

Estimates of Tritium Produced Ratio in the Blanket of Fusion Reactors

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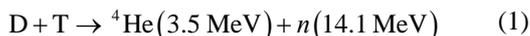
ABSTRACT

For the preparation of tritium fuel as the main and rare fuel of reactors in the fusion reactors, the reactor blanket must be designed so that it provides enough tritium breeding ratio. The tritium breeding ratio, TBR, in the blanket of reactors should be greater than one, ($TBR > 1$), by applying lithium blanket. The calculations for proposed parameters (t_d, f_b, η and t_p), indicate that the estimated tritium breeding ratio is greater than one. The calculated $TBR = 1.04$ satisfies the tritium provision condition.

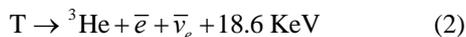
Keywords: Tritium Breeding Ratio; Reactor Blanket; Lithium; Fusion

1. Introduction

The first generation of controlled fusion devices reactive for the release of energy reaction is the following:

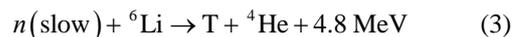


The reaction free energy is 17.6 MeV, about 80 and 20 percent of the energy carried by neutrons and alpha particles respectively [1]. More neutrons are produced in the blanket of lithium where they produce tritium. The neutrons may escape through the blanket or are absorbed by structural material. Tritium is naturally limited resources and also radioactive. Radioactive tritium spontaneously decays to ${}^3\text{He}$, an electron \bar{e} and an antineutrino $\bar{\nu}_e$, with a half life of 12.3 years in process of beta-decay:

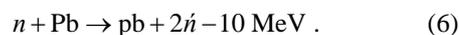
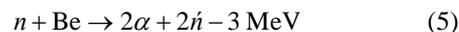


Unlike the stable isotope deuterium, which makes up 156 ppm of hydrogen on earth, tritium has a relatively short shelf life because of the radioactive decay (Equation (2)), so tritium is most efficiently used a few years after its manufacture. Consequently, its production requires special management. For a fusion reaction to be economically profitable, tritium breeding ratio (TBR), should be greater than one. The tritium breeding ratio is defined as the average number of tritium atoms bred per tritium atom burn in the reaction 1. One of methods can be achieved in fusion reactor TBR greater than one, using of lithium containing blankets with neutron multiplier. The interaction of neutrons with lithium blanket with pure tri-

tritium will be produced as follows [2,3]:



The lithium abundance in Equations (3) and (4) are 7.5 and 92.5 respectively. With proper selection of materials blanket structure and geometry, the loss of neutrons can be minimized by absorption or escape from the blanket. Nearly all of the neutrons that slow down to thermal energies from their initial 14.1 MeV, can be absorbed by ${}^6\text{Li}$ and can generate tritium. Another way to produce large tritium breeding ratio is to include neutron multipliers such as beryllium and lead, in lithium blankets. Neutron multiplier will occur as the following:



The variation of cross section versus energy for the important reactions for tritium breeding are shown in **Figure 1**.

2. Theory

Tritium mass consumption rate can be estimated for a reactor with the power P_F GW, during a year as follows:

$$\frac{dM_1}{dt} = \frac{A}{N_0} \times \frac{dN}{dt} \approx 56 P_F \text{ Kgy}^{-1} \text{ GWT}^{-1} \quad (7)$$

A and N_0 are Tritium mass numbers and Avogadro's number respectively. According to Equation (7), it 56 Kg

