

Real Eigenvalue of a Non-Hermitian Hamiltonian System

Ram Mehar Singh

Department of Physics, Chaudhary Devi Lal University, Sirsa, India

Email: dixit_rammehar@yahoo.co.in, rmd06@rediffmail.com

Received August 12, 2012; revised September 12, 2012; accepted September 20, 2012

ABSTRACT

With a view to getting further insight into the solutions of one-dimensional analogous Schrödinger equation for a non-hermitian (complex) Hamiltonian system, we investigate the quasi-exact \mathcal{PT} -symmetric solutions for an octic potential and its variant using extended complex phase space approach characterized by $x = x_1 + ip_2$, $p = p_1 + ix_2$, where (x_1, p_1) and (x_2, p_2) are real and considered as canonical pairs. Besides the complexity of the phase space, complexity of potential parameters is also considered. The analyticity property of the eigenfunction alone is found sufficient to throw light on the nature of eigenvalue and eigenfunction of a system. The imaginary part of energy eigenvalue of a non-hermitian Hamiltonian exist for complex potential parameters and reduces to zero for real parameters. However, in the present work, it is found that imaginary component of the energy eigenvalue vanishes even when potential parameters are complex, provided that \mathcal{PT} -symmetric condition is satisfied. Thus \mathcal{PT} -symmetric version of a non-hermitian Hamiltonian possesses the real eigenvalue.

Keywords: Analogous Schrödinger Equation; Complex Hamiltonian; \mathcal{PT} -Symmetry

1. Introduction

In the recent years, one-dimensional complex Hamiltonians $H(x, p)$ have generated lot of interest for the understanding of several newly discovered phenomena in various science-streams [1,2], but such studies in mathematical terms have not been reached to the desired extent. Several attempts have been made to obtain the solutions of Schrödinger equation (SE) for different anharmonic potentials in real domain. However, the study of complex octic potential has become of considerable interest due to the peculiar nature of the eigenvalue spectrum. Further, besides some general studies of complex Hamiltonians in nonlinear domain [1,3], efforts have been made to study both classical and quantum aspects [4-7] of a system. At the classical context, $H(x, p)$ becomes the function of two complex variables and the analyticity property of $H(x, p) = H_1(x_1, x_2, p_1, p_2) + iH_2(x_1, x_2, p_1, p_2)$ leads to a class of integrable systems in the associated two-dimensional real systems H_1 and H_2 , where, H_1 act as new Hamiltonian and H_2 is a second integral of motion. The possible connection between H_1 and H_2 is sought in terms of anti-Bäcklund transformation [5]. In the quantum context, as $p = -i\hbar \frac{\partial}{\partial x}$ which implies

$$p_1 = \frac{-1}{2} \frac{\partial}{\partial p_2} \quad \text{and} \quad x_2 = \frac{-1}{2} \frac{\partial}{\partial x_1}, \text{ the analyticity of}$$

$H(x, p)$ is translated into complex potential $V(x)$. While a complex Hamiltonian is no longer hermitian and ordinarily does not guarantee for real eigenvalues, however, in \mathcal{PT} -symmetric version [8-10], the system is found to exhibit real and bounded eigenvalue spectrum. The reality of the spectrum is a consequence of combined action of the parity and time reversal invariance of Hamiltonian [6]. Recently, following the work of C. M. Bender *et al.* [8,9], one-dimensional Hamiltonian systems have been studied rigorously through combined parity and time reversal operators. The parity operator \hat{P} and time reversal operator \hat{T} defined by the action of position and momentum operators are $\hat{P}:(x, p, i) \rightarrow (-x, -p, i)$; $\hat{T}:(x, p, i) \rightarrow (x, -p, -i)$. The combined action of parity-time operator is

$$\hat{P}\hat{T}:(x, p, i) \rightarrow (-x, p, -i), \quad (1)$$

where, $\hat{P}^2\hat{T}^2 = 1$. Here, the operators \hat{x} and \hat{p} are real, the commutator $[x, p] = i$ is invariant under operators \hat{P} and \hat{T} . It is interesting to note that commutation relation still remain invariant even if \hat{x} and \hat{p} becomes complex, provided that above transformation hold. There are various ways of complexifying a given Hamiltonian [11], but here we use the scheme given by Xavier and de Aguir [12,13], used to develop an algorithm for the computation of semiclassical coherent state propagator to transform potentials in extended complex

phase space approach (ECPSA). The real and imaginary parts of x and p are introduced as

$$x = \text{Re } x + i \text{Im } x, \quad p = \text{Re } p + i \text{Im } p,$$

if, we define $x_1 = \text{Re } x$, $x_2 = \text{Im } x$, $p_1 = \text{Re } p$, $x_2 = \text{Im } p$, then x and p can be defined as

$$x = x_1 + ip_2, \quad p = p_1 + ix_2. \tag{2}$$

The presence of variables (x_1, p_2, x_2, p_1) in the above transformations may be regarded as some sort of co-ordinate momentum interactions of a dynamical system. Note that, in this complexifying scheme, the degrees of freedom of the underlying system just become double. The \mathcal{PT} -symmetric condition for the above transformation becomes

