, is $t_{\text{unit}}^{\text{tq}} \sim 10^{-27}$ s. Then, as it comes from formula (3), the time $t_{\text{trans}} \sim 10^{-21}$ s would be required to form bounded states with a characteristic size of about 10^{-15} m. The produced evaluation complies with the life time of unstable resonances, therefore, t -quark does not form the stable hadron. This result complies with the common notion, according to which a t -quark is the “only quark that is ‘born and dies free’” [12].

In the area of nuclear scales the suggested method allows to evaluate the time of approach to equilibrium basing on the typical mass numbers and respective nuclide sizes. This problem was solved in [20] with regard to the research on the kinetics of deep-inelastic interaction between the beam of copper ions and the gold target at the collision energy of 365 MeV [9]. It was shown that the formation time for the interaction products with the average mass numbers $\langle A \rangle = 60$ and 100, typical for the

process under consideration, is about of 10^{-20} s. The produced value is much less than the lifetimes of 10^{-16} - 10^{-14} s which are typical for the intermediate compound nucleus. This makes it possible to speak about comparatively quick (explosion) character of the process of deep inelastic heavy-ions interaction.

6. Conclusion

The developed cluster model makes it possible to produce adequate evaluations of space-and-time and mass characteristics for the processes in the subnuclear (quark) matter and for the intranuclear processes of approach to equilibrium.

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