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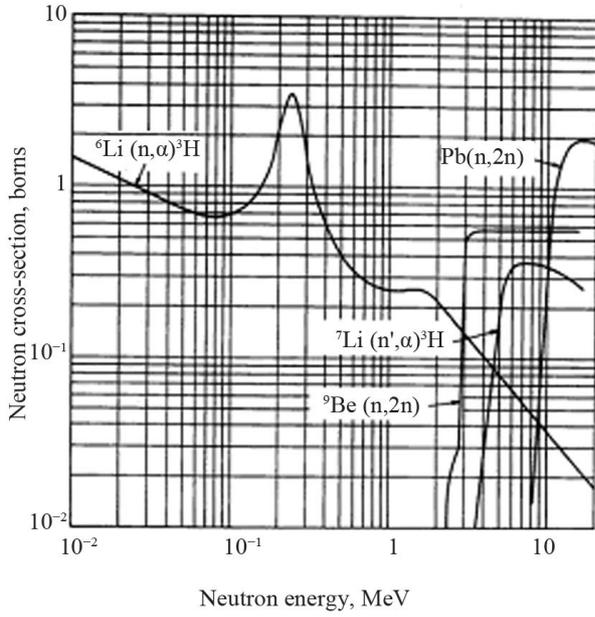


Figure 1. The variation of cross sections versus of energy for the important reactions for tritium breeding [2].

of tritium should be burned per year for producing 1 GW energy. The burnup fraction consumption parameter, (f_b), is required for estimating the tritium. Also, The tritium is not consumed, is collected and is used in the burning cycle. For reactors with the steady state power, the values of the tritium breeding ratio should be closed to one. According to reaction Equation (1), the time dependence of number densities of tritium and deuterium in the plasma are given as following [4,5]:

$$\frac{dN_D}{dt} = \frac{dN_T}{dt} = -N_D N_T \langle \sigma v \rangle, \quad (8)$$

where reaction rate, $\langle \sigma v \rangle$, is given by [6]:

$$\langle \sigma v \rangle = \frac{3.68 \times 10^{-18}}{T_i^{\frac{2}{3}}} \exp\left(\frac{-19.94}{T_i^{\frac{1}{3}}}\right) \quad (9)$$

where T_i is ion temperatures in keV. Suppose the reactor is stoichiometrically fueled so $N_D = N_T = N$, Then one can readily solve Equation (7) to find that the number density N_{τ^*} at time $\tau^* > 0$ is related to the number density N_0 at the time of fuel injection by:

$$N_{\tau^*} = \frac{N_0}{1 + \langle \sigma v \rangle N_0 \tau^*} \quad (10)$$

τ^* is effective confinement time:

$$\tau^* = \frac{\tau_C}{1 - R} \quad (11)$$

τ_C is the ion confinement time, R is fraction of the ions that escape the plasma can recombine. Finally, the tri-

tium burnup fraction is:

$$f_b = 1 - \frac{N_{\tau^*}}{N_0} = \frac{\langle \sigma v \rangle N_0 \tau^*}{1 + \langle \sigma v \rangle N_0 \tau^*} \quad (12)$$

Figure 2 shows the burn up fraction consumption over time. Fraction of fuel consumption depends on three factors τ^* , $\langle \sigma v \rangle$ and N_0 . By considering the tritium radioactive decay rate, $\gamma_s = \frac{\text{Ln}2}{12.3} y^{-1}$ and the loss rate γ_r in the reprocessing of unburn tritium, the inventory of tritium is given by:

$$M_0 = \frac{\frac{dM_1}{dt}}{TBR \times \frac{\eta f_b}{t_p} - \gamma_s - \gamma_r} \quad (13)$$

where t_p , is a mean time to clean up or recycle the tritium and η is improved efficiency of tritium injected into the plasma. if tritium production rate, $\left(TBR \times \frac{\eta f_b}{t_p}\right)$,

is larger than the tritium loss rate, $\gamma_r + \gamma_s$, and amount TBR to be considered one, the mass inventory can approximate by:

$$M_0 \approx \frac{t_p \frac{dM_1}{dt}}{\eta f_b} \quad (14)$$

The Fusion reactor blanket must be designed so that the lost tritium in radioactive decay can be produced and reconstructed [7]. In addition, after the doubling time, t_d , enough extra tritium should be produced to provide the initial inventory for an identical reactor. The total inventory of tritium in the reactor is:

$$M_1 = M_0 + m \quad (15)$$

where m , is produced mass in the reactor blanket. Radio-

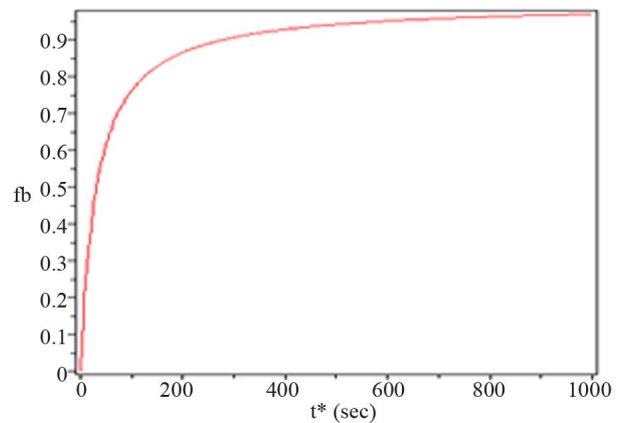


Figure 2. Burnup fraction to effective confinement time with $T_i = 20$ keV and $N_0 = 10^{20} \text{ m}^{-3}$.

active decay at the rate, γ_s , and losses at the rate, γ_r in the reprocessing loop, burning and breeding will cause the refueling mass to change at the rate:

$$\frac{dm}{dt} = \frac{dM_1}{dt} \times TBR - \gamma_s m - (\gamma_s + \gamma_r) M_0 - \frac{dM_1}{dt} \quad (16)$$

In Equation (16), the First term is tritium production rate, Second term is the decay rate of tritium in the blanket, the third sentence is the decay rate and initial mass loss of tritium and the last sentence is tritium mass consumption rate. To calculate TBR the net production rate coefficient, (κ), is required. By replacing $\frac{dM_1}{dt} = \frac{M_0 \eta f_b}{t_p}$,

from Equation (14) in Equation 16, the Equation (16) can be rewritten as:

$$\frac{dm}{dt} = -\gamma_s m + \left(-\gamma_s - \gamma_r + \frac{\eta f_b}{t_p} + TBR \frac{\eta f_b}{t_p} \right) M_0 \quad (17)$$

Using Equation (17), the coefficient of production rate can be calculated as:

$$\kappa = \frac{\eta f_b}{t_p} (TBR - 1) - \gamma_s - \gamma_r \quad (18)$$

Then Equation (17), can be rewritten as:

$$\frac{dm}{dt} = \kappa M_0 - \gamma m \quad (19)$$

If there is no breeding in blanket, $m = 0$ and $\frac{dm}{dt} = 0$,

Equation (16) gives Equation (13). The time depended mass in the reactor blanket can be calculated Equation (19) under the assumption of a positive tritium production rate, ($\kappa > 0$), and assuming $m = 0$ at time $t = 0$, as;

$$m(t) = \frac{KM_0}{\gamma_s} (1 - e^{\gamma_s t}) \quad (20)$$

and at the doubling time, $t_d > 0$, we have:
 $m(t = t_d) = M_0$:

$$m(t = t_d) = M_0 = \frac{KM_0}{\gamma_s} (1 - e^{\gamma_s t_d}) \quad (21)$$

Then the coefficient of net production rate will be defined as;

$$\kappa = \frac{\gamma_s}{1 - e^{-\gamma_s t_d}} \quad (22)$$

For simplicity, consider sufficiently short doubling times, t_d (a few years), that there is negligible radioactive decay of the tritium ($\gamma_s = 0.05 y^{-1}$) then $\gamma_s t_d \ll 1$. Equation (22) can be approximated as follows:

$$\kappa = \frac{\gamma_s}{1 - (1 - \gamma_s t_d + \dots)} = \frac{1}{t_d} \quad (23)$$

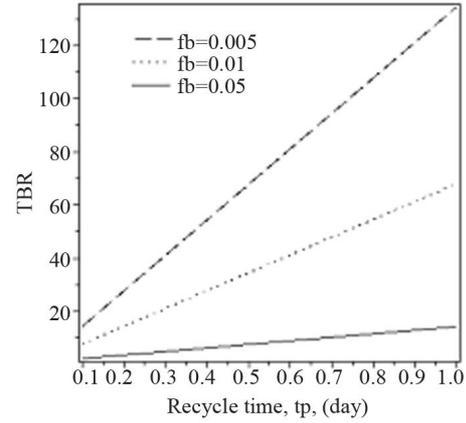


Figure 3. Tritium breeding ratio to recycle time with different of burn up fraction consumption parameters ($f_b = 0.05$, $f_b = 0.01$ and $f_b = 0.1$) and $\eta = 0.5$.