$\hat{\mathcal{P}}\hat{\mathcal{T}} : (x_1, p_2, x_2, p_1, i) \rightarrow (-x_1, p_2, -x_2, p_1, -i)$. Though complex potentials are in practice for a long time, such as in optical model of nucleus, delocalization transition in condensed matter system—such as vortex flux line dippening in type-II superconductors, yet the quantum mechanics of complex potentials has not been studied to a desired level. It is since last few years that the study of complex potentials has become important enough for better theoretical understanding of the detailed properties of some newly discovered phenomena in physics and chemistry, like the phenomena pertaining to resonance scattering in atomic, molecular, and nuclear physics and to some chemical reactions [14-17]. The complex Hamiltonian is used in several other theoretical context like studies of complex trajectories with regard to the calculation of semiclassical coherent-state propagator in the path integral method have attracted particular interest in laser physics [12,13]. The \mathcal{PT} -symmetric non-hermitian Hamiltonians have many applications in various fields of physics-like superconductivity, population biology, quantum cosmology, condensed matter physics, quantum field theory etc.

Transformations similar to Equation (2) have also been used in the study of nonlinear evolution equations in context of amplitude-modulated nonlinear Langmuir waves in plasma [4]. Recently, in some studies, solutions of the Schrödinger wave equation have been reported using the extended complex phase space approach (ECPSA) [11,18,19]. With this motivation and to expand the domain of applications, we investigate the quasi-exact solution of the analogous Schrödinger equation (ASE) for a coupled complex octic potential and its variant in one dimension.

The paper is organized as follows: in Section 2, we are devoted with the mathematical formulation of the ECPSA for computing the ground state and excited state eigenvalue spectra of some one-dimensional complex systems. Under the same mathematical prescription, ground state solutions are presented in Section 3 and excited state solutions are described in Section 4. Finally,

concluding remarks are presented in Section 5.

2. General Results

For a complex Hamiltonian system $H(x, p)$ in one dimension, the ASE (for $\hbar = m = 1$) is given by

$$\hat{H}(x, p)\psi(x) = E\psi(x), \tag{3}$$

where

$$\hat{H}(x, p) = -\frac{1}{2} \frac{d^2}{dx^2} + V(x). \tag{4}$$

Here, Equation (3) departs from the conventional and mathematical setting of the standard Schrödinger equation [20], so Equation (3) is termed as analogous Schrödinger equation (ASE) for a non-hermitian operator $H(x, p)$. The transformation condition (2) implies that

$$\frac{d}{dx} = \frac{1}{2} \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial p_2} \right), \quad \frac{d}{dp} = \frac{1}{2} \left(\frac{\partial}{\partial p_1} - i \frac{\partial}{\partial x_2} \right). \tag{5}$$

Note that, the momentum operator $p = -i\hbar \frac{d}{dx}$ of the conventional quantum mechanics under the transformation (2) reduces to $p_1 + ix_2 = \frac{-i}{2} \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial p_1} \right)$. This relation yields $p_1 = \frac{-1}{2} \frac{\partial}{\partial p_1}$, $x_2 = \frac{-1}{2} \frac{\partial}{\partial x_1}$. Also, the complex co-ordinate transformation (2) preserves the fundamental commutation relations, $[x, p] = i$, which can be easily verified with the help of Equations (2) and (5). To express the ASE (3) into a pair of coupled partial differential equation, the complex forms of $V(x), \psi(x)$ and E are written as

$$V(x) = V_r(x_1, p_2) + iV_i(x_1, p_2), \tag{6a}$$

$$\psi(x) = \psi_r(x_1, p_2) + i\psi_i(x_1, p_2), \tag{6b}$$

$$E = E_r + iE_i, \tag{6c}$$

where, subscripts “ r ” and “ i ” denote the real and imaginary parts of the corresponding quantities and other subscripts to these quantities separated by comma will denote the partial derivatives of the quantity concerned. Thus, after inserting Equations (2), (4) and (6a)-(6c) in Equation (3) and separating the real and imaginary parts in the final expression, one gets the following pair of partial differential equations

$$-\frac{1}{8}(\psi_{r,x_1x_1} - \psi_{r,p_2p_2} + 2\psi_{i,x_1p_2}) + V_r\psi_r - V_i\psi_i = E_r\psi_r - E_i\psi_i, \tag{7a}$$

$$-\frac{1}{8}(\psi_{i,x_1x_1} - \psi_{i,p_2p_2} - 2\psi_{r,x_1p_2}) + V_r\psi_i + V_i\psi_r = E_r\psi_i + E_i\psi_r. \tag{7b}$$

The analyticity property of the wavefunction $\psi(x)$, in terms of Cauchy-Riemann conditions, implies

$$\psi_{r,x_1} = \psi_{i,p_2}, \quad \psi_{r,p_2} = -\psi_{i,x_1}. \tag{8}$$

Under the analyticity condition (8), Equations (7a) and (7b), reduces to

$$-\frac{1}{2}\psi_{r,x_1x_1} + V_r\psi_r - V_i\psi_i = E_r\psi_r - E_i\psi_i, \tag{9a}$$

$$-\frac{1}{2}\psi_{i,x_1x_1} + V_r\psi_i + V_i\psi_r = E_r\psi_i + E_i\psi_r. \tag{9b}$$

The ansatz for the wavefunction $\psi(x)$ is taken as [19]

$$\psi(x) = \phi(x) \exp[g(x)], \tag{10}$$

where, $\phi(x)$ and $g(x)$ are the polynomial functions of the complex variable x , which can be expressed as

$$\phi(x_1, p_2) = \phi_r(x_1, p_2) + i\phi_i(x_1, p_2), \tag{11a}$$

$$g(x_1, p_2) = g_r(x_1, p_2) + ig_i(x_1, p_2). \tag{11b}$$

After utilizing Equations (6b), (9b), (11a) and (11b) in Equation (10), the real and imaginary parts of the wavefunction are expressed as

$$\psi_r(x_1, p_2) = e^{g_r} (\phi_r \cos g_i - \phi_i \sin g_i) \tag{12a}$$

$$\psi_i(x_1, p_2) = e^{g_r} (\phi_i \cos g_i + \phi_r \sin g_i). \tag{12b}$$

In view of the analyticity condition (8), g_r and g_i satisfies the relations

$$g_{r,x_1} = g_{i,p_2}; \quad g_{r,p_2} = -g_{i,x_1}. \tag{13}$$

Therefore, with the help of Equations (12a) and (12b), the Equations (9a) and (9b) yield

$$g_{r,x_1x_1} + (g_{r,x_1})^2 - (g_{i,x_1})^2 + \frac{1}{(\phi_r^2 + \phi_i^2)} \left[\phi_r (\phi_{r,x_1x_1} + 2\phi_{r,x_1} g_{r,x_1} - 2\phi_{i,x_1} g_{i,x_1}) \right. \tag{14a}$$

$$\left. + \phi_i (\phi_{i,x_1x_1} + 2\phi_{r,x_1} g_{i,x_1} + 2\phi_{i,x_1} g_{r,x_1}) \right]$$

$$+ 2(E_r - V_r) = 0,$$

$$g_{i,x_1x_1} + 2g_{r,x_1} g_{i,x_1}$$

$$+ \frac{1}{(\phi_r^2 + \phi_i^2)} \left[\phi_r (\phi_{i,x_1x_1} + 2\phi_{r,x_1} g_{i,x_1} + 2\phi_{i,x_1} g_{r,x_1}) \right. \tag{14b}$$

$$\left. + \phi_i (-\phi_{r,x_1x_1} + 2\phi_{i,x_1} g_{i,x_1} - 2\phi_{r,x_1} g_{i,x_1}) \right]$$

$$+ 2(E_i - V_i) = 0.$$

It is to be noted that for given functional forms of $\phi(x)$ and $g(x)$, the rationalization of Equations (14a) and (14b) yield the real and imaginary components of the energy eigenvalue spectrum for the excited state of a

system. On the other hand, if $\phi(x)$ is chosen as constant, then Equations (14a) and (14b) reduces to ground state solutions as

$$g_{r,x_1x_1} + (g_{r,x_1})^2 - (g_{i,x_1})^2 + 2(E_r - V_r) = 0, \tag{15a}$$

$$g_{i,x_1x_1} + 2g_{r,x_1} g_{i,x_1} + 2(E_i - V_i) = 0. \tag{15b}$$

With a suitable ansatz for $g(x)$, satisfying the analyticity condition (19), the rationalization of Equations (15a) and (15b), provides ground state solutions of the ASE for a given complex potential.

3. Ground State Solutions

Here, we look for the ground state solutions of one-dimensional complex octic potential and its variant as:

3.1. Generalized Octic Potential

Consider a generalized octic potential of the form

$$V(x, y) = a_1x + a_2x^2 + a_3x^3 + a_4x^4 + a_5x^5 + a_6x^6 + a_7x^7 + a_8x^8, \tag{16}$$

where, the coupling parameters $(a_k, 1 \leq k \leq 8)$ are complex *i.e.* $a_k = a_{kr} + ia_{ki}$ and a_{kr}, ia_{ki} are constants.

Under the \mathcal{PT} -symmetric condition (1), the potential (24) reduces to

$$V(x, y) = a_{1i}x + a_{2r}x^2 + a_{3i}x^3 + a_{4r}x^4 + a_{5i}x^5 + a_{6r}x^6 + a_{7i}x^7 + a_{8r}x^8. \tag{17}$$

By implying the transformation (2) on the potential (17), the real and imaginary parts of the potential turn out to be

$$V_r = -a_{1i}p_2 + a_{2r}(x_1^2 - p_2^2) + a_{3i}(p_2^3 - 3x_1^2p_2) + a_{4r}(x_1^4 + p_2^4 - 6x_1^2p_2^2) + a_{5i}(-5x_1^4p_2 - p_2^5 + 10x_1^2p_2^3) + a_{6r}(x_1^6 - p_2^6 + 15x_1^2p_2^4 - 15x_1^4p_2^2) + a_{7i}(p_2^7 - 21x_1^2p_2^5 - 7x_1^6p_2 + 35x_1^4p_2^3) + a_{8r}(x_1^8 - p_2^8 - 28x_1^2p_2^6 - 28x_1^6p_2^2 + 70x_1^4p_2^4), \tag{18}$$

$$V_i = a_{1i}x_1 + 2a_{2r}x_1p_2 + a_{3i}(x_1^3 - 3x_1p_2^2) - 4a_{4r}(x_1p_2^3 - x_1^3p_2) + a_{5i}(5x_1p_2^4 + x_1^5 - 10x_1^3p_2^2) - a_{6r}(-6x_1p_2^5 - 6x_1^5p_2 + 20x_1^3p_2^3) + a_{7i}(x_1^7 - 7x_1p_2^6 - 21x_1^5p_2^2 + 35x_1^3p_2^4) - a_{8r}(8x_1p_2^7 - 8x_1^7p_2 + 56x_1^5p_2^3 - 56x_1^3p_2^5). \tag{19}$$