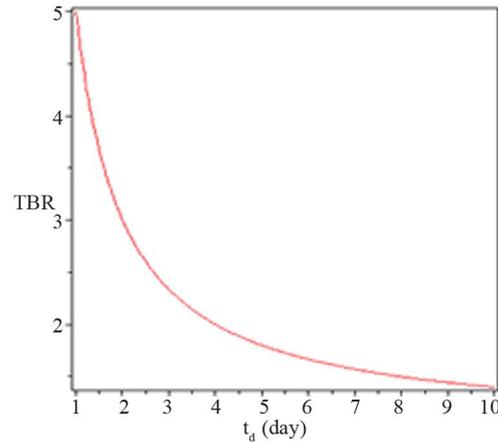


Figure 4. Tritium breeding ratio to doubling time with parameters $f_b = 0.05$, $\eta = 0.5$ and $t_p = 1$ day .

the tritium production rate is larger than the radioactive decay rate and tritium loss rate. Then we can approximate Equation (22) by $\kappa = \frac{1}{t_d}$ and neglect γ_s and γ_r

in Equation (18) to find that the required tritium breeding rate is:

$$TBR - 1 \approx \frac{t_p}{\eta f_b t_d} \quad (24)$$

3. Conclusion

In this scheme, the ratio of tritium breeding in reactors with lithium blanket is estimated by considering of the possible reactions in the plasma and reactor blanket, Tritium breeding ratio based on the recycle time and doubling time are plotted in **Figures 3** and **4**. **Figure 3** shows that TBR increases by the increasing of recycle time with different of burn up fraction consumption pa-

rameters ($f_b = 0.05, f_b = 0.01, f_b = 0.1$ and $\eta = 0.5$). Also, it can be seen that with increasing the doubling time, tritium breeding ratio is always greater than one (**Figure 4**). For example, with parameters $f_b = 0.05, \eta = 0.5$ and $t_p = 1$ day with a doubling time, $t_d = 3$ year, the tritium breeding ratio is approximately 1.04. This amount will secure the condition of fuel supply, then for a reactor design with $TBR > 1$, it is necessary that the doubling time, the tritium burn fraction and the injection efficiency must be longer and the recycling time must be smallest.

REFERENCES

- [1] J. Yu and G. Yu, "Fission-Fusion Neutron Source," *Journal of Nuclear Material*, Vol. 386-388, 2009, pp. 949-953.
- [2] T. Nakagawa, H. Kawasaki and K. Shibata, "Curves and Tables of Neutron Cross Section in JENDL-3.3(Part I and II)," Japan Atomic Energy Research Institute, Fukushima, 2002.
- [3] C. Q. Sh and Y. X. Ln, "Tritium and Materials for Protection of Tritium Permeation," Atomic Energy Publishing House, Beijing, 2002.
- [4] A. Harms, K. F. Schoepf, G. H. Miley and D. R. Kingdon, "Principles of Fusion Energy," World Scientific Publishing Co. Pte. Ltd., Singapore City, 2002.
- [5] S. Atzeni and J. Meyer-ter-Vehn, "Inertial Fusion," Oxford University Press, Oxford, 2004.
- [6] S. Glasstone and R. Loveberg, "Controlled Thermonuclear Reactions," van Norstrand, Princeton, 1960.
- [7] M. Nishikawa, "Tritium Balance in a D-T Fusion Reactor," *Fusion Science and Technology*, Vol. 59, No. 2, 2011, pp. 350-362.