The polynomial forms of $g_r(x_1, p_2)$ and $g_i(x_1, p_2)$, in conformity with (13) are written as

$$g_r = \frac{1}{2}\alpha_2(x_1^2 - p_2^2) - \frac{1}{3}\beta_3(p_2^3 - 3x_1^2 p_2) + \frac{1}{4}\alpha_4(x_1^4 + p_2^4 - 6x_1^2 p_2^2) - \frac{1}{5}\beta_5(-5x_1^4 p_2 - p_2^5 + 10x_1^2 p_2^3), \tag{20a}$$

$$g_i = \alpha_2 x_1 p_2 - \frac{1}{3}\beta_3(x_1^3 - 3x_1 p_2^2) - \alpha_4(x_1 p_2^3 - x_1^3 p_2) - \frac{1}{5}\beta_5(5x_1 p_2^4 + x_1^5 - 10x_1^3 p_2^2), \tag{20b}$$

where, α_i and β_i are real. Now, inserting the above forms of $g_r(x_1, p_2)$ and $g_i(x_1, p_2)$ in Equations (15a) and (15b), the rationalization of the resultant expression yields the following set of non-repeating equations

$$E_r = -\frac{1}{2}\alpha_2, \tag{21a}$$

$$E_i = 0, \tag{21b}$$

$$\beta_3 = -a_{1i}, \tag{21c}$$

$$\alpha_2^2 + 3\alpha_4 = 2a_{2r}, \tag{21d}$$

$$2\beta_5 + \alpha_2\beta_3 = -a_{3i}, \tag{21e}$$

$$-\beta_3^2 + 2\alpha_2\alpha_4 = 2a_{4r}, \tag{21f}$$

$$\beta_5\alpha_2 + \beta_3\alpha_4 = -a_{5i} \tag{21g}$$

$$\alpha_4^2 - 2\beta_3\beta_5 = 2a_{6r}, \tag{21h}$$

$$\alpha_4\beta_5 = -a_{7i} \tag{21i}$$

$$\beta_5^2 = -2a_{8r}. \tag{21j}$$

Here, Equations (21c)-(21j) except (21e) give rise to the constraining relations among the potential parameters. However, the Equations (21e), (21h) and (21j) can be immediately solved for four arbitrary constants *i.e.* α_2 , β_3 , α_4 and β_5 . Whereas Equations (21d)-(21i) can be solved for some negative values of a_{8r} say $a_{8r} = -\bar{a}_{8r}$ in the potential (17). The results obtained are

$$\alpha_2 = \pm \frac{4\bar{a}_{8r}[-a_{3i}\sqrt{2\bar{a}_{8r}} \mp 4\bar{a}_{8r}]}{4\bar{a}_{8r}a_{6r} - a_{7i}^2}, \tag{22}$$

$$\beta_3 = \pm \frac{1}{\sqrt{2\bar{a}_{8r}}}\left[a_{6r} - \frac{a_{7i}^2}{4\bar{a}_{8r}}\right],$$

$$\alpha_4 = \mp \frac{a_{7i}}{\sqrt{2\bar{a}_{8r}}}, \beta_5 = \pm \sqrt{2\bar{a}_{8r}}, \tag{23}$$

where, \bar{a}_{8r} is real positive. The constraining relations obtained from Equations (21c), (21d), (21f) and (21g) are

$$\begin{aligned} \pm(4\bar{a}_{8r}a_{6r} - a_{7i}^2) &= -2a_{1i}(2\bar{a}_{8r})^{3/2}, \\ \frac{16\bar{a}_{8r}^2[-a_{3i}\sqrt{2\bar{a}_{8r}} \pm 4\bar{a}_{8r}]^2}{4\bar{a}_{8r}a_{6r} - a_{7i}^2} \mp \frac{3a_{7i}}{\sqrt{2\bar{a}_{8r}}} &= 2a_{2r}, \\ \mp \frac{1}{2\bar{a}_{8r}}\left[a_{6r} - \frac{a_{7i}^2}{4\bar{a}_{8r}}\right] & \\ \mp \frac{2}{\sqrt{2\bar{a}_{8r}}}a_{7i}\left[\frac{4\bar{a}_{8r}[-a_{3i}\sqrt{2\bar{a}_{8r}} \pm 4\bar{a}_{8r}]}{4\bar{a}_{8r}a_{6r} - a_{7i}^2}\right] &= 2a_{4r}, \\ \pm 2(2\bar{a}_{8r})^{3/2}\left[\frac{-a_{3i}\sqrt{2\bar{a}_{8r}} \mp 4\bar{a}_{8r}}{4\bar{a}_{8r}a_{6r} - a_{7i}^2}\right] & \\ \mp \frac{a_{7i}}{2\bar{a}_{8r}}\left[a_{6r} - \frac{a_{7i}^2}{4\bar{a}_{8r}}\right] &= -a_{5i}. \end{aligned}$$

The presence of these constraining relations, makes the problem quasi solvable. Such relations can be helpful in definition and approximate sub domain in complex parametric space in which a given complex potential provides real spectra. As from Equation (21b), imaginary part of the energy eigenvalue is zero, while the real part of the energy eigenvalue obtained from Equation (21a) turns out to be

$$E_r^{(0)} = \frac{\mp 2\bar{a}_{8r}[a_{3i}\sqrt{2\bar{a}_{8r}} + 4\bar{a}_{8r}]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)}. \tag{24}$$

The corresponding eigenfunction becomes

$$\begin{aligned} \psi^{(0)}(x) &= \exp\left[-\frac{2\bar{a}_{8r}[a_{3i}\sqrt{2\bar{a}_{8r}} + 4\bar{a}_{8r}]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)}x^2\right. \\ &\quad \left.- \frac{i[4\bar{a}_{8r}a_{6r} - a_{7i}^2]}{6(2\bar{a}_{8r})^{3/2}}x^3 + \frac{a_{7i}}{4\sqrt{2\bar{a}_{8r}}}x^4 \mp \frac{i\sqrt{2\bar{a}_{8r}}}{5}x^5\right] \end{aligned} \tag{25}$$

3.2. Variant of Octic Potential

Here, we consider one-dimensional octic potential (16) with the inverse harmonic term as

$$V(x, y) = a_1x + a_2x^2 + a_3x^3 + a_4x^4 + a_5x^5 + a_6x^6 + a_7x^7 + a_8x^8 + \frac{A}{x^2}, \tag{26}$$

where, the potential parameters a_i and A are complex constants. By implying the \mathcal{PT} -symmetric condition (2) on the potential (26), one gets

$$V(x) = a_{1i}x + a_{2r}x^2 + a_{3i}x^3 + a_{4r}x^4 + a_{5i}x^5 + a_{6r}x^6 + a_{7i}x^7 + a_{8r}x^8 + \frac{A_r}{x^2}. \tag{27}$$

Implying the transformation (2), the real and imaginary parts of the potential (26) are written as

$$V_r = V_{r1} + A_r \frac{(x_1^2 - p_2^2)}{(x_1^2 + p_2^2)^2}, \tag{28a}$$

$$V_i = V_{i1} - A_r \frac{(2x_1 p_2)}{(x_1^2 + p_2^2)^2}, \tag{28b}$$

where, V_{r1} and V_{i1} are same as given by Equations (18) and (19). The functional form of $g_r(x_1, p_2)$ and $g_i(x_1, p_2)$ complying with the analyticity condition (13) are written as

$$g_r = \frac{1}{2}\alpha_2(x_1^2 - p_2^2) - \frac{1}{3}\beta_3(p_2^3 - 3x_1^2 p_2) + \frac{1}{4}\alpha_4(x_1^4 + p_2^4 - 6x_1^2 p_2^2) - \frac{1}{5}\beta_5(-5x_1^4 p_2 - p_2^5 + 10x_1^2 p_2^3) - \delta \frac{1}{2} \ln(x_1^2 + p_2^2), \tag{29a}$$

$$g_i = \alpha_2 x_1 p_2 - \frac{1}{3}\beta_3(x_1^3 - 3x_1 p_2^2) - \alpha_4(x_1 p_2^3 - x_1^3 p_2) - \frac{1}{5}\beta_5(5x_1 p_2^4 + x_1^5 - 10x_1^3 p_2^2) + \delta \tan^{-1}\left(\frac{x_1}{p_2}\right). \tag{29b}$$

As before, using these forms of g_r and g_i in Equations (15a) and (15b), the rationalization of the resultant expression yield a set of equations in addition to Equations (21f)-(21j) as

$$E_r = a_r - \frac{1}{2}\alpha_2 + \alpha_2 \delta, \tag{30a}$$

$$E_i = 0, \tag{30b}$$

$$\beta_3(1 - \delta) = -a_{i1}, \tag{30c}$$

$$\alpha_2^2 + 3\alpha_4 - 2\alpha_4 \delta = 2a_{2r}, \tag{30d}$$

$$\alpha_2 \beta_3 + 2\beta_5 - \beta_5 \delta = -a_{3i}, \tag{30e}$$

$$\delta + \delta^2 = 2A_r. \tag{30f}$$

The pair of Equations (30e) and (30f) immediately lead us to

$$\delta = \frac{-1 \pm \sqrt{1+8A_r}}{2},$$

$$\alpha_2 = \mp \frac{4\bar{a}_{8r}}{4\bar{a}_{8r}a_{6r} - a_{7i}^2} \left[a_{3i} \sqrt{2\bar{a}_{8r}} \mp (5 \mp \sqrt{1+8A_r}) \bar{a}_{8r} \right]. \tag{31a}$$

The additional constraining relations given by Equations (30c) and (30d) are

$$\mp \left[4\bar{a}_{8r}a_{6r} - a_{7i}^2 \right] (3 \mp \sqrt{1+8A_r}) = -4a_{i1} (2\bar{a}_{8r})^{3/2},$$

$$2(2\bar{a}_{8r})^{5/2} \left[a_{3i} \sqrt{2\bar{a}_{8r}} \mp (-5 \mp \sqrt{1+8A_r}) \bar{a}_{8r} \right]^2$$

$$\pm (4 \mp \sqrt{1+8A_r}) a_{7i} = a_{2r} (4\bar{a}_{8r}a_{6r} - a_{7i}^2) \sqrt{2\bar{a}_{8r}}.$$

Under the similar prescription as adopted in previous case, the imaginary component of the energy eigenvalue vanishes, whereas the real component of the energy is

$$E_r^{(0)} = \frac{\mp 2(-2 \pm \sqrt{1+8A_r}) \bar{a}_{8r} \left[-a_{3i} \sqrt{2\bar{a}_{8r}} \mp (-5 \pm \sqrt{1+8A_r}) \bar{a}_{8r} \right]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)}. \tag{32}$$

The ground state eigenfunction for the potential (44) turn out to be

$$\psi^{(0)}(x) = (x_1^2 + p_2^2)^{-\frac{\delta}{2}} \exp \left[\mp \frac{2\bar{a}_{8r} \left[a_{3i} \sqrt{2\bar{a}_{8r}} \mp (5 \pm \sqrt{1+8A_r}) \bar{a}_{8r} \right]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)} x^2 \right]$$

$$\mp \frac{i \left[4\bar{a}_{8r}a_{6r} - a_{7i}^2 \right]}{6(2\bar{a}_{8r})^{3/2}} x^3 + \frac{a_{7i}}{4\sqrt{2\bar{a}_{8r}}} x^4$$

$$\pm \frac{i\sqrt{2\bar{a}_{8r}}}{5} x^5 + i\delta \tan^{-1}\left(\frac{x_1}{p_2}\right). \tag{33}$$

4. Excited State Solutions

Here, we elaborate viability of the prescription laid down in Sections 2 and 3 to compute eigenvalue and corresponding eigenfunction for the first excited state. The functional form of $\phi(x)$ for the first excited state is taken as

$$\phi(x) = \alpha x + \eta. \tag{34}$$

Then, under the transformation (2), the above equation reduces to

$$\phi_r(x_1, p_2) = \alpha x_1 + \eta, \phi_i(x_1, p_2) = \alpha p_2, \tag{35}$$

where, α and η are considered as real constants. In order to compute the corresponding energy eigenvalue and eigenfunction for the first excited state of potential (16), we use the same functional forms of g_r and g_i as mentioned in Equations (20a) and (20b). After inserting the Equations (20a), (20b) and (35) in Equations (14a) and (14b), then equating the coefficients of x_1, p_2 and their various products to zero, one gets the following set of non-repeating equations in addition to Equations (21f)-(21j)

$$E_r^{(1)} = a_r - \frac{3}{2}\alpha_2, \tag{36a}$$

$$E_i^{(1)} = 0, \tag{36b}$$

$$2\beta_3 = -a_{1i}, \tag{36c}$$

$$5\alpha_4 + \alpha_2^2 = 2a_{2r}, \tag{36d}$$

$$\alpha_2\beta_3 + 3\beta_5 = -a_{3i}. \tag{36e}$$

Inserting Equations (21h)-(21j) in Equation (36e), we have

$$\alpha_2 = \mp \frac{4a_{8r} [a_{3i}\sqrt{2\bar{a}_{8r}} \mp 6\bar{a}_{8r}]}{4\bar{a}_{8r}a_{6r} - a_{7i}^2}. \tag{37}$$

The other potential parameters are same as described in earlier section, whereas the constraining relations obtained from Equations (36c) and (36d) are

$$4\bar{a}_{8r}a_{6r} - a_{7i}^2 = -a_{1i} (2\bar{a}_{8r})^{3/2},$$

$$\frac{16\bar{a}_{8r}^2 (-a_{3i}\sqrt{2\bar{a}_{8r}} \mp 6\bar{a}_{8r})^2}{4\bar{a}_{8r}a_{6r} - a_{7i}^2} \mp \frac{5a_{7i}}{\sqrt{2\bar{a}_{8r}}} = 2a_{2r} (4\bar{a}_{8r}a_{6r} + a_{7i}^2).$$

Using the various ansatz parameters in Equation (36a), the real component of energy eigenvalue is written as

$$E_r^{(1)} = \frac{6\bar{a}_{8r} [a_{3i}\sqrt{2\bar{a}_{8r}} + 6\bar{a}_{8r} 2]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)}, \tag{38}$$

whereas the eigenfunction is given by

$$\psi^{(1)}(x) = \alpha x \exp \left[-\frac{2\bar{a}_{8r} [a_{3i}\sqrt{2\bar{a}_{8r}} + 6\bar{a}_{8r}]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)} x^2 - \frac{i [4\bar{a}_{8r}a_{6r} - a_{7i}^2]}{6(2\bar{a}_{8r})^{3/2}} x^3 + \frac{ia_{7i}}{4\sqrt{2\bar{a}_{8r}}} x^4 - \frac{i\sqrt{2\bar{a}_{8r}}}{5} x^5 \right] \tag{39}$$

Variant of Octic Potential

Again to compute energy eigenvalue and corresponding eigenfunction for the first excited state of potential (26), we use the same functional forms of g_r and g_i as mentioned in Equations (29a) and (29b). Then implying Equations (20a), (20b) and (35) in Equations (14a) and (14b), the rationalization of the final expression yields the following set of non-repeating equations in addition to Equations (21f)-(21j)

$$E_r = a_r - \frac{3}{2}\alpha_2 + \alpha_2\delta, \tag{40a}$$

$$E_i = 0, \tag{40b}$$

$$\beta_3(2 - \delta) = -a_{1i}, \tag{40c}$$

$$\alpha_2^2 + 5\alpha_4 - 2\alpha_4\delta = 2a_{2r}, \tag{40d}$$

$$\alpha_2\beta_3 + 3\beta_5 - \beta_5\delta = -a_{3i}, \tag{40e}$$

$$-\delta + \delta^2 = 2A_r. \tag{40f}$$

The Equations (40e) and (40f) lead us to

$$\delta = \frac{1 \pm \sqrt{1 + 8A_r}}{2}, \tag{41a}$$

$$\alpha_2 = \frac{-4\bar{a}_{8r}}{4\bar{a}_{8r}a_{6r} - a_{7i}^2} [a_{3i}\sqrt{2\bar{a}_{8r}} \mp (5 - \sqrt{1 + 8A_r})\bar{a}_{8r}]. \tag{41b}$$

The other ansatz parameters are same as for ground state solutions. However, the additional constraining relations given by Equations (40c) and (40d) are

$$\left[4\bar{a}_{8r}a_{6r} - a_{7i}^2 \right] \left(\frac{3\mp\sqrt{1+8A_r}}{2} \right) = \mp 4a_{1i} (2\bar{a}_{8r})^{3/2},$$

$$2(2\bar{a}_{8r})^{5/2} [a_{3i}\sqrt{2\bar{a}_{8r}} \mp (5\mp\sqrt{1+8A_r})\bar{a}_{8r}]^2$$

$$\mp (4\mp\sqrt{1+8A_r})(4\bar{a}_{8r}a_{6r} - a_{7i}^2)a_{7i}$$

$$= 2a_{2r}\sqrt{2\bar{a}_{8r}}(4\bar{a}_{8r}a_{6r} - a_{7i}^2).$$

Under the similar prescription as in previous case, the energy eigenvalue and eigenfunction are given by

$$E_r^{(1)} = \frac{(-2 + \sqrt{1 + 8A_r})\bar{a}_{8r} [a_{3i}\sqrt{2\bar{a}_{8r}} + (5 - \sqrt{1 + 8A_r})\bar{a}_{8r}]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)}, \tag{42}$$

$$\psi^{(1)}(x) = (x_1^2 + p_2^2)^{-i(1 \pm \sqrt{1 + 8A_r})}$$

$$\exp \left[-\frac{2\bar{a}_{8r} [a_{3i}\sqrt{2\bar{a}_{8r}} + (5 - \sqrt{1 + 8A_r})\bar{a}_{8r}]}{(4\bar{a}_{8r}a_{6r} - a_{7i}^2)} x^2 - \frac{i [4\bar{a}_{8r}a_{6r} - a_{7i}^2]}{6(2\bar{a}_{8r})^{3/2}} x^3 - \frac{a_{7i}}{4\sqrt{2\bar{a}_{8r}}} x^4 - \frac{i\sqrt{2\bar{a}_{8r}}}{5} x^5 + \frac{i(1 \pm \sqrt{1 + 8A_r})}{2} \times \tan^{-1} \left(\frac{x_1}{p_2} \right) \right]. \tag{43}$$

5. Concluding Remarks

In the present work, we have investigated quasi-exact \mathcal{PT} -symmetric solutions of the ASE for one-dimensional octic potential and its variants using ECPSA. Besides complexity of the phase space produced by transformation (2), complexity of the potential parame-

ters is also taken into account and ground state as well as excited states solutions are worked out. It is also emphasized that solutions of the ASE in the above said cases are obtained only in the presence of certain constraining relations among potential parameters, such constraining relations give rise to bound states of a system. It is found that imaginary part of the energy eigenvalue always vanishes for the solvable case of ASE, as long as all potential parameters are real. However, for \mathcal{PT} -symmetric potentials, energy eigenvalues are found real, even if concerned potentials possess complex parameters. The interesting aspect of this method is an account of complex coupling coefficients of potential in addition to complex phase space. Thus present method suggests another degree of freedom to obtain the real spectra for non-hermitian operator.

6. Acknowledgements

The author expresses his gratitude to Prof. S. C. Mishra and Dr. Fakir Chand, Department of Physics, Kurukshetra University, Kurukshetra (India), for their valuable suggestions regarding the manuscript.

REFERENCES

- [1] H. Feshbach, C. E. Porter and V. F. Weisskopf, "Model for Nuclear Reactions with Neutrons," *Physical Review*, Vol. 96, No. 2, 1954, pp. 448-464. [doi:10.1103/PhysRev.96.448](https://doi.org/10.1103/PhysRev.96.448)
- [2] R. S. Kaushal, "Classical and Quantum Mechanics of Noncentral Potentials," Narosa Publishing House, New Delhi, 1998.
- [3] F. Verheest, "Nonlinear Wave Interaction in a Complex Hamiltonian Formalism," *Journal of Physics A: Mathematical and General*, Vol. 20, No. 1, 1987, pp. 103-110. [doi:10.1088/0305-4470/20/1/019](https://doi.org/10.1088/0305-4470/20/1/019)
- [4] N. N. Rao, B. Buti and S. B. Khadkikar, "Hamiltonian Systems with Indefinite Kinetic Energy," *Pramana: Journal of Physics*, Vol. 27, No. 4, 1986, pp. 497-505.
- [5] R. S. Kaushal and H. J. Korsch, "Some Remarks on Complex Hamiltonian Systems," *Physics Letters A*, Vol. 276, No. 1-4, 2000, pp. 47-51. [doi:10.1016/S0375-9601\(00\)00647-2](https://doi.org/10.1016/S0375-9601(00)00647-2)
- [6] R. S. Kaushal and S. Singh, "Construction of Complex Invariants for Classical Dynamical Systems," *Annals of Physics*, Vol. 288, No. 2, 2001, pp. 253-276. [doi:10.1006/aphy.2000.6108](https://doi.org/10.1006/aphy.2000.6108)
- [7] C. M. Bender and A. Turbiner, "Analytic Continuation of Eigenvalue Problems," *Physics Letters A*, Vol. 173, No. 6, 1993, pp. 442-446. [doi:10.1016/0375-9601\(93\)90153-Q](https://doi.org/10.1016/0375-9601(93)90153-Q)
- [8] C. M. Bender and S. Boettcher, "Real Spectra in Non-Hermitian Hamiltonians Having \mathcal{PT} -Symmetry," *Physical Review Letters*, Vol. 80, No. 24, 1998, pp. 5243-5246. [doi:10.1103/PhysRevLett.80.5243](https://doi.org/10.1103/PhysRevLett.80.5243)
- [9] C. M. Bender, S. Boettcher and P. N. Meisinger, " \mathcal{PT} -Symmetric Quantum Mechanics," *Journal of Mathematical Physics*, Vol. 40, No. 5, 1999, pp. 2201-2229. [doi:10.1063/1.532860](https://doi.org/10.1063/1.532860)
- [10] F. M. Fernandez, R. Gujardiola, J. Ross and M. Zonjil, "Strong Coupling Expansion for the \mathcal{PT} -Symmetric Oscillator $V(x) = a(ix) + b(ix)^2 + c(ix)^3$," *Journal of Physics A: Mathematical and General*, Vol. 31, No. 50, 1998, pp. 10105-10112. [doi:10.1088/0305-4470/31/50/008](https://doi.org/10.1088/0305-4470/31/50/008)
- [11] R. S. Kaushal, "On the Quantum Mechanics of Complex Hamiltonian Systems in One Dimension," *Journal of Physics A: Mathematical and General*, Vol. 34, No. 49, 2001, pp. L709-L714. [doi:10.1088/0305-4470/34/49/104](https://doi.org/10.1088/0305-4470/34/49/104)
- [12] A. L. Xavier Jr. and M. A. M. de Aguiar, "Complex Trajectories in the Quartic Oscillator and Its Semiclassical Coherent-State," *Annals of Physics*, Vol. 252, No. 2, 1996, pp. 458-476. [doi:10.1006/aphy.1996.0141](https://doi.org/10.1006/aphy.1996.0141)
- [13] A. L. Xavier Jr. and M. A. M. de Aguiar, "Phase Space Approach to the Tunnel Effect: A New Semiclassical Traversal Time," *Physical Review Letters*, Vol. 79, No. 18, 1997, pp. 3323-3326. [doi:10.1103/PhysRevLett.79.3323](https://doi.org/10.1103/PhysRevLett.79.3323)
- [14] T. J. Hollowood, "Solitons in Affine Toda Theories," *Nuclear Physics B*, Vol. 384, No. 3, 1992, pp. 523-540. [doi:10.1016/0550-3213\(92\)90579-Z](https://doi.org/10.1016/0550-3213(92)90579-Z)
- [15] D. R. Nelson and N. M. Shnerb, "Non-Hermitian Localization and Population Biology," *Physical Review E*, Vol. 58, No. 2, 1998, pp. 1383-1403. [doi:10.1103/PhysRevE.58.1383](https://doi.org/10.1103/PhysRevE.58.1383)
- [16] N. Hatano and D. R. Nelson, "Localization Transitions in Non-Hermitian Quantum Mechanics," *Physical Review Letters*, Vol. 77, No. 3, 1996, pp. 570-573. [doi:10.1103/PhysRevLett.77.570](https://doi.org/10.1103/PhysRevLett.77.570)
- [17] N. Hatano and D. R. Nelson, "Vortex Pinning and Non-Hermitian Quantum Mechanics," *Physical Review B*, Vol. 56, No. 14, 1997, pp. 8651-8673. [doi:10.1103/PhysRevB.56.8651](https://doi.org/10.1103/PhysRevB.56.8651)
- [18] R. S. Kaushal and Parthasarathi, "Quantum Mechanics of Complex Hamiltonian Systems in One Dimension," *Journal of Physics A: Mathematical and General*, Vol. 35, No. 41, 2002, pp. 8743-8761. [doi:10.1088/0305-4470/35/41/308](https://doi.org/10.1088/0305-4470/35/41/308)
- [19] Parthasarathi and R. S. Kaushal, "Quantum Mechanics of Complex Sextic Potential in One Dimension," *Physica Scripta*, Vol. 68, No. 2, 2003, pp. 115-127. [doi:10.1238/Physica.Regular.068a00115](https://doi.org/10.1238/Physica.Regular.068a00115)
- [20] L. I. Schiff, "Quantum Mechanics," Tata McGraw-Hill Publishing Company Limited, New York, 1968.